

1 **Variability of the eastward currents in the equatorial**  
2 **Atlantic during 1993-2011**

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15 To be submitted to Journal of Geophysical Research  
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**Abstract**

We develop, validate and apply a synthetic method to monitor the equatorial currents in the Atlantic. This method uses high-density expendable bathythermograph (XBT) data from the AX08 transect and satellite altimetry to estimate properties of the currents and their variability in seasonal and interannual timescales. The method is well suited for surface currents, such as the North Equatorial Countercurrent (NECC), and has some skill in resolving the variability of North Equatorial Undercurrent (NEUC). The synthetic method fails to describe the variability of the South Equatorial Undercurrent (SEUC), which located in a region of small sea surface height variability. Our results confirm that the NECC shows a strong annual cycle of transport, with high values from July-December, and shows a possible strengthening (1-2 Sv on average) in the 2000 decade in comparison with the 1990s. On interannual time-scales, there is a positive correlation between the NECC transport and an anomalous sea surface temperature north-south gradient and the strengthening of the southwesterly winds. The NEUC shows a stronger transport values up to 10 Sv from January-July, and the interannual variability of the NEUC transport agrees with previous mechanisms such that it is correlated with the sea surface temperature on the Gulf of Guinea and southeastern equatorial bands, and reduction of equatorial winds. The EUC shows strong annual and semi-annual component of its variability in the region. This study shows that for a long-term monitoring system of the region, both altimetry and XBT data are necessary a near-real-time inference of the dynamical and thermodynamical properties of the current system in the tropical Atlantic region.

## 71 **1. Introduction**

72

73 The tropical Atlantic current system is of great importance for both inter-hemispheric and west-  
74 to-east exchange of heat and nutrients, which also impacts the climate and weather in the  
75 surrounding continental areas [e.g., McCready et al., 2002; Goldenberg et al., 2001; Sutton and  
76 Hodson, 2005].

77 In the upper tropical Atlantic, the main conduit for the inter-hemispheric water exchange is the  
78 North Brazil Undercurrent/Current (NBUC/NBC), which is a western boundary current system.  
79 On its northward path, part of the NBC retroflects eastward and feeds into a system of zonal  
80 countercurrents [Schott et al., 1995; Bourles et al., 1999a]. Specifically, into the Equatorial  
81 Undercurrent (EUC) along the equator, and three off-equatorial currents, the North Equatorial  
82 Countercurrent (NECC), the North Equatorial Undercurrent (NEUC), and to a lesser extent the  
83 South Equatorial Undercurrent (SEUC).

84 The EUC has been largely studied for being unique in terms of its dynamics [e.g., Pedlosky,  
85 1987] and for being strongly related to the equatorial upwelling, and the Atlantic zonal mode  
86 [e.g., Goes and Wainer, 2003; Hormann and Brandt, 2007]. The seasonal cycle of the EUC has  
87 been analyzed in a model simulation [Arhan et al., 2006], in which two transport maxima of the  
88 EUC were found in the central part of the basin, one during boreal fall and another during  
89 summer. They suggested that two different dynamical regimes drive the EUC seasonal cycle: in  
90 summer and fall, the simulated EUC is mostly driven by equatorial zonal forcing, and supplied  
91 from the ocean interior; in winter and spring, it is driven by remote forcing through the rotational  
92 wind component, and supplied from the western boundary currents.

93 In regard to the off-equatorial currents, the NECC, for having a strong surface signature, has also  
94 been widely studied. This current is mainly located between 3-10°N [Garzoli and Katz, 1983;  
95 Richardson and Walsh, 1986], with its seasonal variability linked to the migration of the  
96 Intertropical Convergence Zone (ITCZ) [e.g., Fonseca et al., 2004]. The NECC is forced mostly  
97 by wind stress within the equatorial band, either by local wind stress curl (WSC) or wave  
98 mechanisms. In addition, the NECC plays an important role in the global meridional overturning  
99 circulation [Frantantoni et al., 2000], since part of its flow is carried through Ekman transport  
100 northward, contributing to the meridional heat transport.

101 The off-equatorial undercurrents (OEUCs), though, are still not widely understood. They  
102 contribute to the shallow subtropical overturning cells that link the subtropical subduction  
103 regions to the equatorial and coastal upwelling regions. The OEUCs provide cold subthermocline  
104 water for the off-equatorial eastern upwelling regimes along the African coast, more specifically  
105 for the Guinea and Angola Domes [Schott et al. 2004; Doi et al., 2007]. Recent advances in  
106 understanding the Atlantic OEUCs are presented, for example, in Huettl-Kabus and Boening  
107 [2008] and Fischer et al. [2008]. It is important to study the variability of the OEUC, for  
108 instance, because of their contribution to the zonal transport of nutrients such as O<sub>2</sub> through long  
109 distances [Brandt et al. 2008]. The OEUCs are located below the thermocline coincident with the  
110 equatorial thermostat, at about 3–6° of latitude. The SEUC is mostly fed by internal  
111 recirculations [Schott et al. 1995, 1998], and according to Reverdin et al. [1991], its seasonal  
112 cycle at 30°W is characterized by a maximum transport in the boreal fall and a minimum in  
113 boreal spring. Close to its origin, the SEUC flow consists of large standing meanders [Fischer et  
114 al., 2008]. The NEUC is weaker than the SEUC, and also more variable. For instance, Schott et

115 al. [2003] used an average of 13 shipboard sections at 35°W, and stated as uncertain the  
116 existence of the NEUC. Indeed, the SEUC and NEUC show a large intra-seasonal variability  
117 linked to tropical instability wave activity [Jochum et al., 2003], and instabilities are more  
118 prominent in the northern hemisphere [Athie and Marin, 2008; Foltz et al., 2004b]. In the central  
119 basin, the separation between the NECC and NEUC is not very clear, but the NEUC may present  
120 higher transport values in boreal spring. On the western side of the basin, some studies [e.g.,  
121 Bourles et al., 1999; Goes et al., 2005] report only one eastward current core within the  
122 thermostat (between 100–250 m) that is comprised of waters from both the northern and  
123 southern hemispheres. Therefore, the EUC and NEUC may originate from this single core, and  
124 split into two cores with different water mass properties further east.

125 Several mechanisms may be responsible for the variability of the OEUCs. The tropical instability  
126 waves (TIWs) may drive intraseasonal variability of the OEUCs and the variability of their  
127 meridional displacement [Rowe et al., 2002]. Diffusive processes are important to the balance of  
128 the momentum of the OEUCs [McPhaden, 1984; Johnson and Moore, 1997], and the TIWs play  
129 an important role on the maintenance of the OEUCs as they act as a necessary source of heat flux  
130 and momentum, therefore explaining the rise of the OEUCs across the isopycnals [Jochum and  
131 Malanotte-Rizzoli, 2004]. Remote forcings are also vital for the onset and maintenance of the  
132 OEUCs. From these forcings we can include the upwelling in the eastern side of the basin  
133 [McCreary et al., 2002; Doi et al., 2007], the global thermohaline circulation [McCreary et al.,  
134 2002; Furue et al., 2007], and the subtropical cells [Marin et al., 2003; Hua et al., 2003].

135 Due to the high variability in the region, it is imperative to have an observational system that can  
136 resolve mesoscale features as well as the temporal structure of the currents. This can be realized  
137 through current meters and cruise data at a relatively high cost. One potential candidate for a  
138 sustainable observational system that resolves reasonably well the spatial and temporal scales,  
139 and presents a large spatial coverage is satellite altimetry. However, there are substantial upper-  
140 ocean currents with very weak sea surface signatures and spatial variability that cannot be  
141 resolved from surface topography fields alone [Goni and Baringer, 2002]. This deficiency can be  
142 overcome by establishing a relationship between certain current bands and their characteristic sea  
143 surface height signature. This can be performed by using altimetry in combination with other  
144 observational platforms such as high-resolution hydrographic data. Therefore, a reliable  
145 monitoring system for the region requires complementary platforms in order to produce details  
146 of the spatial and temporal variability [Goni and Baringer, 2002].

147 The high-density (HD) XBT project has been active for over 20 years, and aims at sustainably  
148 measuring physical properties of the upper ocean with mesoscale resolution. Its high spatial data  
149 resolution and repeated sampling of the region enable assessing upper-ocean properties such as  
150 temperature variability and heat storage, and permit the characterization of the variability of the  
151 major geostrophic currents.

152 The goal of this study is to quantify the variability of the eastward currents in the equatorial  
153 Atlantic during 1993-2011, and determine properties such as transport, velocity and location of  
154 these currents. More specifically, we focus on the seasonal and interannual signatures of the  
155 NECC, EUC, and OEUCs manifested in observations. To accomplish this goal, we consider a  
156 monitoring system that comprises data from the XBT transect AX08 between the years 2000-  
157 2011, and uses altimetry to synthetically quantify the variability of the eastward currents in the  
158 equatorial Atlantic for the 1993-2011 period. The validity of the synthetic method and future  
159 prospects are also assessed in this study.

160 Henceforth, this paper is structured as follows: first we describe in detail the data to be used and  
161 introduce the methodology (Section 2); in Section 3 we validate the synthetic method and apply  
162 it to study the seasonal and interannual cycles of the eastward equatorial currents, and analyze  
163 the surface response and forcings linked to the variability of these currents. Discussion,  
164 conclusions and recommendations for future work are provided in the final section of the paper.

165

## 166 **2. Data and methods**

167

168 This work uses mainly two observational platforms to infer the variability of the Atlantic  
169 eastward equatorial current system: the AX08 HD XBT transect and satellite altimetry data.  
170 Specifications of each of these datasets as well as additional data used are given below.

171

### 172 a) Hydrographic data

173 The temperature data used in this study are from 39 realizations of the AX08HD XBT transect,  
174 which spans between Cape Town and New York City since 2000, crossing the equator at about  
175 23°W (Figure 1c). An average of four transects per year has been achieved since 2002, with  
176 about 200-300 XBTs deployed on each transect, and a spacing of about 25 km between  
177 consecutive deployments in the tropics.

178 Individual temperature profiles are interpolated linearly to a 2 m resolution in depth for the upper  
179 800 m. The data are quality controlled by excluding the profiles whose horizontal gradients lie  
180 outside the three standard deviation range of all profiles. Next, the sections are horizontally  
181 interpolated to a 25 km resolution (along track) in latitude, using an optimal interpolation  
182 procedure with a Gaussian correlation function with a decorrelation length scale of 200 km and a  
183 low observational error of 0.01. This procedure produces evenly spaced and smoother data.  
184 Salinity is inferred from temperature profiles using the climatological T-S relationships from the  
185 World Ocean Database (WOD01, Conkright et al. [2002]).

186 Salinity profiles are extrapolated to the surface using a slab-layer approximation in the mixed  
187 layer. This is a standard approximation to overcome the non-unique characteristic of the T-S  
188 relationship in the tropical surface waters [Goes et al, 2005; Schott et al., 1998]. Goni and  
189 Baringer, [2002] has shown that differences between in-situ and climatological salinity are of the  
190 order of 0.3-0.4 psu, and the resulting uncertainty is among the largest contributors to the  
191 dynamic height error, with differences as large as 5 cm. This error is on the order of the  
192 sensitivity of our interpolation procedure, as well as on the order of the errors in the altimetric  
193 data [Cheney et al., 1994] used in our synthetic methodology described below.

194 We further explore the effect of the salinity on the steric height variability in the region by  
195 calculating the thermosteric and halosteric contributions to the total surface steric variability in  
196 three different latitudes, 4.5°N, 4.5°S and the equator (Figure 4) along the XBT transect. This  
197 calculation is performed using the methodology described in Foffonof and Froese [1955], and  
198 Tabata et al. [1986], in which departures from the mean sea level are estimated in terms of one  
199 component by keeping the other component fixed at its annual mean. The sum of the  
200 thermosteric and halosteric components is equal to the total steric variability of the XBT data.  
201 We illustrate their variability by showing the surface dynamic height ( $DH_0$ ) referenced to 800 m.

202 In all three latitudes, the contribution due to halosteric variability to the variability of  $DH_0$  is  
203 small, generally of the order of 1-2 cm. The variability of  $DH_0$  closely follows the thermosteric  
204 variability in all latitudes. Salinity has a higher contribution to the total variability at the equator  
205 and at  $4.5^\circ\text{S}$ , where the total  $DH_0$  variability is small. At those latitudes, the halosteric  
206 component count for 20-30% of the total variability. At  $4.5^\circ\text{N}$ , where there is high thermosteric  
207 variability with amplitude of  $\sim 30$  cm, only  $\sim 10\%$  of the  $DH_0$  variability is due to the halosteric  
208 component.

