

1           **Variability of eastward currents in the equatorial**  
2                           **Atlantic during 1993-2010**

3  
4   Marlos Goes <sup>1,2\*</sup>, Gustavo Goni<sup>2</sup>, Verena Hormann<sup>1,2</sup>, and Renellys C. Perez<sup>1,2</sup>

5  
6  
7  
8  
9  
10  
11   Submitted to Journal of Geophysical Research-Oceans

12   18- May-2012

13  
14  
15  
16  
17  
18   <sup>1</sup> Cooperative Institute for Marine and Atmospheric Studies, University of Miami, Miami, FL  
19   33149, USA.

20   <sup>2</sup> Atlantic Oceanographic and Meteorological Laboratory, National Oceanic and Atmospheric  
21   Administration, Miami, FL 33149, USA.

22

23 **Abstract**

24         We have developed, validated, and applied a synthetic method to monitor the eastward  
25 currents in the equatorial Atlantic. This method combines high-density expendable  
26 bathythermograph (XBT) temperature data along the AX08 transect crossing the equator at about  
27 23°W with altimetric sea level anomalies (SLAs) to estimate dynamic height fields from which  
28 mean properties of the currents and their variability can be estimated on seasonal to interannual  
29 timescales. The method is well suited for the surface North Equatorial Countercurrent (NECC)  
30 and reproduces the variability of the North Equatorial and Equatorial Undercurrents (NEUC,  
31 EUC) with considerable skill. The synthetic method is unable to describe variations of the South  
32 Equatorial Undercurrent (SEUC), which is located in a region of low SLA variability. Our results  
33 confirm that the NECC shows a strong annual cycle of volume transport, with largest values  
34 from July to December. On interannual timescales, there is a positive correlation between the  
35 NECC transport, an interhemispheric sea surface temperature (SST) gradient and the  
36 southeasterly winds. The NEUC reveals largest transport values (up to 10 Sv) from January to  
37 July and is correlated on interannual timescales with SSTs in the Gulf of Guinea and  
38 southeastern equatorial Atlantic as well as zonal equatorial winds. The EUC exhibits strong  
39 semi-annual and annual variability. This study shows that for a long-term monitoring system  
40 both altimetry and XBT data are needed for near-real-time inference of dynamic and  
41 thermodynamic properties of the tropical Atlantic current system.

42

43

## 44 **1. Introduction**

45           The tropical Atlantic current system is of great importance for both interhemispheric and  
46 west-to-east exchange of heat and nutrients [e.g., Foltz et al., 2003; Brandt et al., 2008], and also  
47 impacts the climate and weather in the surrounding continental areas [e.g., Goldenberg et al.,  
48 2001; Sutton and Hodson, 2005; Stramma et al., 2005; Brandt et al., 2011], since ocean upper  
49 dynamics play an important role in the modes of sea surface temperature variability in this region  
50 [e.g., Carton and Huang, 1994; Carton et al., 1996]. In the upper equatorial Atlantic, the main  
51 conduit for the interhemispheric water exchange is the western boundary current system of the  
52 North Brazil Undercurrent/Current (NBUC/NBC). On its northward path, the NBC partly  
53 retroflects eastward and feeds into a system of zonal countercurrents, namely the Equatorial  
54 Undercurrent (EUC) along the equator, and three off-equatorial currents, the North Equatorial  
55 Countercurrent (NECC), the North Equatorial Undercurrent (NEUC), and to a lesser extent the  
56 South Equatorial Undercurrent (SEUC) [e.g., Metcalf and Stalcup, 1967; Cochrane et al., 1979;  
57 Peterson and Stramma, 1991; Schott et al., 1995]. Moreover, the eastward equatorial currents  
58 play an important role in the global meridional overturning circulation by connecting the western  
59 boundary regime to the interior circulation [e.g., Frantantoni et al., 2000], and in the shallow  
60 subtropical cells (STCs) by providing water for the off-equatorial eastern upwelling regimes  
61 along the African coast, specifically for the Guinea and Angola Domes [e.g., Hisard and Henin,  
62 1987; Zhang et al., 2003; Marin et al., 2003; Hua et al., 2003; Schott et al., 2004; Doi et al.,  
63 2007].

64           The primary equatorial branch of the STCs is the EUC, which is unique in its dynamics  
65 [e.g., Moore and Philander, 1986; Pedlosky, 1988; Wacongne, 1989] and has been found to be  
66 linked with the Atlantic zonal mode on interannual timescales [e.g., Goes and Wainer, 2003;

67 Hormann and Brandt, 2007]. The seasonal cycle of the EUC has been widely analyzed in model  
68 simulations [e.g., Philander and Packanowski, 1986; Hazeleger et al., 2003; Arhan et al., 2006;  
69 Hormann and Brandt, 2007] which generally agree that the EUC shows two transport maxima in  
70 the central part of the basin, one during boreal summer/fall and another during boreal  
71 winter/spring. Arhan et al. [2006] suggested that two different dynamic regimes drive the  
72 seasonal cycle of the EUC: In boreal summer/fall, the simulated EUC is mostly driven locally by  
73 equatorial zonal wind forcing, and is supplied from the ocean interior; in boreal winter/spring, it  
74 is driven by remote forcing through the rotational wind component, and is supplied from the  
75 western boundary currents. Observational studies, however, have drawn different conclusions,  
76 either by showing one sole maximum in late boreal summer/autumn [Hisard and Henin, 1987],  
77 or no significant seasonal variations [Weisberg et al., 1987].

78         Due to its strong surface expression, the NECC has been the most extensively studied of  
79 the three eastward off-equatorial currents. The NECC is mainly located between 3°-10°N [e.g.,  
80 Garzoli and Katz, 1983; Richardson and Walsh, 1986], and splits seasonally into two discernible  
81 cores that are dynamically linked to the widening of the zonal wind stress band [Urbano et al.,  
82 2006]. The seasonal variability of the NECC is strongly tied to the migration of the Intertropical  
83 Convergence Zone (ITCZ) [Katz and Garzoli, 1982; Katz, 1987; Garzoli and Richardson, 1989;  
84 Garzoli, 1992]. In addition to local wind stress curl and Rossby waves from the eastern basin  
85 [e.g., Garzoli and Katz, 1983; Katz, 1987; Garzoli and Richardson, 1989], wind stress  
86 fluctuations in the equatorial band have been found to play an important role in driving the  
87 seasonal cycle of the NECC [Yang and Joyce, 2006]. The NECC also exhibits considerable year-  
88 to-year variability [e.g., Katz, 1993; Fonseca et al., 2004] and Hormann et al. [2012] recently

89 showed that its interannual variability is linked to the two dominant tropical Atlantic climate  
90 modes.