209 The typical distribution of temperature and salinity along the AX08 section is exemplified by the  
210 months of January and July of 2010 (Figure 2). This region is characterized by a warm well-  
211 mixed layer in the top 100 m, followed by a sharp temperature gradient of the upper thermocline  
212 of approximately  $0.1^\circ\text{C}/\text{m}$ . North and south of the displayed domain are characterized by high  
213 salinity values around 100 m, within the upper thermocline waters. The high salinity waters are  
214 characteristic of the Subtropical Underwaters (SUW). SUWs are formed in the subtropics and  
215 advected equatorward by the North and South Equatorial Currents (NEC and SEC, respectively).  
216 Underneath the SUW are the central waters, characterized by the nearly straight line in the T-S  
217 space. These waters form the thermostad between  $15^\circ\text{C}$  and  $12^\circ\text{C}$  found in the equatorial region  
218 [Reverdin et al., 1991], which is more pronounced in the southern hemisphere between  $5^\circ\text{S}$  and  
219 the equator at around 200 m depth (Figure 2), and are coincident with the position of the OEUCs.  
220 Therefore, the central waters are the main source of waters for the OEUCs, particularly the  
221 SEUC. The NEUC is known to have also a strong contribution from the high salinity SUW. This  
222 is an important feature of the region, and confirms the higher contribution of the western  
223 boundary to the NEUC, and higher contribution from eastern waters to the SEUC. In fact,  
224 previous works [e.g., Schott et al. 1995, 1998] suggest that the source waters for the SEUC  
225 originate from a recirculation gyre.

226

## 227 b) Remote sensed data

228 Satellite altimetry provides global coverage of the sea surface height anomaly (SSA) relative to  
229 the 1993-1999 mean period. Here we use the AVISO delayed mode product  
230 (<http://www.aviso.oceanobs.com>), which consists of gridded SSA relative to the 1992-1998  
231 period mean, obtained from a multi-satellite mission [Le Traon et al., 1998]. The delayed mode  
232 data are available continuously in a weekly temporal resolution since October 1992, on a  $1/3^\circ$   
233 horizontal grid, with precision of 2 cm [Ducet et al, 2000; Cheney et al., 1994]. For the purpose  
234 of this study, we use data from October 1992-December 2010, subtracted the 2000-2011 period  
235 mean from the SSA field, and interpolate the data onto a Mercator grid with  $0.25^\circ \times 0.25^\circ$   
236 resolution. Time and spatial linear interpolation is further applied when comparing to the XBT  
237 track locations.

238 We perform additional analysis using gridded sea surface temperature (SST) and pseudo-wind  
239 stress anomalies, averaged to monthly means and interpolated onto a  $1^\circ \times 1^\circ$  mercator grid. The  
240 SST data used is the optimum interpolation (OISST-v2) analysis produced weekly on a one-  
241 degree grid [Reynolds et al., 2002]. The reanalysis uses in situ and satellite SST plus SST  
242 simulated by sea ice cover. Before the reanalysis is computed, the satellite data is adjusted for  
243 biases using the method of Reynolds [1988] and Reynolds and Marsico [1993], which improves  
244 the large scale accuracy of the OI.

245 The pseudo-wind stress data is taken from cross-calibrated, multi-platform (CCMP), multi-  
246 instrument ocean surface wind velocity data set, available since July 1, 1987. This data set  
247 combines data derived from SSM/I, AMSRE, TRMM TMI, Quikscat and other missions using a  
248 variational analysis method (VAM) to produce a consistent record of ocean surface vector winds  
249 at a 25 km resolution [Atlas et al. 1996].

250

### 251 *Velocity calculation*

252 In order to study the variability of the eastward equatorial currents, we restrict the observations  
253 to a region with similar dynamical characteristics. We define a criterion that selects the sections  
254 whose mean longitude lie within the 68 percentile around their median value between 10°S-10°N  
255 (Figure 1a). The median longitude of the transect, which is about 40° oblique with respect to a  
256 true meridional section, is approximately 23°W. Our selection reduces the number of transects  
257 from 39 to 32, but assures that we are dealing with comparable data.

258 Geostrophic currents are inferred from the hydrographic data using the thermal wind relationship  
259 (cf. Figure 3), which is the standard dynamical method to calculate velocities from this type of  
260 data. The cross-sectional velocity field is calculated from the horizontal gradients of dynamic  
261 height. We estimate absolute velocities by using the XBT dynamic heights relative to 800 m  
262 ( $DH_{XBT}$ ), which the maximum depth reached by the XBTs, plus information about the absolute  
263 dynamic height at 800 m ( $DH_{IPRC}$ ):

$$264 \quad DH(z) = DH_{XBT}(z) + DH_{IPRC}(800), \quad (1)$$

265 where  $DH_{IPRC}(800)$  is taken from the monthly of International Pacific Research Center (IPRC)  
266 climatology (<http://apdrc.soest.hawaii.edu>). This dataset is available on a 1°x1° horizontal grid at  
267 27 depth levels derived from ARGO floats and altimetry observations, with mean sea level  
268 referred to MDT\_CNES-CLS09 [Rio et al., 2011]. The inclusion of  $DH_{IPRC}(800)$  reduces  
269 considerably the mean DH (~ 70 cm), but the DH gradients are not greatly affected (< 1 cm). In  
270 the equatorial region, geostrophy has an inflection point, and it is necessary to use the equatorial  
271 beta approximation, which relies on the calculation of higher order derivatives of the momentum  
272 equation near the equator. In this study, we apply the method of Vianna and Menezes [2003] for  
273 velocity calculations within  $\pm 3^\circ$  off the equator.

274 The a priori errors involved in the our methodology of geostrophic velocity estimation, which  
275 uses horizontal gradients of DH, are due to temperature precision of  $\pm 0.1^\circ\text{C}$  in the XBT  
276 measurements (O[1 cm/s]); salinity estimates from the T-S relationship (O[5 cm/s]); uncertainty  
277 in the level of known motion (O[1 cm/s]); and the angle of the transect with the true meridional  
278 (O[4 cm/s]) or smaller than 10% error in the region (Reverdin et al., 1991), assuming a priori  
279 small zonal pressure gradients. These errors add up to  $\pm 11$  cm/s, which is the order or smaller  
280 than the intrinsic mesoscale variability of the region. Therefore the standard monthly error will  
281 be used as a confidence interval for our calculations. At the equator, the beta plane  
282 approximation generates errors one order of magnitude greater (O[10 cm/s]) than the other  
283 components, and this caveat has to be taken into consideration in our analysis.

284

### 285 *Synthetic method*

286 Altimetry can be used along with other observational platforms to provide a broader four  
287 dimensional (spatial and temporal) coverage of the tropical Atlantic. Several works have used  
288 altimetric and hydrographic data to infer properties of the upper ocean, such as velocity,  
289 temperature, and salinity [e.g., Gilson et al., 1998; Ridgway, 2002; Phillips and Rintoul, 2002;  
290 Carnes et al., 1990].

291 Here, we apply a synthetic method to produce a hindcast of the velocity and density fields along  
292 the mean AX08 transect position (red line in Figure 1c). Since we are interested in velocity  
293 sections along isopycnal layers, we use as predictands potential density ( $\sigma_\theta$ ) and dynamic (steric)  
294 height (DH) from the surface to 800 m. The synthetic method uses altimetric information of sea  
295 surface height anomalies to generate the predictors. For this, we first calculate the absolute DH  
296 (Equation 1) for the 32 selected XBT sections. Anomalies of  $\sigma_\theta$  and DH ( $\sigma'_\theta$  and  $DH'$ ,  
297 respectively) are calculated by subtracting one mean annual field, defined by the mean WOD01  
298 density field along the mean section, and the mean dynamic height of all sections. For the latter,  
299 we first calculate monthly averages to reduce potential sampling biases towards any specific  
300 month. These results are mostly insensitive to the mean fields, since they just subtract a constant  
301 from the temporal fields, but subtracting a mean is known to reduce the variance of residuals in a  
302 linear regression. Anomalies of surface DH ( $DH'_0$ ) are linearly regressed onto  $DH'$  and  $\sigma'_\theta$  as a  
303 function of location and depth. Finally, the  $DH'_0$  is linearly regressed onto SSA, thus forming the  
304 link between the satellite and hydrographic observations. The  $DH'_0$  and SSA are well correlated  
305 (Figure 5), with an R-squared coefficient  $R^2 = 0.72$ , which is equal to the squared correlation in a  
306 simple linear regression, and a negligible intercept. This strong and unbiased relationship  
307 between the two variables show that the SSA captures well the baroclinic structures in the  
308 region, in specially the first mode. The highest variance of SSA occurs between 5-10°N, in the  
309 same region where the highest SSH gradients are located, and which are closely related to the  
310 dynamics of the NECC. Less variability is seen south of the equator, which can be due to  
311 compensating effects in the water column [Mayer et al., 2001]. The differences found here  
312 between SSA and  $DH'_0$  can arise from a number of factors, such as temporal and spatial  
313 sampling, tides not completely removed, as well as variations in the barotropic flow and changes  
314 in the entire depth of the ocean [Rintoul et al., 2002].

315 One indicator of the skill of this method to monitor the variability of the predictands  $DH'$  and  $\sigma'_\theta$   
316 is the temporal correlation between  $DH'_0$  and the predictands at each depth and latitude. The  
317 correlations (Figure 6) exhibit, as expected, opposite signs for  $DH'$  and  $\sigma'_\theta$ , since dynamic height  
318 is calculated from steric volume anomalies, which are inversely related to density [Pond and  
319 Pickard, 1983]. Apart from their sign, both fields show similar relationships, with highest  
320 correlations in the upper 200 m of the water column, and decreasing magnitude with depth. In  
321 certain latitudes, such as around 5°N, high correlations ( $R > 0.8$ ) can reach depths of 800 m,  
322 similarly to features described in Roemmich and Gilson [2001].

323

### 324 **3. Results**

325

326 The central tropical Atlantic has a very characteristic seasonal variability of the surface dynamic  
327 height. December through May are characterized by an almost flat  $DH_0$  between 10°S-10°N,  
328 implying small surface geostrophic velocities (Figure 3a,b). From June through November  
329 (Figure 3c,d) there is a stronger gradient of  $DH_0$  north of 5°N, and the NECC core is well



330 defined. A double NECC core is developed seasonally around 8°N, especially during the boreal  
331 fall, which is merged back into one single and stronger core generally during late summer/early  
332 fall. There is often a connection between the surface NECC with the thermocline layer, and the  
333 NEUC can be found as a lobe attached to the NECC. Two cores are also found in the  
334 thermocline layer in the north of the section, and the northernmost one is referred as the northern  
335 NECC core (nNECC; e.g., Didden and Schott [1992], Urbano et al. [2006]). During the spring,  
336 however, the NEUC is clearly detached from the nNECC. The isopycnal signature of the NEUC  
337 on the  $\sigma_\theta = 26.5$  isopycnal is seen at about 5°N, but does not show a very sharp gradient [c.f.,  
338 Bourles et al, 2002; Schott et al., 2003]. In contrast, the isopycnal signature in the south around  
339 4-5°S, where the SEUC is located around 200 m, has a very distinguished southward elevation of  
340 the  $\sigma_\theta = 26.5$  isopycnal. This is an indication that the meridional pressure gradient is a major  
341 driver of the SEUC variability. The isopycnal signature of the SEUC and its well defined core  
342 are visible throughout the year; therefore the SEUC is a permanent feature in the tropical  
343 Atlantic. The EUC core is located mostly south of the equator, and its core is generally located  
344 between the surface and 100 meters.

345 We define the location of the equatorial eastward currents by assigning a latitudinal band to each  
346 of the currents (Figure 3), similarly to Hüttl-Kabus and Böning [2008]: the NECC is defined  
347 between 3°-10°N; the NEUC between 3°-6°N; the SEUC between 3°-6°S; and the EUC between  
348 2.5°S and 2.5°N. We further characterize these currents by selecting isopycnal layers to define  
349 their vertical boundaries: one upper or surface layer, from the surface to  $\sigma_\theta = 24.5 \text{ kg m}^{-3}$ , and one  
350 lower or thermocline layer, from  $\sigma_\theta = 24.5 - 26.8 \text{ kg m}^{-3}$  [Schott et al. 1998]. The upper layer  
351 contains the NECC and part of the EUC, and the lower layer comprises OEUCs and deeper  
352 contribution of the EUC.

353

### 354 *Synthetic method validation*

355 We validate the synthetic method by comparing the transports of the eastward equatorial currents  
356 estimated from the 32 XBT transect realizations with their synthetic estimates along the same  
357 sections (Figure 7). The agreement between the transport estimated synthetically and from  
358 hydrography is high for the NECC (Figure 7a, b). The transport differences are in general less  
359 than 1 Sv, and in most of the cases the values are coincident with each other. A strong linear  
360 relationship is found between the two estimates, with the corresponding  $R^2 = 0.95$  for the linear  
361 fit. The EUC (Figure 7g, h), which has transport contributions from the upper and lower layers,  
362 and also accounts for artifacts of the equatorial beta approximation, shows a moderate-to-high  
363 agreement between the two estimates ( $R^2 = 0.55$ ).

364 Regarding the OEUCs, the synthetic method has still some skill for the NEUC ( $R^2 = 0.33$ ), but  
365 the transports using XBT data are generally higher than the synthetic ones (Figure 7c, d). The  
366 method fails for the SEUC, and the variance of the synthetically-derived SEUC transports barely  
367 resemble the ones from hydrography (Figure 7e). However, a mean transport of around 7 Sv is  
368 found on both methods. Due to the low agreement between the two SEUC transport estimates,  
369 the analysis of the synthetic method will only be applied to the EUC, NEUC and NECC.

370

### 371 3.1 Seasonal variability

372

373 Here we analyze the seasonal variability of the NECC, SEUC, NEUC and EUC along the AX08  
374 transect in terms of across transect transport, location and velocity at the core, using for this  
375 monthly averages of these variables. The selected 32 sections provide high resolution coverage  
376 along the central Atlantic of all months except February (Figure 1b). We apply the method of  
377 Foltz et al. [2004a] to the XBT-derived quantities to estimate the monthly mean and variance of  
378 transport, core velocity and position of the currents. This method fits the first two harmonics  
379 (annual and semi-annual) to the data in a least square fashion. It accounts for observational  
380 uncertainties by calculating a diagonal covariance matrix, with each diagonal term corresponding  
381 to a month, calculated as the sum of the fitting residual error and the observational error given by  
382 the standard deviation of each month. Fitting the first two annual harmonics to the data filter a  
383 large part of the mesoscale features. Below we analyze the each of the currents separately.

384

#### 385 a) NECC

386 The hydrographic data estimates show a strong annual cycle of the NECC, which alone  
387 represents 61% of the total transport variance, 34% of the maximum velocity and 50% of  
388 position variances (Table 1). Earlier analyses from drifter data also show a dominant contribution  
389 as high as 80% of the annual cycle for off-equatorial surface currents [e.g., Richardson and  
390 Walsh, 1986]. The NECC reaches its southernmost position in April-May ( $\sim 3-4^\circ\text{N}$ ), concomitant  
391 with its lowest transport ( $\sim 1$  Sv), and its northernmost position ( $\sim 7-9^\circ\text{N}$ ) in September-October  
392 (Figure 8i), at the period of its highest transport ( $\sim 10-11$  Sv) reached during summer/fall (Figure  
393 8a). This seasonal variability of the NECC agrees with previous observational findings  
394 [Richardson et al., 1992], and is known to be linked to the north-south migration of the ITCZ  
395 [Katz and Garzoli, 1982; Katz, 1987; Garzoli and Richardson, 1989; Garzoli, 1992], which is  
396 near the equator during boreal spring and farthest north during boreal summer. The NECC core  
397 velocity shows a similar pattern, with lowest values ( $\sim 10-20$  cm/s) in boreal spring and  
398 maximum values ( $\sim 55$  cm/s) in boreal fall (Figure 8e).

399 The synthetic estimates for the comparable 2000-2010 period, given by the blue line in Figure  
400 (8a,e,i), follow closely the ones from hydrography. As the synthetic estimates are averaged using  
401 a much longer time series, and the hydrographic data is subject to higher influence of mesoscale  
402 effects due to sparse temporal sampling (Figure 1b), differences are expected. For instance, the  
403 synthetic method shows generally higher explained variance for the annual cycle, from 78-92%  
404 of the transport variability. Regarding the maximum velocity time series (Figure 8e), the  
405 synthetic estimates show strong semi-annual variability (blue and green curves) that is not  
406 obvious in the hydrographic data (red curve). Our results suggest that, when comparing the  
407 synthetic estimates from the 2000-2010 decade to the ones from the previous decade (1993-  
408 2000), there is an increase in the transport of the NECC of about 1-2 Sv (Figure 8a). This change  
409 is however in the range of difference between the XBT and synthetic estimates.

410 The characteristics of the NECC transport over time can be retrieved by applying a wavelet  
411 transformation onto the NECC transport timeseries for the whole altimetric period (Figure 9a).  
412 The wavelet analysis [Torrence and Compo, 1998] confirms that the annual cycle of the NECC is  
413 the strongest signal. Furthermore, the energy of the annual period seems to be enhanced since  
414 2002, in agreement with the transport increase in the climatological time series (Figure 8a). The  
415 mechanisms related to interannual variability of this current, as well as for the NEUC, are  
416 explored in the section 3.2.

417

418 b) SEUC

419 The XBT-derived transport of the SEUC (Figure 8c) is weaker (~6 Sv) from July to September,  
420 in agreement with Reverdin et al. [1991], and also reduced in January-February (~5 Sv).  
421 Maximum transport values are found from October through December, reaching over 10 Sv, and  
422 the mean SEUC transport is about 7-8 Sv, in close agreement with Brandt et al. [2006] and  
423 Fisher et al. [2008]. The semi-annual harmonic is dominant, and explained variances are 90% for  
424 transport and 90% for maximum velocity, in agreement with the model work of Hüttl-Kabus and  
425 Böning [2008]. The core velocity is about 28 cm/s in May/June and October, and 20 cm/s in  
426 August (Figure 8g). Fisher et al. [2008] estimates a weaker mean SEUC velocity of 17.9-13.4  
427 cm/s in the region comprehending 23°W, but departures from this mean were found to add up to  
428 50 cm/s in its core. The mean position of the SEUC is about  $4.5 \pm 0.5^\circ\text{S}$ , and its position  
429 variability follows closely its transport variability, in that higher transports are associated to a  
430 more northern position and vice versa.

431

432 c) NEUC

433 The NEUC shows higher transport values during boreal spring (Figure 8b), when its transport  
434 ranges from 5-10 Sv, and low transports (~3 Sv) are observed in boreal fall. In comparison to the  
435 NECC transport seasonal cycle (Figure 8a), the NEUC bears an almost inverse relationship with  
436 the NECC, and therefore with the location of the ITCZ. The annual cycle dominates the XBT  
437 transport variability (54% of variance), whereas the semi-annual cycle only explains 1% of  
438 variance. The NEUC location (Figure 8i) cycle resembles the transport one, with more  
439 southward position (~3.5-4°N) in the boreal spring at the time of higher transport. The NEUC  
440 core velocity (Figure 8f) is mostly semi-annual (49% of variance), with core speed of about 25  
441 cm/s in boreal winter/summer, and smaller core velocity of about 10-15 cm/s during September-  
442 October, revealing that the NEUC is still present during these months. Synthetic estimates  
443 broadly agree in magnitude with the XBT ones, and may show some changes between the 1990's  
444 and 2000's decades. The wavelet analysis of the whole altimetric period indicates a possible shift  
445 in energy from a prevailing semi-annual to an annual variability starting in 2000 (Figure 9b).  
446 This energy peak is related to mesoscale features and wave mechanisms that are integral part of  
447 the NEUC dynamics. There is a hint of interannual variability in the NEUC transport, with  
448 increased energy centered in 1995, 2000 and 2006. The resemblance between the variability of  
449 the zonal winds in the equatorial region (Figure 9d) and the NEUC transport variability is  
450 notorious. Possible dynamical mechanisms of the NEUC variability will be explored in the next  
451 section.

452

453 d) EUC

454 Both methods agree well in that the EUC at 23°W is mostly located south of the equator,  
455 between  $-1 \pm 1^\circ\text{S}$  (Figure 8l). Despite the fact that the seasonal variability of the position of the  
456 EUC is small, there is indication of a more southward position is during late boreal fall/winter.  
457 The core velocity from XBT agrees reasonably well with the synthetic method, but higher range,  
458 from 20-45 cm/s for XBT estimate, and 20-30cm/s for the synthetic estimate. These values are  
459 somewhat smaller than direct velocity observations of the EUC, which show higher values (~80-

460 90 cm/s) in the EUC core [Giarolla et al., 2005; Brandt et al., 2006, 2008]. The EUC transport  
461 hinges to a semi-annual variability from both hydrographic and synthetic estimates (Figure 8d),  
462 with lower values during the summer and winter, and higher values during spring and late  
463 summer-early fall. The wavelet analysis (Figure 9c) confirm the observed findings of a strong  
464 semi-annual transport variability of the EUC, in agreement with results from previous model  
465 studies for the central equatorial Atlantic [Arhan et al., 2006; Hormann and Brandt, 2007;  
466 Kolodziejczyk et al., 2009]. The high variability of the XBT estimates in some months, see for  
467 instance the errorbars of transport and core velocity during July (Figures 8d,h), are artifact of the  
468 performance of the equatorial method for specific density structure of the month, and need to be  
469 interpreted carefully.

470

### 471 3.2 Interannual variability

472 In this section, we investigate statistically the interannual signature of the NECC and NEUC on  
473 the sea surface temperature (SST) and surface wind stress. We restrict our analysis to the NECC  
474 and NEUC transports only because the synthetic method does not provide good estimates of the  
475 SEUC variability, and we do not expect that the equatorial beta approximation can produce  
476 reliable estimates of the EUC transport variability on interannual timescales. We use in this  
477 analysis monthly transport anomalies with respect to their monthly climatology, with the intent  
478 of reducing the influence of eddy variability on these currents. Furthermore, we perform a  
479 correlation analyses of the monthly transport anomalies of the NECC and NEUC with both the  
480 gridded monthly anomalies of SST [Reynolds et al., 2002] and pseudo-wind stress (CCMP; Atlas  
481 et al. [1996]) datasets. Correlations are performed for the 1993-2011 period over the tropical  
482 Atlantic, taking only statistically significant values into account ( $p < 0.05$ ). In this analysis, all  
483 data are standardized by subtracting their mean and dividing by their standard deviation, and  
484 smoothed with a 5-month moving average.

485

#### 486 a) NECC

487 The correlation of the NECC transport anomalies with SST anomalies (SSTA) for the altimetric  
488 period produces a very distinctive pattern in the form of an anomalous interhemispheric SST  
489 gradient (Figure 10), with positive phase north, just off equator in the ITCZ region, centered at  
490  $\sim 30^\circ\text{W}/2^\circ\text{N}$  and extending northeastward, and negative phase in the central south of the domain,  
491 with center at  $\sim 20^\circ\text{S}/17^\circ\text{W}$ . The corresponding correlation with pseudo-wind stress anomalies  
492 indicates an anomalous strengthening of the southeasterly trade winds, with largest magnitude in  
493 the western equatorial region, where winds are also strongest. This pattern is reminiscent of the  
494 meridional mode and agrees with an underlying mechanism explained by the wind-evaporation-  
495 SST (WES) feedback, which involves interactions between SST changes and wind-induced  
496 latent heat flux [Xie and Philander, 1994; Chang et al., 1997]. Similar pattern has been also  
497 found for the NECC in a recent study using CEOF analysis surface drifter data [Hormann et al.,  
498 2012].

499 The SSTA gradient index can be defined by subtracting area averages of positive ( $35^\circ\text{W}$ -  
500  $17^\circ\text{W}/0^\circ$ - $7^\circ\text{N}$ ) and the negative ( $30^\circ\text{W}$ - $15^\circ\text{W}/11^\circ\text{S}$ - $23^\circ\text{S}$ ) correlation regions. The correlation  
501 between the monthly NECC transport anomalies and the SST index is  $R=0.45$ , with maximum  
502 correlation at zero lag. The meridional wind (y-wind) index that is centered in the northern box

503 above described, which focus on the variability of the meridional wind stress near the central  
504 equatorial region, shows a positive correlation of also  $R=0.45$ , with the NECC leading the wind  
505 strengthening within one month lag. This relationship indicates that there is a fast response of the  
506 NECC to temperature and wind anomalies that might be explained by either the fast adjustment  
507 time of the ocean through equatorial waves [e.g., Ma, 1996] or with anomalies connected with  
508 the simply strengthening of the surface retroflexion of the NBC.