91         The variability of the off-equatorial undercurrents (OEUCs), though, is still not well  
92 understood. This is mostly related to observational constraints, since the OEUCs are weaker,  
93 diffusive, and highly variable on intraseasonal timescales due in large part to tropical instability  
94 wave (TIW) activity [Johnson and Moore, 1997; Rowe et al., 2002; Jochum et al., 2003; Jochum  
95 and Malanotte-Rizzoli, 2004]. The OEUCs are located below the thermocline, coincident with  
96 the equatorial thermostat, at about  $\pm 3^{\circ}$ – $6^{\circ}$  of latitude [Cochrane et al., 1979]. The SEUC is  
97 mostly fed by internal recirculations [Schott et al. 1995, 1998], and according to Reverdin et al.  
98 [1991], its seasonal cycle at  $30^{\circ}$ W is characterized by a transport maximum in boreal fall and a  
99 minimum in boreal spring. Close to its origin, observations show that the SEUC flow consists of  
100 large standing meanders [Fischer et al., 2008]. The NEUC is weaker than the SEUC, and more  
101 variable due to enhanced TIW activity in the northern hemisphere [e.g., Foltz et al., 2004b; Athie  
102 and Marin, 2008]. For instance, Schott et al. [2003] analyzed an average of 13 shipboard sections  
103 at  $35^{\circ}$ W, and stated as uncertain the existence of the NEUC. Moreover, the separations between  
104 the NEUC and EUC in the western part of the basin [e.g., Bourlès et al., 1999; Goes et al., 2005]  
105 and the NEUC and NECC in the central part of the basin [Brandt et al., 2010] are not very  
106 distinct.

107         All the eastward currents in the equatorial Atlantic are modified by wind forcing and by  
108 oceanic mesoscale phenomena, such as TIWs [e.g., Düing et al., 1975; Weisberg and Weingartner,  
109 1988; Menkes et al., 2002; Grodsky et al., 2005] that can alias estimates of seasonal and  
110 interannual variability obtained from observational systems that are not continuous in time,  
111 posing a significant sampling challenge. An existing observational system that can potentially

112 resolve the short spatial and temporal scales, and allow for quantification of the seasonal to  
113 interannual variability of the eastward currents is satellite altimetry. Although upper-ocean  
114 currents in the tropical Atlantic having very weak sea surface height signatures cannot be  
115 resolved from surface topography fields alone [Goni and Baringer, 2002], a relationship between  
116 certain current features and their characteristic sea surface height signatures can be established to  
117 overcome this challenge. This can be achieved by combining altimetry with other observational  
118 platforms, such as high-resolution hydrographic data. Therefore, a reliable monitoring system for  
119 the region requires complementary platforms in order to produce details of the spatio-temporal  
120 variability [Goni and Baringer, 2002].

121         The goal of this study is to quantify the variability of the eastward currents in the  
122 equatorial Atlantic during 1993-2010, and determine their properties, such as transport, velocity,  
123 and location. Specifically, we focus on the seasonal to interannual signatures of the NECC, EUC,  
124 and OEUCs manifested in the observations. To accomplish this goal, we consider a monitoring  
125 system that relies upon data from the cross-equatorial High Density eXpendable  
126 BathyThermograph (HD XBT) transect AX08 during 2000-2010, and use satellite altimetry  
127 observations to synthetically quantify the variability of the eastward currents in the equatorial  
128 Atlantic for the 1993-2010 period. The HD XBT project has been active for over 20 years, and  
129 aims at sustainably measuring physical properties of the upper ocean with mesoscale resolution.  
130 Its high spatial data resolution and repeated sampling of the region enable assessment of upper-  
131 ocean temperature and heat storage variability, and permit characterizing of the variability of the  
132 major geostrophic currents.

133         This paper is structured as follows: First we describe in detail the data to be used (Section  
134 2) and introduce the synthetic methodology (Section 3); in Section 4, we validate the synthetic

135 method, apply it to examine the seasonal and interannual variability of the eastward equatorial  
136 currents, and analyze the surface response and forcings linked to this variability. Discussion,  
137 conclusions, and recommendations for further development of the methodology are provided in  
138 Section 5.

139

## 140 **2. Data**

141 This work mainly uses two observational platforms to infer the variability of the eastward  
142 currents in the equatorial Atlantic: a high-density XBT transect and satellite altimetry.  
143 Specifications for each of these datasets, as well as for additional data used in the analysis are  
144 given below.

145

### 146 a) Hydrographic data

147 The temperature data used in this study are from 39 realizations of the AX08 XBT  
148 transect, which is carried out between Cape Town and New York City and crosses the equator at  
149 about 23°W (Figure 1a). The first section was obtained in December 2000, and an average of  
150 four transects per year has been achieved since 2002. Approximately 200-300 XBTs are  
151 deployed during each transect realization, with a nominal spacing of 25 km between casts in the  
152 tropics and a maximum depth of roughly 800 m.

153 In order to study the variability of the eastward equatorial currents, we restrict the  
154 observations to a region with similar dynamic characteristics. We define a criterion that selects  
155 the sections whose mean longitude lies within the 68 percentile around their median value  
156 between 10°S-10°N (Figure 1b). The median longitude of the transect, which is about 40°  
157 oblique with respect to a true meridional section, is approximately 23°W (Figure 1a).

158 Additionally we exclude the September 2004 section because its derived surface dynamic height  
159 shows an offset in comparison to the other sections. The applied constraints reduce the number  
160 of transects from 39 to 31, but assures that we are working with comparable data.

161 XBT measurements are performed by sampling temperature and the time of descent of  
162 the XBT probes. The elapsed time of descent is converted to depth by applying a fall rate  
163 equation [Hanawa et al., 1995]. Individual temperature profiles are linearly interpolated onto a 2  
164 m grid. The data are quality controlled by excluding the outlier profiles, chosen as the profiles  
165 whose both forward and backward horizontal gradients of the surface dynamic height lie outside  
166 the three standard deviation range of all profiles. Next, the sections are horizontally interpolated  
167 to a 25 km resolution in latitude, using an optimal interpolation scheme based on a Gaussian  
168 correlation function with a decorrelation length scale of 200 km and a low noise-to-signal ratio  
169 of 0.01. Salinity is inferred from the temperature profiles using climatological temperature-  
170 salinity (T-S) relationships extracted from the World Ocean Database (WOD01, Conkright et al.  
171 [2002]). In the mixed layer, salinity profiles are extrapolated using a slab-layer approximation  
172 (Figure 2c, d), which is a standard approximation to overcome the non-unique characteristic of  
173 the T-S relationship in the tropical surface waters [Schott et al., 1998; Goes et al., 2005]. The  
174 differences between in-situ and climatological salinity have been shown to be of the order of 0.3-  
175 0.4 psu, and the resulting uncertainty to be among the largest contributors to the dynamic height  
176 error, with differences as large as 5 cm [Goni and Baringer, 2002].

177 Typical examples of temperature and salinity distributions along the AX08 section in  
178 boreal winter and summer are shown in Figure 2. This region is characterized by a warm well-  
179 mixed layer in the top 100 m, followed by a sharp temperature gradient of approximately  
180  $0.1^{\circ}\text{C}/\text{m}$  that marks the upper thermocline. The waters to the north and south of the displayed



181 domain are characterized by high salinity values at around 100 m, within the upper thermocline  
182 waters. These high salinity waters are characteristic of the Subtropical Underwater (SUW).  
183 SUWs are formed in the subtropics and advected equatorward by the North and South Equatorial  
184 Currents. Underneath the SUW is the central water (CW), characterized by a nearly straight line  
185 in the T-S space [e.g., Stramma and England, 1999]. The CW forms the thermostat between 12-  
186 15°C found in the equatorial region [Reverdin et al., 1991], which is more pronounced in the  
187 southern hemisphere between 5°S and the equator at around 200 m depth (Figure 2), and  
188 coincides with the position of the OEUCs [e.g., Cochrane et al., 1979; Schott et al., 1995, 1998].

189

#### 190 b) Altimetric data

191 Here we use the AVISO reference (“ref”) delayed mode product  
192 (<http://www.aviso.oceanobs.com>) obtained from a multi-satellite mission [Le Traon et al., 1998],  
193 which provides gridded, optimally interpolated, and cross-calibrated global coverage of sea level  
194 anomalies (SLAs) relative to the 1993-1999 mean. These delayed mode data are continuously  
195 available on a 1/3° horizontal grid with weekly temporal resolution since October 1992, and SLA  
196 precision of 2 cm [Ducet et al, 2000; Cheney et al., 1994]. In this study, we use data from  
197 October 1992-December 2010 and subtract the 2000-2010 mean from the SLA field for  
198 consistency with the XBT dataset. We further interpolate SLA linearly onto to the location and  
199 time of the individual XBT sections to estimate regression parameters for the synthetic method,  
200 and onto the mean AX08 transect to create a hindcast of the currents for the whole altimetric  
201 period.

202

#### 203 c) Additional data

204 We use referenced dynamic height from the monthly climatology of the International  
205 Pacific Research Center (IPRC) (<http://apdrc.soest.hawaii.edu>). This dataset is available on a 1°  
206 x 1° horizontal grid at 27 depth levels and is derived from Argo floats and altimetry  
207 observations, with mean sea level referred to the mean dynamic topography MDT\_CNES-CLS09  
208 [Rio et al., 2011].

209 We perform additional analyses using gridded sea surface temperature (SST) and pseudo-  
210 wind stress anomalies with respect to their climatological means. These data are monthly  
211 averaged for the considered altimetric period, and interpolated onto a 1° x 1° Mercator grid. The  
212 SST data are extracted from the optimum interpolation (OISST-v2) analysis [Reynolds et al.,  
213 2002], which is available weekly on a 1° horizontal grid since November 1981. The pseudo-wind  
214 stress data are obtained from the cross-calibrated, multi-platform (CCMP), multi-instrument  
215 ocean surface wind velocity data set. This product combines data derived from SSM/I, AMSRE,  
216 TRMM TMI, Quikscat, and other missions using a variational analysis method to produce a  
217 consistent record of ocean surface vector winds at a 25 km resolution and is available since July  
218 1987 [Atlas et al., 2011].

219

### 220 **3. Methodology**

#### 221 a) *Velocity calculation*

222 We calculate the cross-sectional absolute geostrophic currents from horizontal gradients  
223 of absolute dynamic height (DH(z)) using the thermal wind relationship. To obtain DH(z), we  
224 reference the XBT-derived dynamic height (DH<sub>XBT</sub>) to the IPRC monthly climatology of  
225 absolute dynamic height at 800 m (DH<sub>IPRC</sub>(800m)) for each XBT transect:

$$226 \quad DH(z) = DH_{XBT}(z) + DH_{IPRC}(800m) \quad (1)$$

227 Although the inclusion of  $DH_{IPRC}(800m)$  alters the mean DH significantly, DH gradients are not  
228 greatly affected ( $O(10^{-7})$  cm/km change). Geostrophy has an inflection point at the equator and it  
229 is necessary to use the equatorial beta approximation, which relies on the calculation of higher  
230 order derivatives of the  $DH(z)$  near the equator. In this study, we apply the method of Lagerloef  
231 et al. [1999] for velocity calculations within  $\pm 3^\circ$  off the equator, except here we use a 6th order  
232 polynomial rather than a 3rd order polynomial.

233 We further explore the effect of salinity on the steric variability in the XBT data by  
234 calculating the thermosteric and halosteric contributions to the total surface dynamic height  
235 variability at three different latitudes,  $4.5^\circ N$ ,  $4.5^\circ S$ , and the equator (Figure 3). These latitudes  
236 roughly coincide with the core of the eastward currents. The calculation of the thermosteric and  
237 halosteric contributions is performed using the methodology described by Fofonoff and Froese  
238 [1958] and Tabata et al. [1986] in which departures from the mean sea level are estimated in  
239 terms of one component by keeping the other component fixed at its annual mean. The sum of  
240 the thermosteric and halosteric components is approximately equal to the total steric variability  
241 of the XBT data. We illustrate their contributions to the surface dynamic height ( $DH_0$ ) referenced  
242 to 800 m and at all three latitudes (Figure 3). Results indicate that the variability of total  $DH_0$  is  
243 largely explained by its thermosteric component, and that the halosteric contribution to the  
244 variability of  $DH_0$  is rather small, generally of the order of 1-2 cm. Halosteric effects have a  
245 larger contribution to the total  $DH_0$  variability at the equator and  $4.5^\circ S$  (Figure 3b,c), where the  
246 total  $DH_0$  variability is small compared to  $4.5^\circ N$ . At those latitudes, the halosteric component  
247 accounts for 20-30% of the total variability, while at  $4.5^\circ N$  (Figure 3a), where there is high  
248 thermosteric variability with amplitudes of up to 20 cm, only about 10% of the  $DH_0$  variability is  
249 due to the halosteric component.

250           The a priori errors in geostrophic velocity resulting from our methodology include XBT  
251 temperature precision of  $\pm 0.1^\circ\text{C}$  ( $O[1 \text{ cm/s}]$ ), the use of salinity estimates from climatological T-  
252 S relationships ( $O[5 \text{ cm/s}]$ ), and the uncertainty in the level of known motion ( $O[1 \text{ cm/s}]$ ). An  
253 additional error results from the assumption that the cross-sectional geostrophic velocity along an  
254 oblique transect (Figure 1a) is approximately equal to the zonal geostrophic velocity estimated  
255 from a meridional transect. This error has been previously quantified in a similar region as  
256 smaller than 10% ( $O[4 \text{ cm/s}]$ ) by Reverdin et al., [1991]. We performed tests with satellite  
257 altimetry data (not shown here) and concluded that 10% error is a conservative estimate due to  
258 the large random variability in the region. Assuming the errors above are random and  
259 uncorrelated, their total contribution would result in an error of  $\pm 6.5 \text{ cm/s}$ , which is on the order  
260 of or smaller than the intrinsic mesoscale variability of the region [e.g., Grodsky et al., 2005;  
261 Brandt et al., 2006]. Therefore, the standard errors of the monthly averages will be used as a  
262 confidence interval for our calculations. Near the equator, the beta plane approximation  
263 introduces errors up to one order of magnitude greater ( $O[10 \text{ cm/s}]$ ) than the other components,  
264 and this caveat has to be taken into consideration in the analysis of the EUC results.