509

#### 510 b) NEUC

511 For the NEUC transport anomalies, the pattern arising from its correlation with SST and wind  
512 stress anomalies (Figure 11) can be described as follows: while negative SST coefficients prevail  
513 in the northeastern part of the basin, positive correlations are found along the equator and on the  
514 southwestern African coast, in addition to reduced trades in the western to central equatorial  
515 Atlantic. These two regions of high interannual SST variability are sites where the thermocline is  
516 rather shallow, in which anomalous SSTs are highly correlated with thermocline depths [Carton  
517 et al., 1996], and therefore are closely linked to the equatorial mode [Zebiak, 1993]. Some  
518 studies [c.f. Chiang et al., 2001; Foltz et al., 2011] define the difference of anomalous SST  
519 averaged within these two regions as the Atlantic Meridional Mode (AMM). This pattern is tied  
520 to the seasonal cycle [Servain et al., 1999], in which anomalous negative/positive SST gradient  
521 in boreal spring is associated with westerly/easterly equatorial wind anomalies. Therefore, winds  
522 at the equatorial region respond to this negative pattern of the interhemispheric SST gradient by  
523 reducing their strength, therefore shifting the ITCZ south [Moura and Shukla, 1996]. This could  
524 reinforce the positive anomaly in the equatorial region by decreasing upwelling there [Foltz et  
525 al., 2011]. Similar pattern was also found in model studies for the upwelling regions of virtual  
526 lagrangian floats released inside the NEUC [Stramma et al., 2005; Huettl-Kabus and Boening,  
527 2008]. Therefore a relationship between the overlying atmospheric variability and the NEUC  
528 dynamics can be drawn, taking into account the current knowledge of the variability of the  
529 region. According to the model study of Huettl-Kabus and Boening [2008], part of the NEUC  
530 can merge on the EUC, contributing for the equatorial upwelling. Since the EUC and NEUC are  
531 known to pertain to only one core at their origin in the western Atlantic, their synergy with  
532 respect to the equatorial upwelling may be expected. In addition, Stramma et al. [2005] shows a  
533 direct trajectory from virtual lagrangian particles from the NEUC towards the Guinea Dome.  
534 Some mechanisms have been proposed to describe the link between these two regions of high  
535 SST variability [e.g., Doi et al., 2009], in which both coastal upwelling in the Guinea Dome and  
536 WES feedback related to the Atlantic meridional mode are involved. In the Guinea Dome region,  
537 where the NEUC transport is negatively correlated with SSTA, prevailing wind anomalies are  
538 directed to the southwest, proper for the increased coastal upwelling in the region. This link  
539 between the NEUC and the Guinea Dome has long been proposed [e.g., Voituriez, 1981], since  
540 the uplifting of the thermal structure in the dome extends much further down the thermocline.  
541 The strong upwelling in the region is related, for instance, to the outcropping of the  $\sigma_{\theta} = 24.5$   
542 layer. The dynamics underlying this relationship is that a cooler Guinea Dome and warmer  
543 equatorial region would increase the north-south density gradient in the NEUC region, and  
544 strengthen its core. Some studies indeed suggest that the upwelling in the Guinea Dome region  
545 can drive the NEUC variability [MCready et al., 2002]. Likewise we have produced for the  
546 NECC analysis, we create one interhemispheric SST index, now taking into account the  
547 difference between the SSTA in the Guinea Dome region ( $15^{\circ}$ - $25^{\circ}$ N/ $15^{\circ}$ - $30^{\circ}$ W) and in the

548 southeast Atlantic (0-8°S/15°W-10°E), along with one meridional wind stress index inside the  
549 Guinea Dome box. The correlation between the NEUC transport and the SST index is  $R = -0.50$   
550 at zero lag, and  $R = -0.31$  with the y-wind index in the Guinea Dome. This significant link  
551 between the variability of the AMM and the equatorial mode with the NEUC might explain the  
552 interannual pattern observed in the wavelet analysis (Figure 9b), such that the anomalously  
553 strong SST gradient during 1and increased Guinea Dome upwelling strengthened the NEUC  
554 during those periods.

555

#### 556 **4. Discussion and Conclusions**

557

558 In the present study, we use a combination of HD XBT transect data along with altimetry to  
559 investigate the variability of the eastward surface and subsurface currents in the equatorial  
560 Atlantic. The AX08 XBT transect is carried out on average four times a year since 2000. The  
561 high spatial resolution of the hydrographic data and its repeated sampling rate of the region  
562 enable us to assess the dynamical and thermodynamical properties of the upper ocean, and  
563 permit the characterization of the seasonal cycle of major geostrophic zonal currents, such as the  
564 NECC and NEUC. However, due to strong intra-seasonal variability generated by the eddy  
565 variability and the passage of TIWs in the region [Weisberg and Weingartner, 1988; Jochum et  
566 al., 2004], it is necessary to carry out a high number of sections for one to be able to produce a  
567 good statistical estimate of the seasonal cycle of these currents using hydrography only. For  
568 instance, from the 32 XBT sections analyzed, observations during the months of January, May  
569 and June were only carried once in each month, and there is yet no section available for  
570 February.

571 Merging altimetry and XBT data with this synthetic methodology reproduces the main features  
572 of the transect with depth that cannot be seen with altimetry alone, and overcome the sampling  
573 restriction and produce a smoother estimate of the long term average of the properties of the  
574 currents. The best performance of the synthetic profile methodology for the AX08 region resides  
575 in the upper 200-300 m of the water column, where the temporal correlations of SSA and  
576 anomalies of density and dynamic height are above 0.8.

577 The seasonal cycle of the eastward equatorial currents derived from our analyses are in good  
578 agreement with previous works [e.g., Peterson and Stramma, 1990; Bourles et al., 2002]. The  
579 NECC contains a strong seasonal cycle, with transport ranging from 1 to 11 Sv and maximum  
580 speed greater than 50 cm/s. Its seasonal variability is related to the migration of the ITCZ (not  
581 shown), with stronger transport positively correlated to the northward shift of the ITCZ. On  
582 interannual timescales, the NECC strength is linked to the strengthening of the southwesterly  
583 trade winds and an interhemispheric SST gradient pattern (positive north and negative south).  
584 The NECC strengthening during positive interhemispheric SST gradient might act as a positive  
585 feedback, since it would increase the eastward transport of warmer western waters in the region  
586 of the SST maximum gradient. This SST pattern is known to be influenced by the Atlantic  
587 Meridional Oscillation (AMO), and North Atlantic Oscillation (NAO), and connections from the  
588 eastern Pacific [Enfield and Mayer, 1996] There is indication of a recent strengthening of the  
589 climatological NECC of ~1-2 Sv, and attribution can be related to the recent decrease in sulphate  
590 aerosols emissions over the North Atlantic, and the associated increase in the SST there and  
591 consequent trend northward of the ITCZ [Chang et al., 2011].

592 The NEUC annual cycle is characterized by stronger transport values from January-July up to 10  
593 Sv, therefore in opposite phase to the NECC transport cycle. The NEUCs core is located between  
594 4-6°N with maximum velocities of about 30 cm/s in June-July. Synthetic estimates suggest that  
595 the semi-annual range of the spectrum used to be stronger during the 1990's, and that there was a  
596 shift to a more annual period, which shows to be intermittent, since 2000. This agrees with the  
597 pattern variability of the Atlantic meridional mode (AMM) described in Foltz et al. [2011]. We  
598 show that the variability of the NEUC transport is statistically related to the upwelling in the  
599 AMM and the variability Guinea Dome region [Schott et al. 2004].

600 NEUC and NECC bear an inverse relationship of transport in interannual timescales. Both are  
601 linked to the variability of the cross equatorial wind stress, but the NECC has positive values for  
602 increased southwesterlies whereas the NEUC is increased to negative southwesterlies anomalies.  
603 The interannual modes of tropical Atlantic variability are strongly tied to their seasonal cycle  
604 [Nobre and Shukla, 1996], therefore this also explains the inverse relationship between NEUC  
605 and NECC that holds for seasonal timescales.

606 The SEUC calculated from hydrography is located on average at 4.5°S, and presents a weaker  
607 seasonal cycle, with mean transport of about 8 Sv, higher values from October-December, and a  
608 secondary peak during May-June. Velocities during the stronger months range from 25-30 cm/s.  
609 The synthetic method so far represented well the mean flow but not its variability. This may be a  
610 shortfall of our methodology since the variability of SSA in the southern tropical Atlantic is  
611 somewhat weak (Figure 1c), and the surface signature of the current is masked by compensating  
612 effects in the water column (Figure 6). The compensating effects can be, for example, driven by  
613 buoyancy and wind forcing components of the same magnitude and with opposing signals  
614 [Mayer et al., 2001]. Therefore, a regular sampling from the HR XBT transect is particularly  
615 important for monitoring the SEUC.

616 The EUC can be resolved by the equatorial beta approximation using hydrographic data. The  
617 EUC seasonal cycle is well explained using the first and second harmonics, with peaks in boreal  
618 fall and boreal spring, therefore the semi-annual period is important for its variability in the  
619 region. The mean EUC transport is around 15 Sv, close to direct observations [e.g., Brandt et al.,  
620 2006], but the maximum velocity at the core ranges within 30-40 cm/s, about 50 % of direct  
621 mooring observations [e.g., Giarolla et al., 2005], which is a known artifact of the equatorial beta  
622 plane approximation [e.g., Langerlof et al., 2000]. A comparison with EUC transport estimates  
623 using mooring data in the equatorial region (Andreas Funk, personal communication) shows  
624 similar annual variability as the one derived from the method of Vianna and Menezes [2003].

625 Results from this study are subject to specific caveats that provide avenues for future studies and  
626 improvements to the methodology. First, we use a simple statistical method to infer relationship  
627 between the surface height and the ocean properties at depth. Using an improved statistical  
628 method may allow the inclusion of additional information, such as latitudinal cross-correlation  
629 between and autocorrelation of the residuals with depth, and use if additional constraints derived  
630 from co-located observations. Second, we use climatological values of salinity calculated from  
631 the mean T-S relationship at each location, and also climatological values of the absolute  
632 dynamic height at 800 m. Available observations from ARGO, for example, could reduce errors  
633 in the methodology. However, these observations are mostly restricted to the last 5-6 years.  
634 Third, along track data might provide better agreement between altimetric sea surface height and  
635 hydrographic-based dynamic height at the surface, since it is less smoothed by optimal  
636 interpolation procedures. Finally, high resolution modeling can fill the gaps of the observational

637 results, and confirm if our conclusions in a dynamically consistent fashion. These possibilities  
638 will be explored in a future study.

639

## 640 **Acknowledgements**

641 The authors would like to thank Greg Foltz and Rick Lumpkin for insightful discussions. This  
642 research was carried out under the auspices of the Cooperative Institute for Marine and  
643 Atmospheric Studies (CIMAS), a cooperative institute of the University of Miami and the  
644 National Oceanic and Atmospheric Administration, cooperative agreement #NA17RJ1226. Any  
645 errors and opinions are those of the authors alone.

646

## 647 **References**

648 Atlas, R., R. N. Hoffman, S. C. Bloom, J. C. Jusem, and J. Ardizzone (1996), A multiyear global  
649 surface wind velocity data set using SSM/I wind observations. *Bull. Amer. Meteor. Soc.*, 77,  
650 869–882.

651 Bourles, B., Y. Gouriou, and R. Chuchla (1999), On the circulation in the upper layer of the  
652 western equatorial Atlantic, *J. Geophys. Res.*, 104, C9, 21.151-21.170.

653 Bourles, B., M. D'Orgeville, G. Eldin, Y. Gouriou, R. Chuchla, Y. DuPenhoat, and S. Arnault  
654 (2002), On the evolution of the thermocline and subthermocline eastward currents in the  
655 Equatorial Atlantic, *Geophys. Res. Lett.*, 29(16), 1785, doi: 10.1029/2002GL015098.

656 Brandt, P., V. Hormann, B. Bourle`s, J. Fischer, F. A. Schott, L. Stramma, and M. Dengler  
657 (2008), Oxygen tongues and zonal currents in the equatorial Atlantic, *J. Geophys. Res.*, 113,  
658 C04012, doi:10.1029/2007JC004435.

659 Chang, C-Y, J.C.H. Chang, M.F. Wehner, A.R. Friedman and R. Ruedy (2011), Sulfate Aerosol  
660 Control of Tropical Atlantic Climate over the Twentieth Century, *J. Climate*, 24, 2540-2555,  
661 DOI: 10.1175/2010JCLI4065.1.

662 Cheney, R., L. Miller, R. Agreen, N. Doyle and J. Lillibridge (1994), TOPEX/Poseidon: The 2-  
663 cm solution, *J. Geophys. Res.*, 99, 24,555-24,563.

664 Conkright, M. E., et al. (2002), *World Ocean Database 2001*, vol. 1, Introduction, NOAA Atlas  
665 NESDIS 42, 159 pp., U.S. Gov. Print. Off., Washington D. C.

666 Didden, N., F. Schott (1992), Seasonal variations in the Western Tropical Atlantic: surface  
667 circulation from geosat altimetry and WOCE model results. *J. Geophys. Res. C* 97 (3), 3529–  
668 3541.

669 Ducet, N., P.-Y. Le Traon, and G. Reverdin (2000), Global high resolution mapping of ocean  
670 circulation from Topex/Poseidon and ERS-1 and -2, *J. Geophys. Res.*, 105 (C8), 19477-19498.

671 Fischer, J., V. Hormann, P. Brandt, F. A. Schott, B. Rabe, and A. Funk (2008), South Equatorial  
672 Undercurrent in the western to central tropical Atlantic, 35, L21601, doi:  
673 10.1029/2008GL035753.

674 Foltz, G., S. Grodsky, and J. Carton (2004a), Seasonal salt budget of the northern tropical  
675 Atlantic Ocean along 38°W, *J. Geophys. Res.*, 109, C03052, doi: 10.1029/2003JC002111.



676 Foltz, G., J. Carton, E.P. Chassignet (2004b), Tropical instability vortices in the Atlantic Ocean,  
677 *J. Geophys. Res.*, 109, C03029, doi: 10.1029/2003JC001942.

678 Fonseca, C., G. Goni, W. E. Johns, and E. J. D. Campos (2004), Investigations of the North  
679 Brazil Current retroflexion and North Equatorial Countercurrent variability, *Geophys. Res. Lett.*,  
680 31, L21304, doi: 10.1029/2004GL020054.

681 Frantantoni, D., W. E. Johns, T. L. Townsend, and H. E. Hurlburt (2000), Low latitude  
682 circulation and mass transport pathways in a model of the tropical Atlantic Ocean. *J. Phys.*  
683 *Oceanogr.*, 30, 1944–1966.

684 Goes, M., and I. Wainer (2003), Equatorial Currents transport changes for extreme warm and  
685 cold events in the Atlantic Ocean, *Geophys. Res. Lett.*, 30 (5), doi: 10.1029/2002GL015707.

686 Goes, M., R. Molinari, I. C. A. Silveira, and I. Wainer (2005), Retroflexions of the North Brazil  
687 Current during February 2002. *Deep Sea Research Part 1. Oceanographic Research Papers*, v. 52,  
688 n. 3, pp. 647-667, doi: 10.1016/j.dsr.2004.10.010.

689 Goldenberg, S. B., C. Landsea, A. M. Mestas-Nunez, and W. M. Gray (2001), The recent  
690 increase in Atlantic hurricane activity, *Science*, 293, 474– 479.

691 Hüttl-Kabus, S., and C. W. Böning (2008), Pathways and variability of the off-equatorial  
692 undercurrents in the Atlantic Ocean *Journal of Geophysical Research - Oceans*, 113. C10018.  
693 ISSN 0148-0227. DOI: 10.1029/2007JC004700.

694 Jochum, M., P. Malanotte-Rizzoli, and A. Busalacchi (2004a), Tropical instability waves in the  
695 Atlantic Ocean, *Ocean Modell.*, 7, 146–163.

696 Jochum, M., and P. Malanotte-Rizzoli (2004), A new theory for the generation of the equatorial  
697 subsurface undercurrents, *J. Phys. Oceanogr.*, 34, 755–771.

698 Lagerloef, G., G. T. Mitchum, R. B. Lukas, and P. Niiler (1999), Tropical Pacific near surface  
699 currents estimated from altimeter, wind and drifter data, *J. Geophys. Res.*, 104, 23,313-23,326.

700 Le Traon P.Y., F. Nadal, N. Ducet (1998), An improved mapping method of multisatellite  
701 altimeter data. *J. Atmos. Oceanic Technol.*, 15, 522-533.

702 Locarnini, R. A., A. V. Mishonov, J. I. Antonov, T. P. Boyer, and H. E. Garcia (2006), *World*  
703 *Ocean Atlas 2005, Volume 1: Temperature*. S. Levitus, Ed. NOAA Atlas NESDIS 61, U.S.  
704 Government Printing Office, Washington, D.C., 182 pp.

705 Mayer, D., R.L. Molinari, and M. O. Baringer (2001), Transition regions and their role in the  
706 relationship between sea surface height and subsurface temperature structure in the Atlantic  
707 Ocean, *Geophysical Research Letters*, 28, 20, 3943-3946.

708 McCreary, J.P., P. Lu, and Z. Yu (2002), Dynamics of the Pacific subsurface countercurrents. *J.*  
709 *Phys. Oceanogr.*, 32, 2379–2404.

710 Nobre, P., and J. Shukla (1996), Variations of sea surface temperature, wind stress and rainfall  
711 over the tropical Atlantic and South America, *J. Clim.*, 9, 2464– 2479.

712 Picaut, J., and R. Tournier (1991), Monitoring the 1979-1985 equatorial Pacific current  
713 transports with expendable bathythermograph data, *J. Geophys. Res.*, 96, 3263-3277.

714 Pond, S. and G. Pickard (1983), *Introductory Dynamical Oceanography*, 2<sup>nd</sup> Ed., Pergamon Press  
715 Inc., NY 10523, USA.

716 Reverdin, G., P. Rual, Y. du Penhoat, and Y. Gouriou (1991), Vertical structure of the seasonal  
717 cycle in the central equatorial Atlantic Ocean: XBT sections from 1980 to 1988, *J. Phys.*  
718 *Oceanogr.*, 21, 277-291.

719 Reynolds, R.W., N.A. Rayner, T.M. Smith, D.C. Stokes, and W. Wang (2002), An improved in  
720 situ and satellite SST analysis for climate. *J. Climate*, 15, 1609-1625.

721 Richardson, P.L., and D. Walsh (1986), Mapping climatological seasonal variations of surface  
722 currents in the tropical Atlantic using ship drifts, *Journal of Geophysical Research*, 91, 10537-  
723 10550.

724 Ridgway, K. R., R. C. Coleman, R. J. Bailey, and P. Sutton (2008), Decadal variability of East  
725 Australian Current transport inferred from repeated high-density XBT transects, a CTD survey  
726 and satellite altimetry, *J. Geophys. Res.*, 113, C08039, doi:10.1029/2007JC004664.

727 Rintoul, S. R., S. Sokolov, and J. Church (2002), A 6 year record of baroclinic transport  
728 variability of the Antarctic Circumpolar Current at 140°E derived from expendable  
729 bathythermograph and altimeter measurements, *J. Geophys. Res.*, 107(C10), 3155,  
730 doi:10.1029/2001JC000787.

731 Rio, M. H., S. Guinehut, and G. Larnicol (2011), New CNES-CLS09 global mean dynamic  
732 topography computed from the combination of GRACE data, altimetry, and in situ  
733 measurements, *J. Geophys. Res.*, 116, C07018, doi:10.1029/2010JC006505.

734 Roemmich, D., J. Gilson (2001), Eddy Transport of Heat and Thermocline Waters in the North  
735 Pacific: A Key to Interannual/Decadal Climate Variability? *J. Phys. Oceanogr.*, 31, 675–687.

736 Schott, F., L. Stramma and J. Fischer (1995), The warm water inflow into the western tropical  
737 Atlantic boundary regime, spring 1994, *Journal of Geophysical Research* 100, pp. 24745–24760.

738 Schott, F., J. Fisher, and L. Stramma (1998), Transports and pathways of the upper-layer  
739 circulation in the western tropical Atlantic, *Journal of Physical Oceanography*, 28, pp. 1904–  
740 1928.

741 Schott, F. A., M. Dengler, P. Brandt, K. Affler, J. Fischer, B. Bourles, Y. Gouriou, R. L.  
742 Molinari, and M. Rhein (2003), The zonal currents and transports at 35°W in the tropical  
743 Atlantic, *Geophys. Res. Lett.*, 30, 7, 1349, doi:10.1029/2002GL016849.

744 Sutton, R. T. and Hodson, D. L. R. (2005), Atlantic Ocean forcing of North American and  
745 European summer climate. *Science*, 309 (5731), 115-118, doi: 10.1175/JCLI4038.1.

746 Torrence, C and G. P. Compo (1998), *A Practical Guide to Wavelet Analysis.*, *Bulletin of the*  
747 *American Meteorological Society*, 79, 1, 61-78. Software available at URL:  
748 <http://atoc.colorado.edu/research/wavelets/>

749 Urbano, D. F., M. Jochum, and I. C. A. Silveira (2006), Rediscovering the second core of the  
750 Atlantic NECC, *Ocean Modell.*, 12, 1–15, doi:10.1016/j.ocemod.2005.04.003.

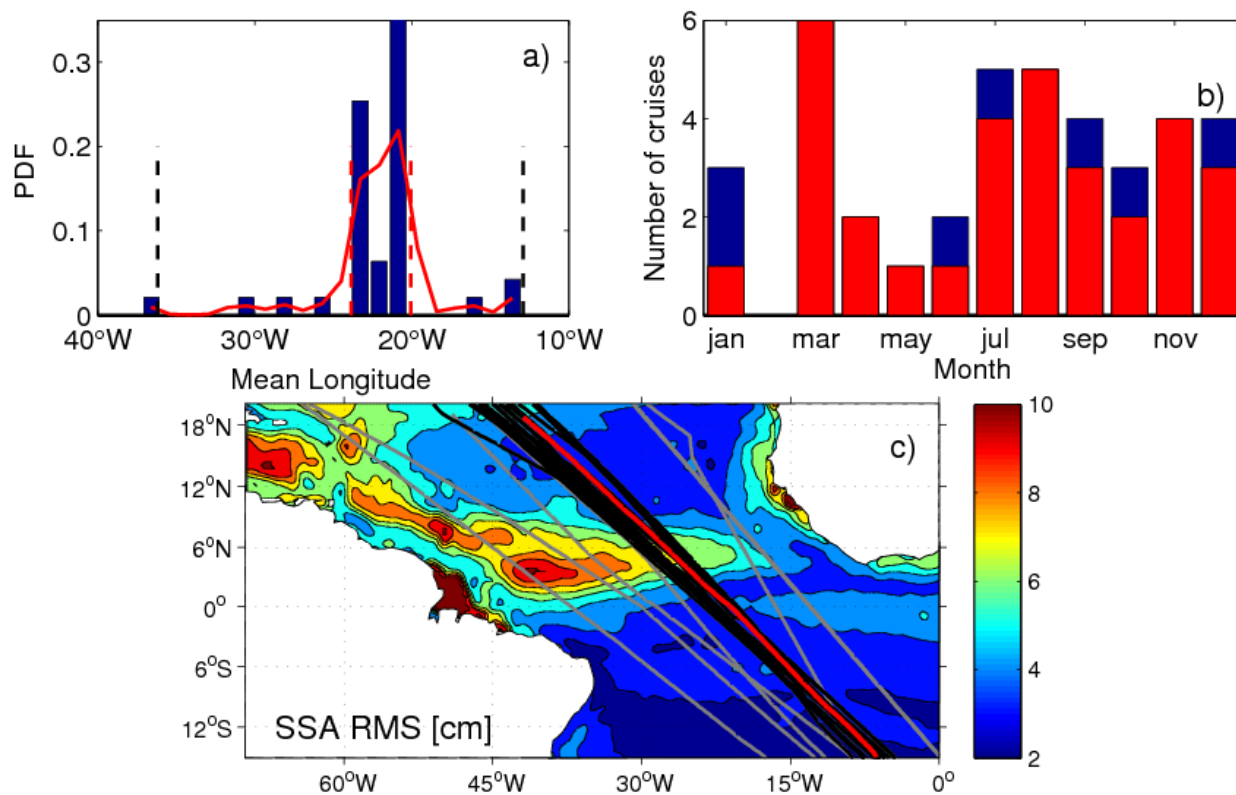
751 Vianna, M., and V. V. Menezes (2003), A seasonal and interannual study of the western  
752 equatorial Atlantic upper thermocline circulation variability, In: *Interhemispheric Water*  
753 *Exchange in the Atlantic Ocean*, G.Goni & P.Malanotte-Rizzoli, Editors, 2003: 137-174.

754 Voituriez, B. (1981), Les sous courants equatoriaux Nord et Sud et la Formation des domes  
755 thermiques tropicaux, *Oceanolog. Acta.*, 5, 497-506.

756 Weisberg, R., and T. Weingartner (1988), Instability waves in the equatorial Atlantic Ocean, *J.*  
757 *Phys. Oceanogr.*, 18, 1641– 1657.