265

#### 266 b) *Synthetic method*

267           Altimetry data is used here along with those of other observational platforms to provide a  
268 broader four-dimensional (i.e., spatial and temporal) coverage of the tropical Atlantic. Several  
269 studies have combined altimetric and hydrographic data to infer properties of the upper ocean,  
270 such as velocity, temperature, and salinity [e.g., Gilson et al., 1998; Ridgway, 2002; Phillips and  
271 Rintoul, 2002].

272 Here, we apply a synthetic method to produce a hindcast of the velocity and density fields  
273 along the mean AX08 transect position (red line in Figure 1a). Since we are interested in velocity  
274 sections along isopycnal layers, we seek as predictands potential density ( $\sigma_\theta$ ) and DH from the  
275 surface to 800 m and use altimetric SLA data as predictors. To determine the relationship  
276 between the predictors and predictands, we first calculate the absolute DH (Equation 1) for the  
277 31 selected XBT sections. Anomalies of  $\sigma_\theta$  and DH ( $\sigma'_\theta$  and  $DH'$ , respectively) are calculated by  
278 subtracting a mean annual field, defined by the mean WOD01 density field along the mean  
279 section and the mean dynamic height field of all sections. For the latter, we first calculate  
280 monthly averages to reduce potential sampling biases toward any specific month (Figure 1c).  
281 These results are largely insensitive to the mean fields since only a constant is subtracted from  
282 the temporal fields, but removing a mean is known to also reduce the variance of residuals in a  
283 linear regression. Anomalies of  $DH_0$  ( $DH'_0$ ) are linearly regressed onto  $DH'$  and  $\sigma'_\theta$  as a function  
284 of latitude and depth. Finally, the  $DH'_0$  is linearly regressed onto SLA forming the link between  
285 the altimetric and hydrographic observations. The  $DH'_0$  and SLA are well correlated (Figure 4d),  
286 with  $R^2 = 0.75$  (the R-squared coefficient corresponds to the squared correlation in a simple  
287 linear regression), a root mean square error of  $RMSE = 2.40$  cm, and a negligible intercept. This  
288 strong and unbiased relationship between the two variables shows that the SLA captures well the  
289 baroclinic structures in the region, especially the first mode [Gilson et al., 1998; McCarthy et al.,  
290 2000; Guinehut et al., 2006]. The highest SLA variance occurs between  $5^\circ$ - $10^\circ$ N, coinciding  
291 with the highest SLA gradients which are closely related to the dynamics of the NECC. Less  
292 variability is observed south of the equator, which may be due to compensating effects in the  
293 water column [Mayer et al., 2001]. The differences found here between SLA and  $DH'_0$  can arise  
294 from a number of factors, such as temporal and spatial sampling, incomplete removal of tides, as

295 well as barotropic flow variations and changes in the entire depth of the ocean [Rintoul et al.,  
296 2002].

297 Here we use the temporal correlation between  $DH'_0$  and the predictands at each depth and  
298 latitude as an indicator of the skill of this method to monitor the variability of the predictands  
299  $DH'$  and  $\sigma'_\theta$  (Figure 5). As expected, these correlations are predominantly positive for  $DH'$  and  
300 negative for  $\sigma'_\theta$  since dynamic height is calculated from specific volume anomalies, which are  
301 inversely related to density [Pond and Pickard, 1983]. Apart from their sign, both fields show  
302 similar relationships, with highest correlations in the upper 200 m of the water column, and  
303 decreasing correlation values with depth. In certain latitudes, such as in the vicinity  $5^\circ\text{N}$ , high  
304 correlations ( $R \approx 0.8$ ) can reach depths of 800 m, similar to values encountered in the North  
305 Pacific [Roemmich and Gilson, 2001].

306

## 307 **4. Results**

### 308 4.1 Identification of upper and lower layer currents.

309 The central tropical Atlantic has a very characteristic seasonal variability of the surface  
310 dynamic height. During December through May,  $DH_0$  is nearly flat between  $10^\circ\text{S}$ - $10^\circ\text{N}$ ,  
311 implying small surface geostrophic velocities (Figure 6a, b). From June to November (Figure 6c,  
312 d), the  $DH_0$  gradients are larger north of  $5^\circ\text{N}$  and as a result the NECC core is stronger and well  
313 defined. A second NECC core develops in late boreal summer/early boreal fall around  $8^\circ$ - $10^\circ\text{N}$   
314 (Figure 6c, d) and the NECC shows a maximum double peak structure during boreal fall (Figure  
315 6d) which is in agreement with previous results [e.g., Didden and Schott, 1992; Urbano et al.,  
316 2006]. Results obtained indicate that there is often a connection between the surface NECC and  
317 the flow in the thermocline layer, when the NEUC can be found as a lobe attached to the

318 southern NECC branch. During boreal spring (Figure 6b), however, the NEUC is clearly  
319 detached from the NECC. A signature of the NEUC on the  $\sigma_\theta = 26.8 \text{ kg m}^{-3}$  isopycnal is  
320 observed at approximately  $5^\circ\text{N}$ , but it does not show a very sharp gradient [c.f., Bourlès et al.,  
321 2002; Schott et al., 2003]. In contrast, the potential density structure at about  $4^\circ\text{-}5^\circ\text{S}$ , where the  
322 SEUC is located at around 200 m, exhibits a very distinguished southward elevation of the  $\sigma_\theta =$   
323  $26.8 \text{ kg m}^{-3}$  surface. This is an indication that the meridional pressure gradient plays an important  
324 role in SEUC dynamics. The isopycnal signature of the SEUC and its well-defined core are  
325 visible throughout the year, indicating that the SEUC is a permanent feature in the tropical  
326 Atlantic. The EUC core is generally located south of the equator and between the surface and  
327 100 m, with eastward flow typically between  $2.5^\circ\text{S}$  and  $2.5^\circ\text{N}$ , (Figure 6).

328         We define the location of the equatorial eastward currents by assigning a latitudinal band  
329 to each of them (Table 1), similar to the bands used in previous studies [e.g., Hüttl-Kabus and  
330 Böning, 2008]: The NECC is defined between  $3^\circ\text{-}10^\circ\text{N}$ , the NEUC between  $3^\circ\text{-}6^\circ\text{N}$ , the SEUC  
331 between  $3^\circ\text{-}6^\circ\text{S}$ , and the EUC between  $2.5^\circ\text{S}$  and  $2.5^\circ\text{N}$ . We further characterize these currents  
332 by selecting isopycnal layers to define their vertical boundaries: an upper or surface layer from  
333 the surface to  $\sigma_\theta = 24.5 \text{ kg m}^{-3}$ , and a lower or thermocline layer between  $\sigma_\theta = 24.5 - 26.8 \text{ kg m}^{-3}$   
334 [e.g., Schott et al., 1998]. The NECC is restricted to the upper layer, and the OEUCs to the lower  
335 layer. The EUC has significant contributions in both layers, thus it is considered to occupy both  
336 the surface and lower layers.

337

#### 338 4.2 Synthetic method validation

339         We validate the synthetic method by comparing the transports of the eastward equatorial  
340 currents estimated from the 31 XBT transect realizations with their synthetic counterparts

341 (Figure 7). The agreement between the synthetic and XBT transport estimates is high for the  
342 NECC (Figure 7a, b), with  $RMSE = 0.72$  Sv. A strong linear relationship is found between the  
343 two estimates, with a corresponding linear-correlation coefficient of  $R^2 = 0.95$ . The EUC (Figure  
344 7g, h), which has transport contributions from the upper and lower layers, shows a moderate-to-  
345 high agreement between the synthetic and XBT transport estimates ( $R^2 = 0.55$ ,  $RMSE = 7.0$  Sv).  
346 Regarding the OEUCs, the synthetic method has considerable skill for the NEUC ( $R^2 = 0.34$ ), but  
347 the transports using XBT data are generally higher than the synthetic ones (Figure 7c, d), with  
348  $RMSE = 1.2$  Sv. The synthetic method is able to reproduce a comparable mean SEUC transport  
349 of 7 Sv to the one derived from the XBT data. However, the method fails to reproduce the SEUC  
350 transport variations (Figure 7e), with low correlations between the synthetically-derived and  
351 XBT-derived SEUC transports ( $R^2 = 0.04$ ,  $RMSE = 0.43$  Sv), and the variance of the  
352 synthetically-derived SEUC transport barely resembles the ones from hydrography (Figure 7e).  
353 Due to the low agreement between the two SEUC transport estimates, analysis of the variability  
354 of the synthetic fields will focus on the NECC, NEUC, and to a lesser extent the EUC.

355

#### 356 4.3 Seasonal variability of the eastward equatorial currents

357 Here we analyze the seasonal variability of the NECC, NEUC, SEUC, and EUC along  
358 the AX08 transect in terms of cross-sectional volume transport, core velocity, and core latitude  
359 using their monthly average values. The selected 31 AX08 temperature sections provide high-  
360 resolution coverage of the central tropical Atlantic for all months except February (Figure 1c).  
361 We apply to the XBT-derived quantities a previously used methodology [Foltz et al., 2004a] to  
362 estimate monthly means and variances of transport, core velocity, and position of the currents.  
363 This method fits the first two harmonics (i.e., annual and semi-annual) to the data in a least



364 squares fashion. It accounts for observational uncertainties by calculating a diagonal covariance  
365 matrix, with each diagonal term corresponding to a month, calculated as the sum of the fitting  
366 residual error and the observational error given by the standard deviation of each month. Fitting  
367 the first two annual harmonics to the data filters a large fraction of the mesoscale signal.

368

#### 369 *a) NECC*

370 The hydrographic estimates show a strong annual cycle of the NECC, which alone  
371 represents 58% of the transport, 41% of the maximum velocity, and 52% of the position  
372 variances (Table 2). Earlier analyses based on drifter data also show a dominant contribution, as  
373 high as 80%, of the annual cycle for the off-equatorial surface currents [e.g., Richardson and  
374 Walsh, 1986; Lumpkin and Garzoli, 2005]. The NECC (Figure 8a, i) reaches its southernmost  
375 position in boreal spring ( $\sim 3^{\circ}$ - $4^{\circ}$ N), concurrent with its lowest transport ( $\sim 1$  Sv), and its  
376 northernmost position ( $\sim 7^{\circ}$ - $9^{\circ}$ N) during boreal fall, coincident with its highest transport ( $\sim 10$ - $11$   
377 Sv). The NECC core velocity shows a similar pattern, with lowest values ( $\sim 10$ - $20$  cm/s) in boreal  
378 spring and maximum values ( $\sim 55$  cm/s) in boreal fall (Figure 8e). This seasonal variability of the  
379 NECC agrees with previous observational findings [e.g., Richardson et al., 1992], and it is linked  
380 to the north-south migration of the ITCZ [Katz and Garzoli, 1982; Katz, 1987; Garzoli and  
381 Richardson, 1989; Garzoli, 1992; Fonseca et al., 2004], which is near the equator during boreal  
382 spring and farthest north during boreal summer.

383 The estimates using the synthetic method for the comparable 2000-2010 period (blue  
384 lines in Figure 8a, e, i) follow closely the estimates from hydrography. As the synthetic estimates  
385 are averaged using a much longer time series, and the hydrographic data are subject to higher  
386 influence of mesoscale effects due to sparse temporal sampling (Figure 1c), differences are

387 expected and the synthetic method generally shows higher explained variances by the annual  
388 cycle (Table 2). The synthetically-derived core velocity time series (Figure 8e) reveal stronger  
389 semi-annual variability during 1993-2000 (blue lines) compared to the 2000-2010 decade (green  
390 lines) and hydrographic data (red lines). Additional characteristics of how the NECC transport  
391 varies over time can be retrieved by applying a wavelet transform [Torrence and Compo, 1998]  
392 to the NECC transport time series for the whole altimetric period (Figure 9a). This analysis  
393 confirms that the annual cycle of the NECC is the strongest signal in the synthetically-derived  
394 transport time series. Moreover, the energy of the annual period shows intermittent increases and  
395 a prominent semi-annual signal is observed during an event in 1998-1999, when there is a strong  
396 peak in the NECC transport time series. The latter might be related to a strong wind stress curl  
397 anomaly in the western tropical Atlantic during that period, which has been previously described  
398 [Fonseca et al., 2004]. Mechanisms related to the interannual variability of this current are  
399 explored in Section 4.2.