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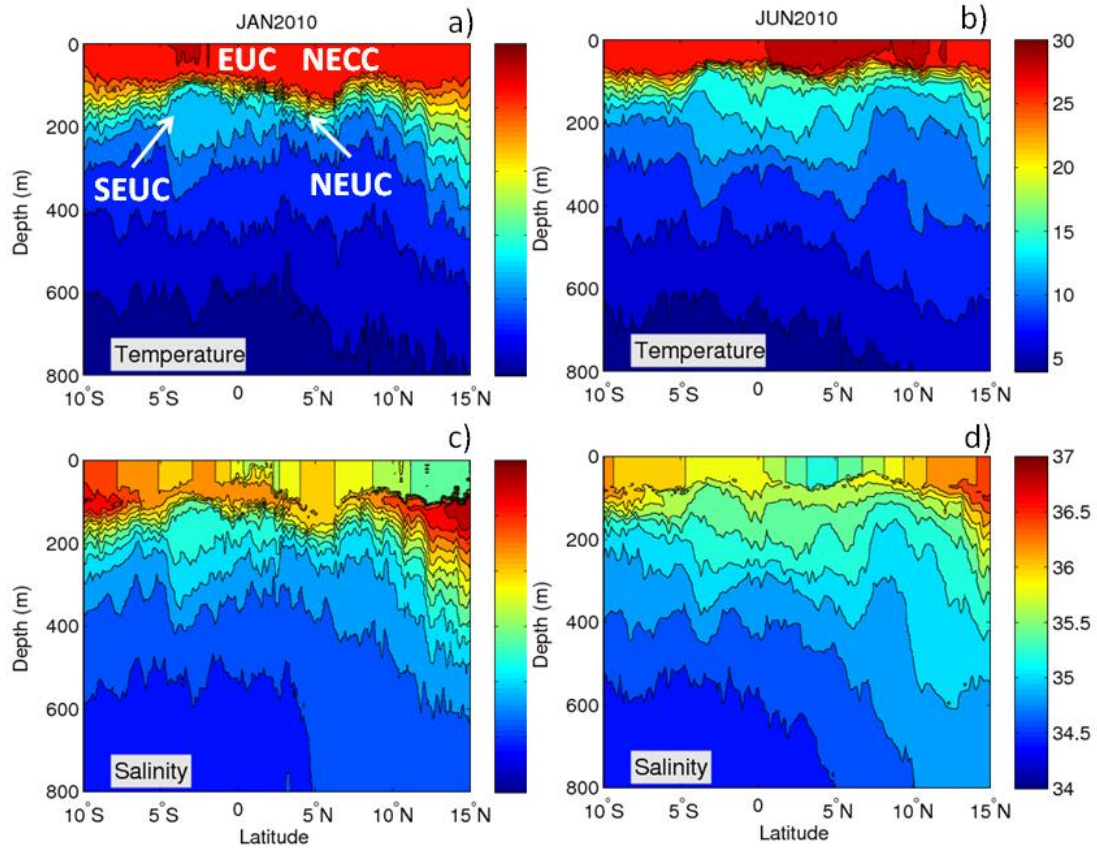
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769 Figure 1: a) Probability density function of the averaged longitude between 10°S-10°N for all the  
770 AX08 XBT sections. b) Histogram of the monthly distribution of the number of the sections  
771 before (blue) and after (red) the selection of the sections that fit into the 68 percentile of a). c)  
772 Root mean square of SSA (cm) in the tropical Atlantic (filled contours), with superimposed  
773 AX08 sections: selected sections (black), sections not included in the analysis (gray), and mean  
774 of the selected sections (red).

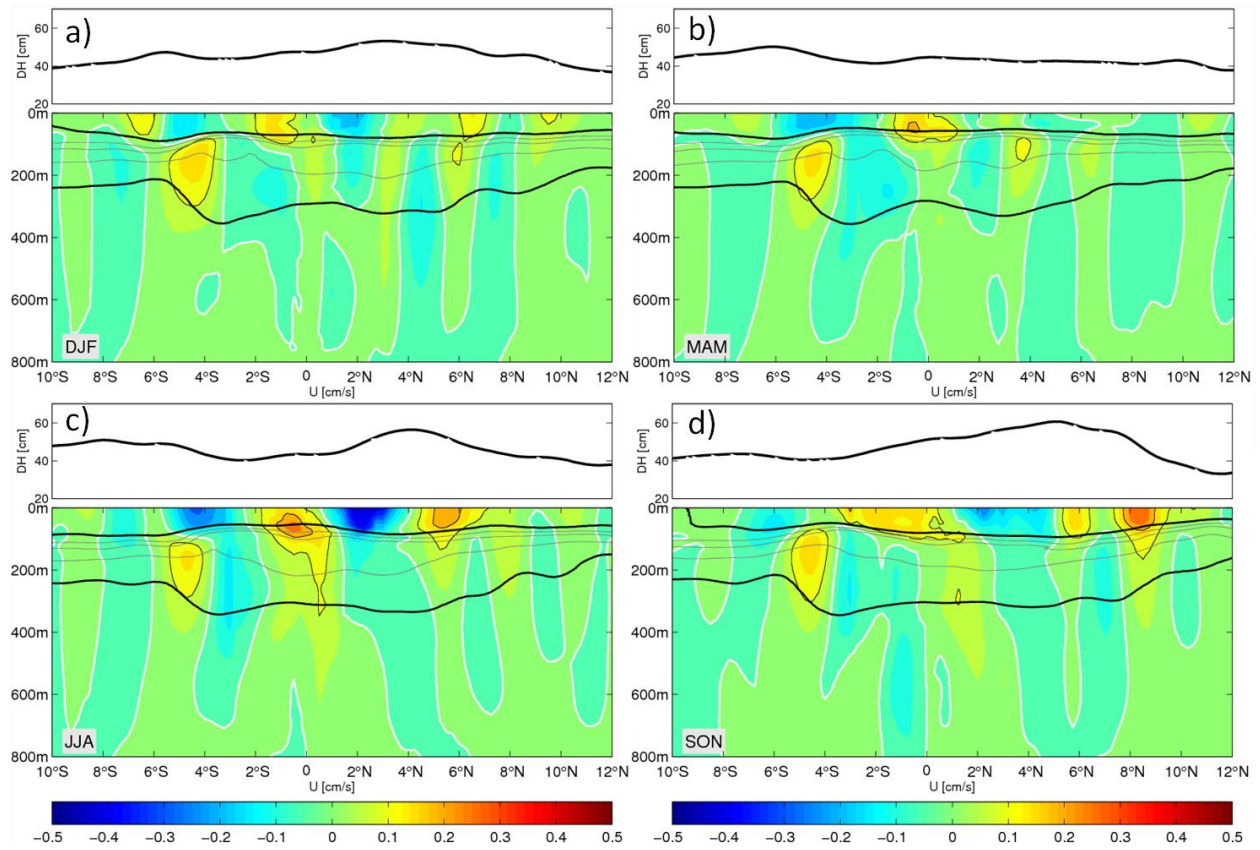
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Figure 2: Example of typical temperature ( $^{\circ}\text{C}$ ) and salinity (psu) sections for January 2010 (left) and June (right). Salinity is derived from the mean T-S relationship using WOD01. The name of the eastward currents are written over their mean positions in panel (a).

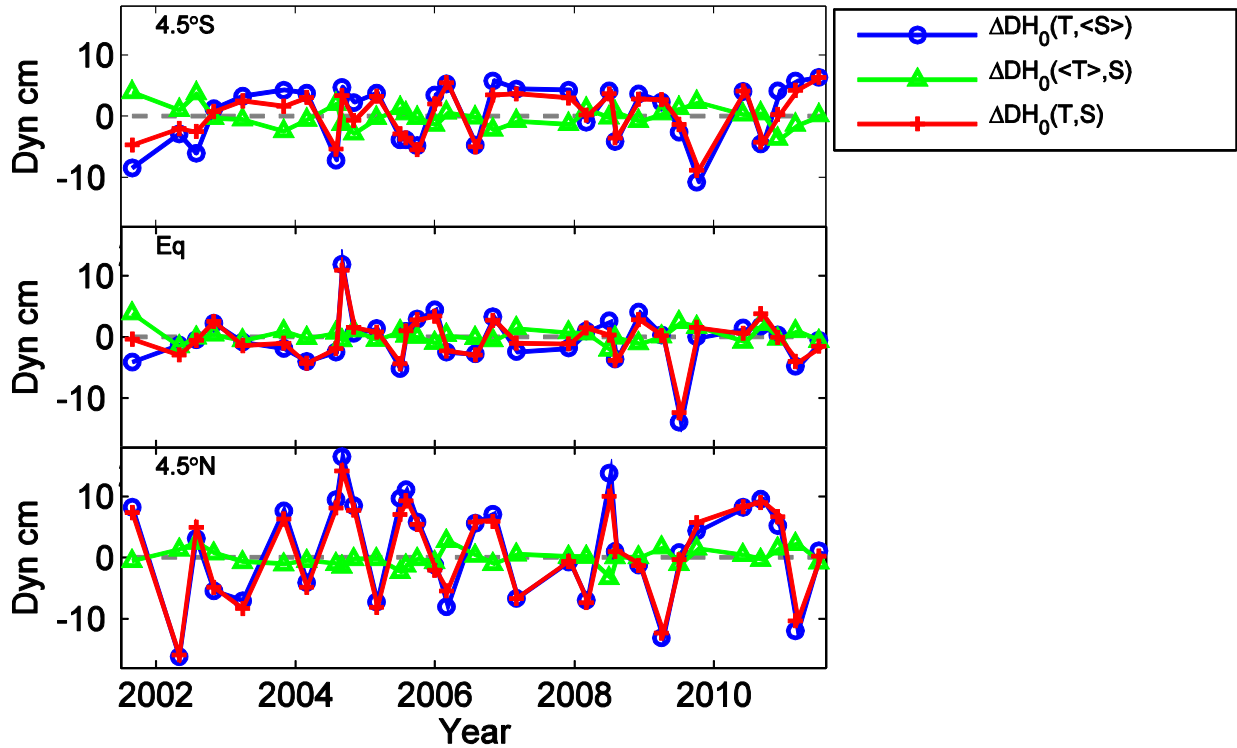
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804 Figure 3: Mean seasonal cross-sections of velocity derived from XBT data within the months of  
805 a) Dec, Jan, Feb, b) Mar, Apr, May, c) Jun, Jul, Aug, d) Sep, Oct, Nov. The black lines along  
806 each section are the depth of the seasonal mean depths of 24.5 and 26.8 isopycnal. The top panel  
807 of each subfigure is the seasonal mean absolute dynamic height.

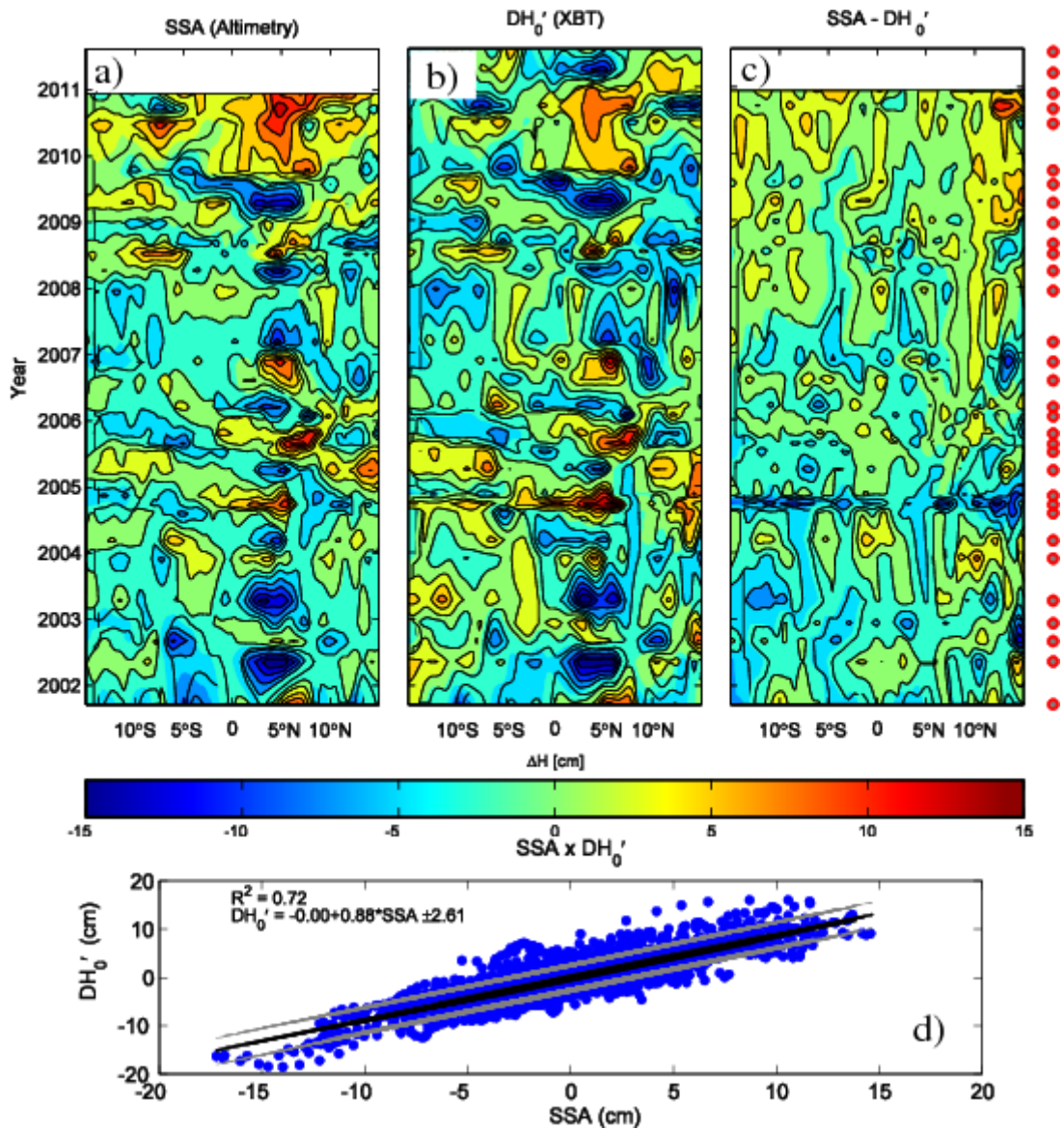
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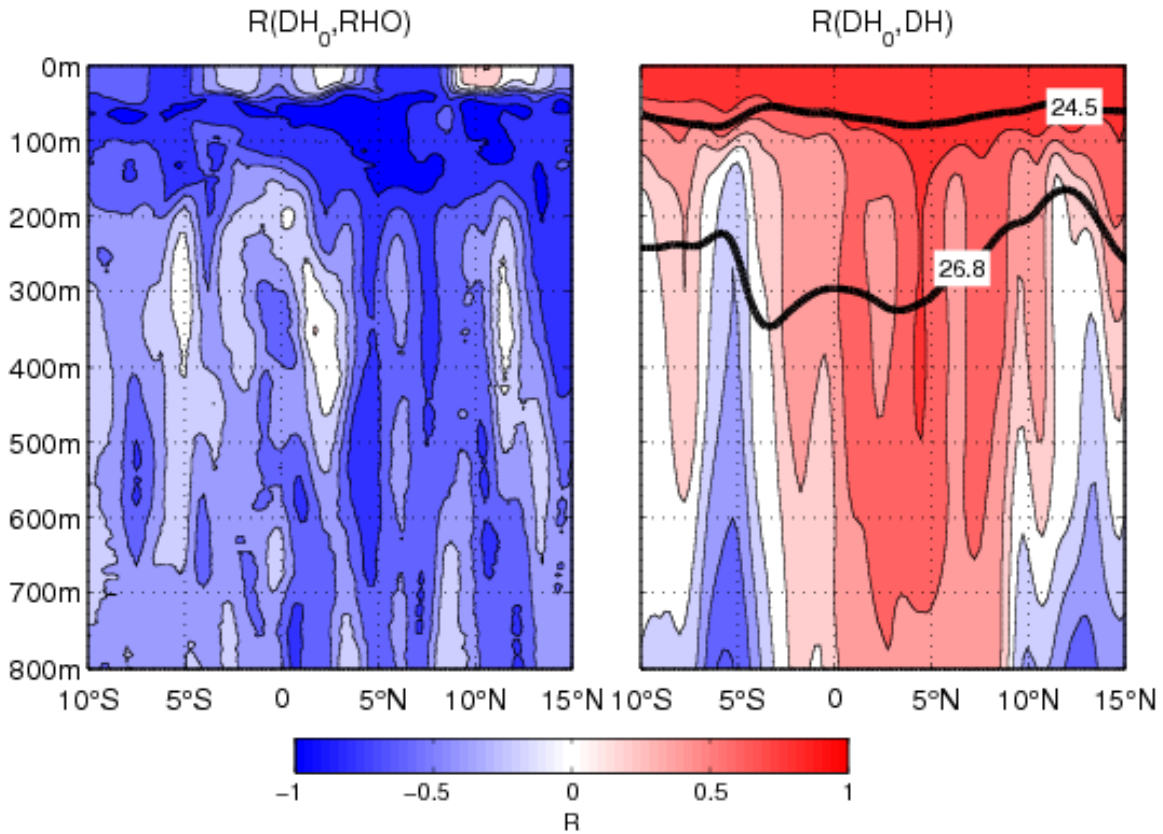
Figure 4: Anomalies of surface DH ( $DH_0'$ ) with respect to the annual mean of the XBT data (red) and their components thermosteric (blue) and halosteric (green). The three panels are for different latitudes,  $4.5^\circ\text{S}$  (upper),  $0^\circ\text{S}$  (middle), and  $4.5^\circ\text{N}$  (lower).





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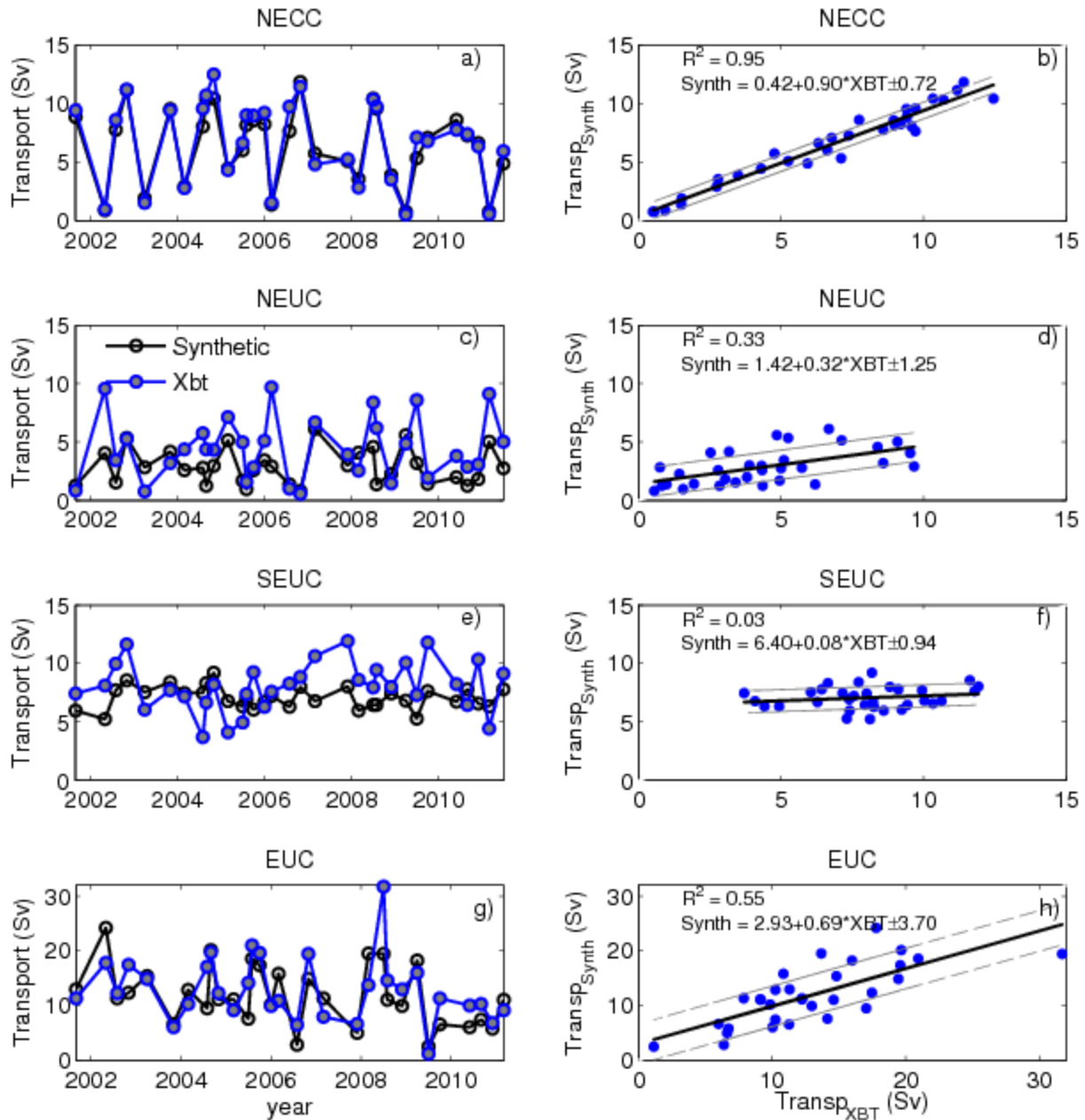
Figure 5: Comparison between SSA and surface dynamic height anomalies ( $DH_0'$ ) from hydrography. Top panels: longitude-time diagrams of (a) SSA, (b)  $DH_0'$ , and (c)  $SSA - DH_0'$ . Dots on the right-hand side of the figure mark the realizations of the XBT transect. Bottom panel: (d) linear fit between  $DH_0'$  and SSA.



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 840 Figure 6: Correlation at each depth and latitude between the dynamic height anomalies (cm) at  
 841 the surface ( $DH_0$ ) and: a) density anomalies ( $\sigma_{\theta}$ ;  $\text{kg m}^{-3}$ ) and b) dynamic height anomalies ( $DH'$ ).  
 842 The thick black lines in b) mark the mean location of the isopycnals  $\sigma_{\theta} = 24.5 \text{ kg m}^{-3}$  and  $\sigma_{\theta} =$   
 843  $26.8 \text{ kg m}^{-3}$  that define the upper and lower dynamic layers used in this study.

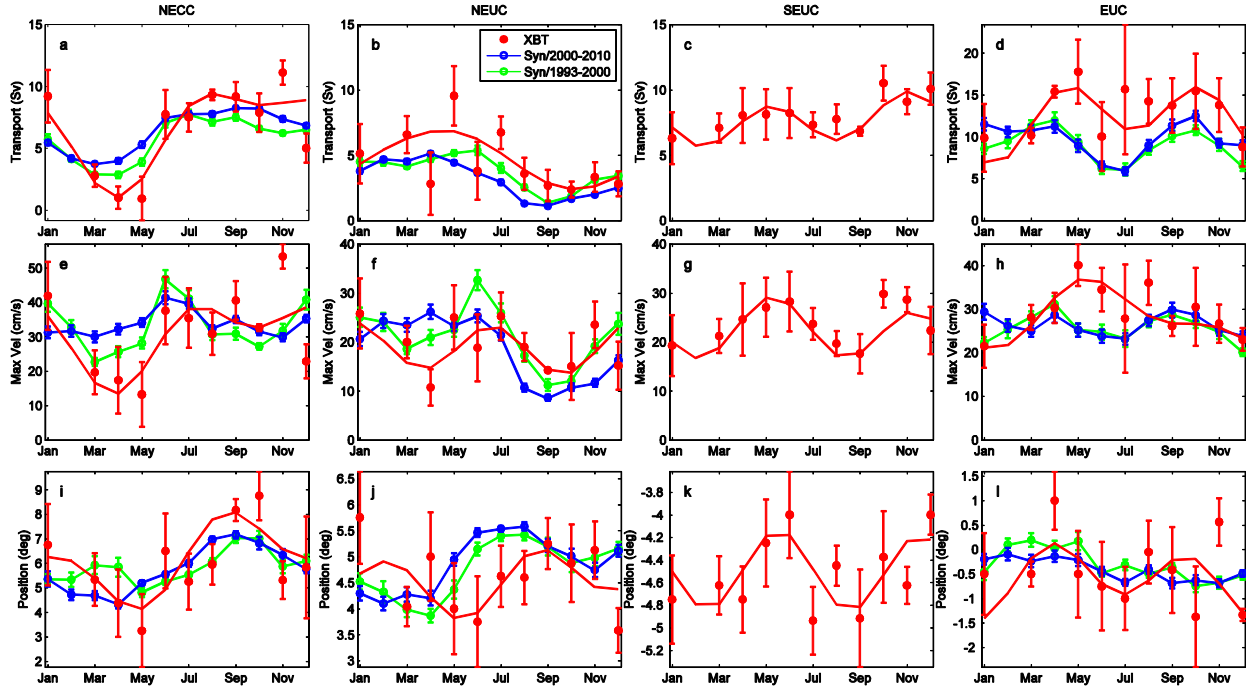
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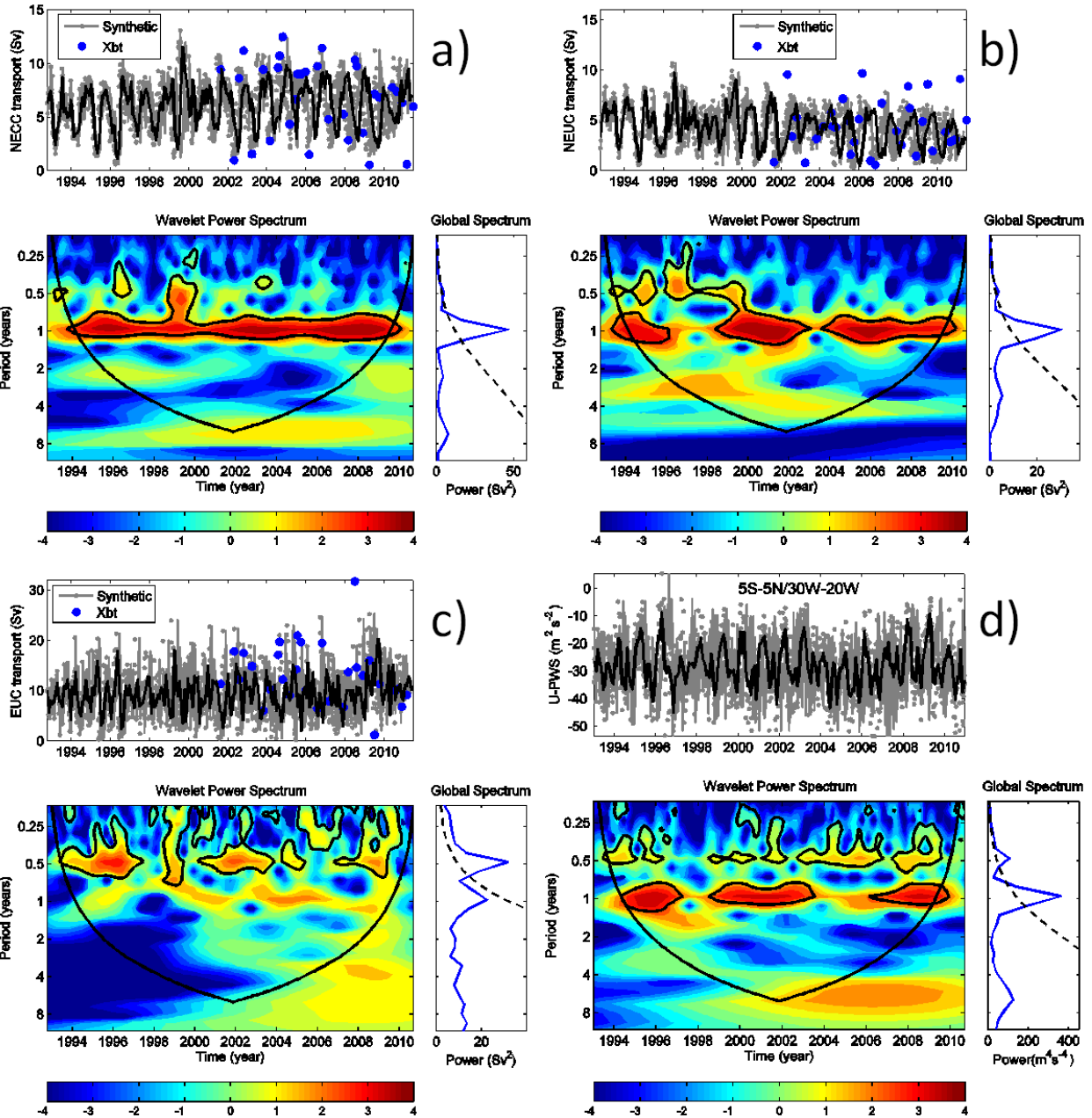
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Figure 7: Left: geostrophic transports estimated by the synthetic method (black line with open dots) and transports estimated using XBT data (blue dots) for (a) NECC, (c) NEUC, (e) SEUC, and (g) EUC. Right: corresponding linear fit between the two transports estimates.