400

#### 401 *b) NEUC*

402 The NEUC shows higher transport values during boreal spring (Figure 8b), when its  
403 transport ranges from 5-10 Sv, and lower transports (~2-3 Sv) in boreal fall. Compared to the  
404 seasonal cycle of NECC transport (Figure 8a), the NEUC bears an almost inverse relationship  
405 with the NECC, and therefore with the ITCZ. The overlapping of the positions of these two  
406 currents from April to May (Figures 8i, j) suggests that the NECC may partly contribute to the  
407 NEUC flow during this period. The annual cycle dominates the XBT transport variability  
408 (explained variance: 63%), whereas the semi-annual cycle only explains 1% of the variance  
409 (Table 2). The seasonal variability of NEUC location (Figure 8j) resembles its transport

410 variability, in that a more southward position ( $\sim 3.5^\circ$ - $4^\circ$ N) occurs during boreal spring at the time  
411 of higher transport. The NEUC core velocity (Figure 8f) exhibits a relatively strong semi-annual  
412 cycle (explained variance: 35%, Table 2), with maximum speed of about 25 cm/s during boreal  
413 winter and summer and smaller core velocity of about 10-15 cm/s in September-October,  
414 revealing that the NEUC is present year-round in this latitude band. The synthetic estimates  
415 broadly agree in magnitude with the XBT measurements and may indicate some changes  
416 between the 1990s and 2000s. The wavelet analysis performed for the whole altimetric period  
417 indicates a possible shift in energy from a prevailing semi-annual to an annual variability starting  
418 in 1999-2000 (Figure 9b). There are also indications of interannual variability in the NEUC  
419 transport, with increased energy centered in 1995, 2001 and 2006. The resemblance between the  
420 variability of the zonal winds in the equatorial region (Figure 9d) and the NEUC transport  
421 variability is noteworthy, and mechanisms for this variability will be explored in Section 4.2.

422

### 423 *c) SEUC*

424 The XBT-derived transport of the SEUC (Figure 8c) is weaker ( $\sim 6$  Sv) from July to  
425 September, in agreement with Reverdin et al. [1991], and it is also reduced in January ( $\sim 5$  Sv).  
426 Maximum transport values are found from October through December, reaching up to 10 Sv, and  
427 mean transport is about 7-8 Sv, in good agreement with previous estimates [Brandt et al., 2006;  
428 Fischer et al., 2008]. The semi-annual harmonic component is dominant and explained variances  
429 are 90% and 92% for transport and core velocity (Table 2), respectively, which is consistent with  
430 model studies [e.g., Hüttl-Kabus and Böning, 2008]. The core velocity is about 28 cm/s in  
431 May/June and October, and 20 cm/s in August (Figure 8g). Fischer et al. [2008] reported a  
432 weaker mean SEUC velocity of 13.4-17.9 cm/s in central tropical Atlantic, but departures from

433 this mean of up to 50 cm/s have been observed in its core. The mean position of the SEUC is  
434 about  $4.5^{\circ}\pm 0.5^{\circ}\text{S}$  and although small, its position variability (Figure 8k) follows closely its  
435 transport variability, in that higher transports are associated with a more northern position and  
436 vice versa.

437

#### 438 *d) EUC*

439 The XBT and synthetic estimates agree well in that the EUC near  $23^{\circ}\text{W}$  is located mostly  
440 south of the equator, between  $1^{\circ}\text{S}\pm 1^{\circ}$  (Figure 8l). Despite the relatively small seasonal variability  
441 of the EUC location, there are indications of a more southward position during boreal fall/winter.  
442 The core velocity (Figure 8h) derived from XBT data is generally higher than from the synthetic  
443 method, between 40-80 cm/s, and shows strong annual variability (explained variance: 72%,  
444 Table 2), whereas the synthetic core velocity estimates range from 30 to 50 cm/s, and exhibit a  
445 more semi-annual periodicity. The XBT estimates are comparable to direct velocity observations  
446 of the EUC, which show mean values of about 70-80 cm/s [e.g., Giarolla et al., 2005; Brandt et  
447 al., 2006, 2008]. The EUC transports (Figure 8d) derived from both hydrographic and synthetic  
448 estimates exhibit strong semi-annual variability (i.e., explained variances between 40-68%, Table  
449 2), with lower values during boreal summer and winter, and higher values during boreal spring  
450 and late summer-early fall. The wavelet analysis (Figure 9c) of the synthetic transport confirms  
451 the strong semi-annual variability of the EUC, in agreement with results from previous model  
452 studies of the central equatorial Atlantic [e.g., Arhan et al., 2006; Hormann and Brandt, 2007].  
453 The high variability of the XBT estimates in some months (see for instance the error bars of  
454 transport and core velocity during July in Figures 8d, h) are an artifact of the equatorial beta  
455 approximation for specific monthly density structures and need to be interpreted carefully.

456

#### 457 4.4 Interannual variability of the synthetic transports of the eastward currents

458           In this section, we investigate the interannual signature of the NECC and NEUC on SST  
459 and surface wind stress using statistical tools. We restrict our analysis to the NECC and NEUC  
460 transports because the synthetic method does not provide good estimates of the SEUC variability  
461 and because we do not expect that the equatorial beta approximation can produce reliable  
462 estimates of the EUC transport variability on interannual timescales. We use here monthly  
463 transport anomalies relative to the monthly climatology, with the intent of reducing the influence  
464 of eddy variability on these currents. Furthermore, we perform correlation analyses of the  
465 monthly NECC and NEUC transport anomalies with the gridded monthly anomalies of both SST  
466 and pseudo-wind stress. Correlation analyses are performed in the tropical Atlantic region for the  
467 1993-2010 period, taking only statistically significant values into account ( $p < 0.05$ ). In this  
468 analysis, all data are standardized by subtracting their mean and dividing by their standard  
469 deviation and low-pass filtered with a 5-month moving average.

470

##### 471 *a) NECC*

472           The correlation between the NECC transport and SST anomalies (SSTA) for the  
473 altimetric period produces a distinct pattern in the form of an anomalous interhemispheric SST  
474 gradient (Figure 10a), with positive phase just north of the equator in the ITCZ region, centered  
475 at approximately  $2^{\circ}\text{N}$ ,  $30^{\circ}\text{W}$ , and extending northeastward, and negative phase in the central  
476 south of the domain, centered at about  $20^{\circ}\text{S}$ ,  $17^{\circ}\text{W}$ . The corresponding correlation with pseudo-  
477 wind stress anomalies indicates an anomalous strengthening of the southeasterly trades, with  
478 largest magnitude in the western equatorial region where winds are also strongest. This pattern is

479 reminiscent of the Atlantic meridional mode [e.g., Chang et al., 2006, and references therein],  
480 which is believed to be driven by the wind-evaporation-SST (WES) feedback mechanism  
481 involving interactions between SST changes and wind-induced latent heat fluxes [Xie and  
482 Philander, 1994; Chang et al., 1997]. A similar pattern has also been found for the NECC in a  
483 recent study using complex empirical-orthogonal-function analysis of surface drifter data  
484 [Hormann et al., 2012].

485         A SSTA gradient index can be defined by subtracting area averages within the northern  
486 ( $35^{\circ}\text{W}$ - $17^{\circ}\text{W}$ / $0^{\circ}$ - $7^{\circ}\text{N}$ ) and southern ( $30^{\circ}\text{W}$ - $15^{\circ}\text{W}$ / $11^{\circ}\text{S}$ - $23^{\circ}\text{S}$ ) boxes marked in Figure 10a. The  
487 correlation between the monthly NECC transport anomalies and the SSTA index is significant ( $R$   
488 = 0.43), with maximum correlation at zero lag. A meridional wind-stress (y-wind) index,  
489 computed in the northern box described above, which focuses on the variability near the central  
490 equatorial region, shows a maximum positive correlation of  $R = 0.44$ , with the NECC lagging  
491 the wind strengthening by one month. This relationship indicates that there is a fast response of  
492 the NECC and SST to wind anomalies that might be explained by either the fast adjustment time  
493 of the ocean through equatorial waves [e.g., Ma, 1996] or by anomalies simply connected to a  
494 strengthening of the surface NBC retroflection.