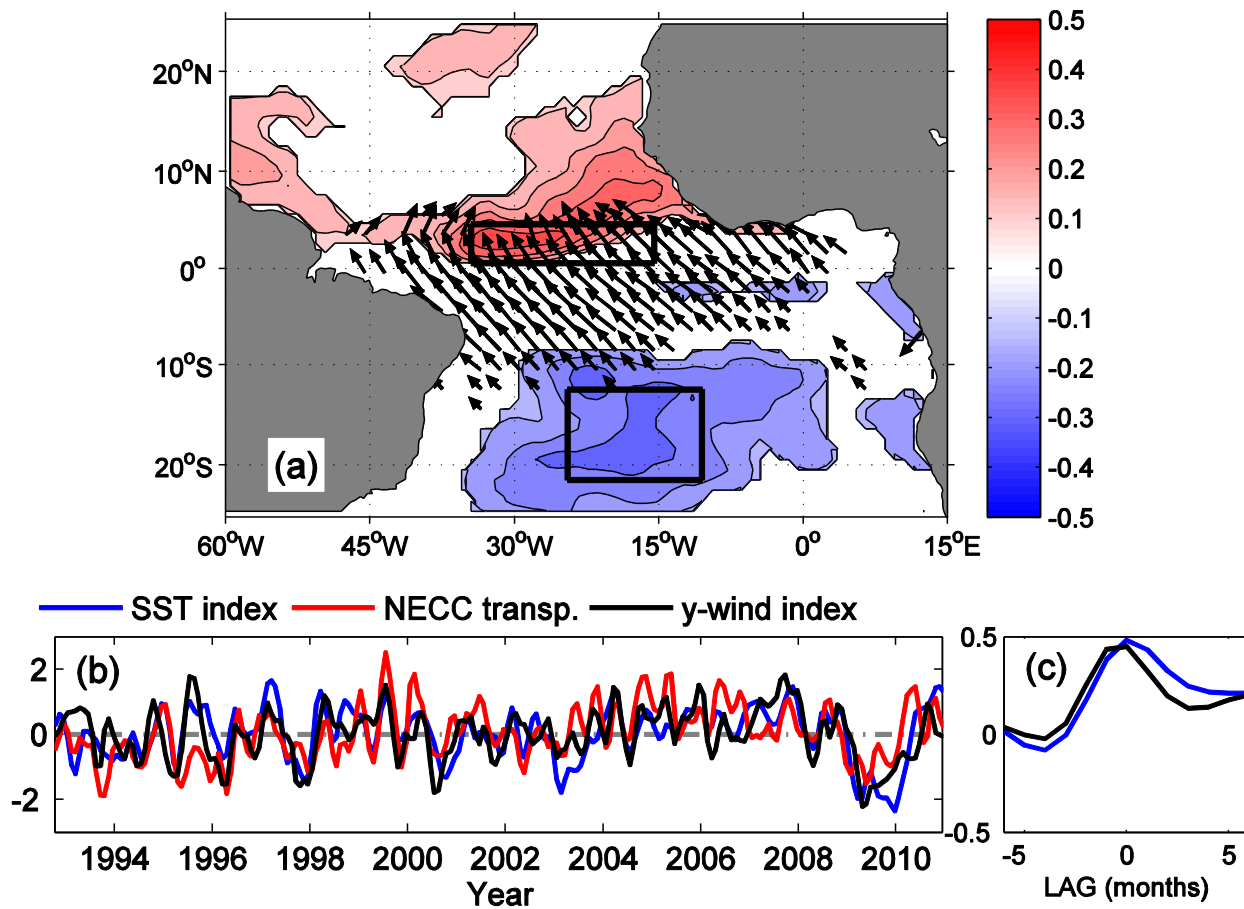
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Figure 8: Climatology of geostrophic transport (Sv) (upper), maximum velocity (cm/s) (middle), and latitudinal position (deg) (lower) for the NECC (a,e,i), NEUC (b,f,j), SEUC (c,g,k) and EUC (d,h,l). Red dots are monthly XBT averages and red line represents the corresponding fit of first and second harmonics (Foltz et al., 2004a) for the XBT measurements; green/blue lines mark the average synthetic estimates for the period 1993-2000/2000-2011. The errorbars are the standard errors for the synthetic estimates, and the standard error plus fitting error as calculated in Foltz et al. (2004a).



896 Figure 9: Time series and respective wavelet transforms for transport of (a) NECC, (b) NEUC (c)  
 897 EUC, and d) zonal pseudo-wind stress averaged inside the 5S-5N/30W-20W box. The current  
 898 transport timeseries (a, b, and c) are generated using the synthetic method (gray) with the  
 899 discrete XBT estimates plotted on top of each timeseries (blue dots). The monthly means  
 900 timeseries (black) are overlaid on each timeseries. The wavelet power spectra are based on a  
 901 Morlet transform on the monthly timeseries and are shown in the contour plots underneath their  
 902 respective timeseries. The wavelet power spectrum panels show the regions above the 95%  
 903 significance level bounded by black closed contours, and the bowl-shaped black lines are the  
 904 cone of influence above which the spectral values can be considered. Also plotted are the global  
 905 averages of the wavelet spectrum (blue) and respective significance level (black dashed), as a  
 906 right side panel of each wavelet power spectrum.  
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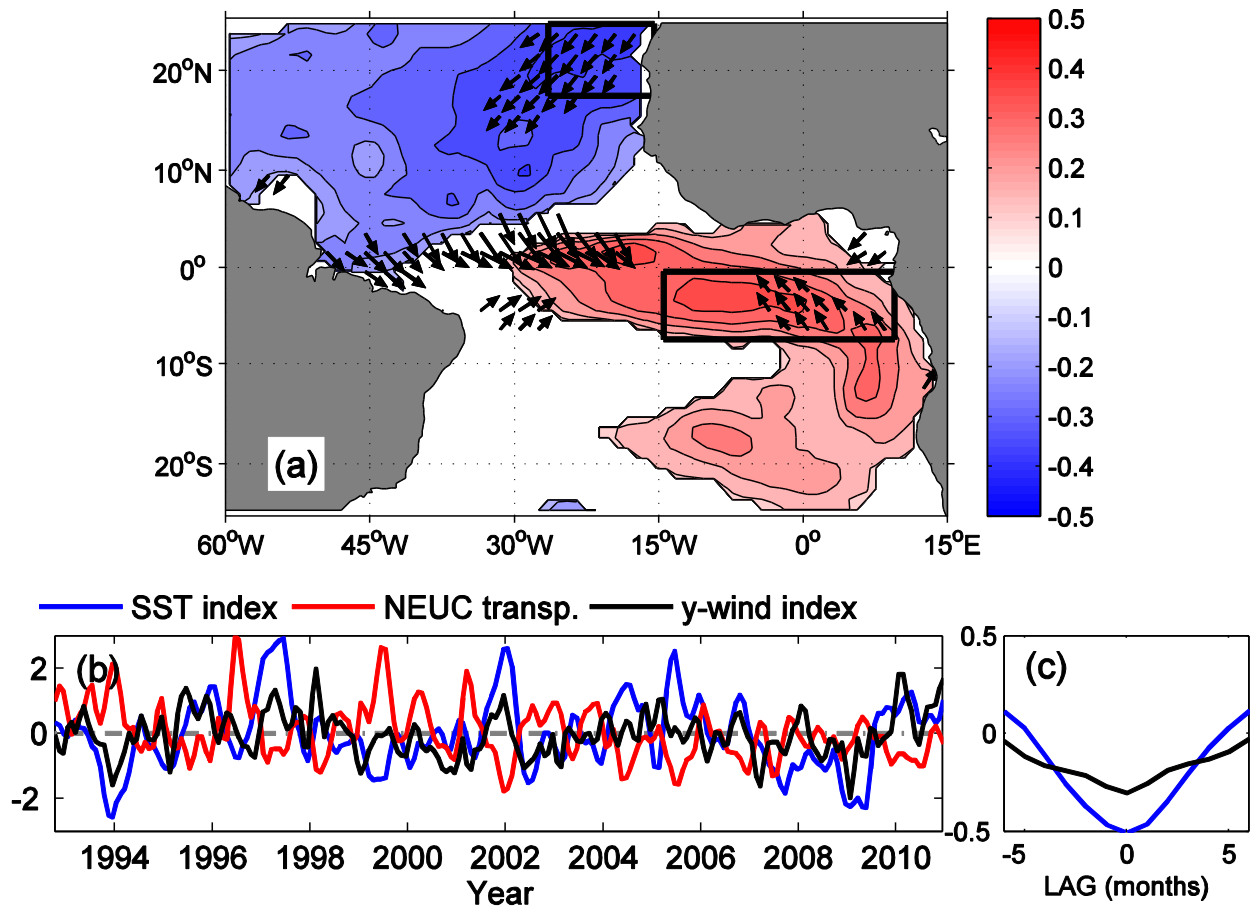
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Figure 10: a) Instantaneous correlation between the NECC transport anomalies timeseries with SSTA (color) and CCMP pseudo-wind stress anomalies (arrows). Only the values that are statistically significant are shown. b) Time series of the standardized NECC transport anomalies (red), the SSTA index (blue), and the meridional wind index (y-wind; black). The SSTA index defined by subtracting the averages within the northern box (35°W-17°W/0-7°N) and the southern box (30°W-15°W/11°S-23°S) as marked in the panel a), and the y-wind index is defined as the average of the meridional pseudo-wind stress anomalies inside the northern box. c) Lagged correlation between the NECC index with the SSTA index (blue) and the y-wind index (black).

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Figure 11: a) Same as figure 9a, but for NEUC transport anomalies. In b) and c) the time series of the SSTA index (blue) is defined as the difference between the northern box over the Guinea Dome region (15°-25°N/15°-30°W) and the southern box (8S-0/15W-10E), and the y-wind index is defined for the northern box as marked in panel a).

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Table 1: Percentage of the variance of transport, core velocity, and position explained by the annual and semi-annual harmonics for each current using XBT estimates, and synthetic method for the 1990-2000 mean (S1990-2000) and 2000-2011 mean (S2000-2011). The synthetic method is not analyzed for the SEUC.

Current	Data	Transport (%)		Velocity (%)		Position (%)	
		Annual	Semi-ann	Annual	Semi-ann	Annual	Semi-ann
NECC	XBT	61	12	34	18	50	15
	S/1990-2000	78	17	3	70	60	20
	S/2000-2011	92	4	27	23	99	0.1
NEUC	XBT	54	1	1	49	11	26
	S/1990-2000	72	26	25	63	75	21
	S/2000-2011	99	1	87	5	78	1
SEUC	XBT	11	90	8	92	1	58
EUC	XBT	25	80	72	17	8	34
	S/1990-2000	10	81	15	75	58	18
	S/2000-2011	30	55	8	40	83	8

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