495

#### 496 *b) NEUC*

497         The pattern arising from the correlation between the NEUC transport with SST and wind  
498 stress anomalies (Figure 11) can be described as follows: Negative SST coefficients prevail in  
499 the northeastern part of the basin, while positive correlations are found along the equator and off  
500 the southwestern African coast. In addition, the trades are reduced in the western to central  
501 equatorial Atlantic and southwesterly wind anomalies prevail over the Gulf of Guinea. The

502 obtained interhemispheric SST pattern suggests a relation to the Atlantic meridional mode, and  
503 the equatorial SST pattern also indicates a relation to the zonal mode [e.g., Chang et al., 2006,  
504 and references therein], further supporting the previously proposed link between these two  
505 modes [Servain et al., 1999; Foltz and McPhaden, 2010]. The wind stress anomalies also agree  
506 with proposed mechanisms in that this anomalous SST pattern shifts the ITCZ south [Moura and  
507 Shukla, 1981], reinforcing the positive SST anomaly along equator [Foltz and McPhaden, 2010].

508 Previous studies that used virtual Lagrangian floats have shown that the NEUC provides  
509 waters for the upwelling in the Guinea Dome and equatorial regions [Stramma et al., 2005;  
510 Hüttl-Kabus and Böning, 2008]. Results obtained here agree that the NEUC transport anomalies  
511 are negatively correlated with SSTAs and northeasterly wind anomalies in the Guinea Dome  
512 region (Figure 11a), which are consistent with an increased coastal upwelling in this region. As  
513 for the NECC, we create an interhemispheric SST index and a meridional wind stress index, and  
514 relate these indices to the NEUC transport. We consider here the difference between SSTA in the  
515 Guinea Dome region ( $10^{\circ}$ - $25^{\circ}$ N/ $15^{\circ}$ - $35^{\circ}$ W) and in the southeast Atlantic ( $8^{\circ}$ S- $5^{\circ}$ N/ $30^{\circ}$ W- $10^{\circ}$ E)  
516 as well as the meridional wind stress average over the Guinea Dome region. We find that the  
517 maximum correlation between the NEUC transport with the indices of SST and wind to be  $R = -$   
518  $0.51$  and  $R = -0.32$ , respectively, both at zero lag (Figure 11c).

519

## 520 **5. Discussion and Conclusions**

521 In the present study, we have used a combination of high-density quarterly AX08 XBT  
522 transect and altimetric data to investigate the variability of the eastward surface and subsurface  
523 currents in the equatorial Atlantic. The high spatial resolution of the hydrographic data and its  
524 repeated sampling of the region enabled us to assess the dynamic and thermodynamic properties

525 of the upper ocean, and permit the characterization of the seasonal cycle of major geostrophic  
526 zonal currents such as the NECC and NEUC. However, due to strong regional intraseasonal  
527 variability generated by eddy effects and the passage of TIWs, it is hypothesized that it would be  
528 necessary to carry out a high number of sections for one to be able to produce a good statistical  
529 estimate of the seasonal cycle of these currents using hydrography only. For instance, out of the  
530 31 analyzed XBT sections only one realization is available for January, May, and June, and no  
531 section is yet available for February.

532         Combining altimetry and XBT data by a synthetic methodology reproduces the main  
533 surface and subsurface features along the equatorial AX08 XBT transect, which could not be  
534 obtained by altimetry alone, overcoming the sampling restrictions, and producing better  
535 estimates of the long-term average of the currents properties. The best performance of the  
536 synthetic methodology for the equatorial AX08 region resides in the upper 200-300 m of the  
537 water column, where temporal correlations between SLA and anomalies of density and dynamic  
538 height are larger than 0.8.

539         The seasonal cycle of the eastward equatorial currents derived from our analyses are in  
540 good agreement with previous works [e.g., Richardson and Reverdin, 1987; Peterson and  
541 Stramma, 1991; Bourlès et al., 2002; Hormann and Brandt, 2007]. The NECC exhibits a strong  
542 seasonal cycle, with transports ranging from 1 to 11 Sv and maximum speed greater than 50  
543 cm/s. Its seasonal variability is related to the migration of the ITCZ [e.g., Garzoli and  
544 Richardson, 1989; Fonseca et al., 2004], with stronger transport positively correlated to a  
545 northward shift of the ITCZ. On interannual timescales, the NECC transport is found to be linked  
546 to a strengthening of the southwesterly trades and a positive interhemispheric SST gradient  
547 pattern (i.e., warmer tropical North Atlantic and colder tropical South Atlantic), consistent with



548 Hormann et al. [2012]. The NECC strengthening associated with such an interhemispheric SST  
549 gradient might act as a positive feedback, since it would increase the eastward transport of  
550 warmer western waters toward the region of the SST maximum gradient. Previous studies have  
551 shown that this tropical Atlantic SST pattern is influenced by the Atlantic Multidecadal  
552 Oscillation and the North Atlantic Oscillation, and teleconnections from the eastern Pacific [e.g.,  
553 Enfield and Mayer, 1997; Czaja et al., 2002]. The recent decrease in sulphate aerosol emissions  
554 over the North Atlantic might be associated with a SST increase there and a consequent  
555 northward trend of the ITCZ [Chang et al., 2011]. Our results indicate a possible recent  
556 strengthening of the climatological NECC of the order of about 1 Sv in the 2000s (Figure 8a),  
557 but this value is small compared to the year-to-year NECC variability.

558         The dominant annual cycle of the NEUC is characterized by stronger transports from  
559 January to July (up to 10 Sv), which is in opposite phase to the NECC transport cycle. The  
560 NEUC core is located between 4°N-6°N, with maximum velocities of about 30 cm/s in June-  
561 July. Synthetic estimates of the NEUC transport suggest that the semi-annual variability used to  
562 be stronger during the 1990s, that there was a shift to a stronger annual periodicity since 2000,  
563 and that throughout the record there has been an interannual modulation of the annual variability.  
564 The latter is also in agreement with anomalous zonal wind stress variability in the central-  
565 western equatorial Atlantic (Figure 9b, d) as well as an index of the Atlantic meridional mode  
566 [Foltz et al., 2012]. Our results further indicate that the interannual variability of the NEUC  
567 transport is statistically related to the upwelling in the equatorial Atlantic and the variability in  
568 the Guinea Dome region. Such a link between the NEUC and the Guinea Dome has long been  
569 proposed [e.g., Voituriez, 1981; Schott et al., 2004], since the uplifting of the thermal structure in  
570 the dome extends much further down than the thermocline. The strong upwelling in this region

571 is, for instance, related to the outcropping of the  $\sigma_\theta = 24.5 \text{ kg/m}^3$  isopycnal. The underlying  
572 mechanism of this relationship might be that a cooler Guinea Dome and warmer equatorial  
573 region increase the north-south density gradient in the NEUC region, and strengthen its core.  
574 Some model studies in the Pacific indeed suggest that the coastal upwelling in the eastern side of  
575 the basin can drive the variability of such a current [e.g., McCreary et al., 2002].

576 The NEUC and NECC exhibit an inverse relationship in transport variability on  
577 interannual timescales. Both are linked to the variability of the cross-equatorial wind stress, but  
578 the NECC is strengthened in association with increased southwesterly winds whereas increases  
579 in the NEUC are associated with reduced southwesterly winds. The interannual modes of tropical  
580 Atlantic variability are strongly tied to their seasonal cycle [Nobre and Shukla, 1996, Servain et  
581 al., 1999], which can also explain why the inverse relationship between the NEUC and NECC  
582 transports also holds for seasonal timescales. Our results suggest that the NECC contributes to  
583 the NEUC, especially during the boreal spring, when the NECC is located further south.

584 The SEUC calculated from hydrography is located on average at about  $4.5^\circ\text{S}$  and exhibits  
585 a rather weak seasonal cycle, with mean transport of about 7 Sv, higher values from October to  
586 December, and a secondary peak during May-June. Velocities during the stronger months are of  
587 the order of 25-30 cm/s. Although the synthetic method produced a comparable SEUC mean  
588 flow, it failed to describe its variability. The limited ability of our methodology to capture the  
589 SEUC is a consequence of the somewhat weak variability of SLA in the southern tropical  
590 Atlantic (Figure 1a), and the resulting surface signature of the current is masked by  
591 compensating effects in the water column (Figure 6). The compensating effects can be explained,  
592 for example, by buoyancy and wind forcing components of the same magnitude and opposing

593 signs [Mayer et al., 2001]. Therefore, regular hydrographic sampling is particularly important for  
594 monitoring the SEUC.

595         The EUC can be resolved by the equatorial beta approximation using hydrographic data.  
596 The seasonal cycle of the EUC is well explained by the annual and semi-annual harmonics, with  
597 peaks in boreal fall and spring indicating a relatively strong semi-annual component. The mean  
598 EUC derived from XBT data transports 24 Sv, with a corresponding core velocity of 56 cm/s,  
599 whereas the synthetic EUC exhibits a mean transport of 14 Sv and core velocity of 34 cm/s.  
600 Based on 11 cross-equatorial ship sections taken in the central equatorial Atlantic, Brandt et al.  
601 [2006] reported a mean EUC transport of 13.8 Sv, and observed mean velocities are of the order  
602 of 70-80 cm/s [e.g., Giarolla et al., 2005; Brandt et al., 2006, 2008]. Therefore, we conclude that  
603 the modified Lagerloef et al. [1999] equatorial beta approximation applied to dynamic height  
604 estimates generates a weaker, more diffuse, and wider EUC compared to direct velocity  
605 observations.

606         Results from this study are subject to specific caveats that provide avenues for future  
607 research: First, we use a simple statistical method to infer the relationship between surface height  
608 and ocean properties at depth. Using an improved statistical method may allow for including  
609 additional information, such as latitudinal cross-correlation between and autocorrelation of the  
610 residuals at depth, as well as the use of additional constraints derived from co-located  
611 observations. Second, we use climatological values of salinity calculated from a mean T-S  
612 relationship at each location, and also climatological values of the absolute dynamic height at  
613 800 m. Available observations from the Argo network, for example, could reduce errors in the  
614 methodology. However, these observations are mostly restricted to the last 5-6 years. Third,  
615 along-track altimetry data might provide better agreement between altimetric SLA and

616 hydrography-based dynamic height at the surface, since it is less smoothed by optimal  
617 interpolation procedures than the gridded altimetry product used here. Finally, high-resolution  
618 modeling can fill gaps in the observational results and confirm the robustness of our conclusions  
619 in a dynamically consistent fashion.

620

621

622 **Acknowledgements**

623           The authors would like to thank Greg Foltz, Rick Lumpkin, and Molly Baringer for  
624 insightful discussions, and George Halliwell for helpful comments. This research was carried out  
625 under the auspices of the Cooperative Institute for Marine and Atmospheric Studies (CIMAS), a  
626 cooperative institute of the University of Miami and the National Oceanic and Atmospheric  
627 Administration (NOAA), cooperative agreement #NA17RJ1226. The XBT transects are funded  
628 by NOAA's Climate Program Office and the authors acknowledge the effort of ship crews and  
629 logistical personnel.

630

631

632 **References**

633

634 Arhan, M., A. M. Treguier, B. Boulès, and S. Michel (2006), Diagnosing the annual cycle of the  
635 Equatorial Undercurrent in the Atlantic ocean from a general circulation model, *J. Phys.*  
636 *Oceanogr.*, 36, 1502-1522.

637 Athie, G., and F. Marin (2008), Cross-equatorial structure and temporal modulation of  
638 intraseasonal variability at the surface of the Tropical Atlantic ocean, *J. Geophys. Res.*, 113,  
639 C08020, doi:10.1029/2007JC004332.

640 Atlas, R., R. N. Hoffman, J. Ardizzone, S. M. Leidner, J. C. Jusem, D. K. Smith, and D. Gombos  
641 (2011), A cross-calibrated, multiplatform ocean surface wind velocity product for  
642 meteorological and oceanographic applications, *Bull. Amer. Meteor. Soc.*, 92, 157-174. doi:  
643 10.1175/2010BAMS2946.1.

644 Boulès, B., Y. Gouriou, and R. Chuchla (1999), On the circulation in the upper layer of the  
645 western equatorial Atlantic, *J. Geophys. Res.*, 104, C9, 21,151-21,170.

646 Boulès, B., M. D'Orgeville, G. Eldin, Y. Gouriou, R. Chuchla, Y. DuPenhoat, and S. Arnault  
647 (2002), On the evolution of the thermocline and subthermocline eastward currents in the  
648 Equatorial Atlantic, *Geophys. Res. Lett.*, 29(16), 1785, doi: 10.1029/2002GL015098.

649 Brandt, P., F. A. Schott, C. Provost, A. Kartavtseff, V. Hormann, B. Boulès, and J. Fischer  
650 (2006), Circulation in the central equatorial Atlantic: Mean and intraseasonal to seasonal  
651 variability, *Geophys. Res. Lett.*, 33, L07609, doi:10.1029/2005GL025498.

652 Brandt, P., V. Hormann, B. Boulès, J. Fischer, F. A. Schott, L. Stramma, and M. Dengler  
653 (2008), Oxygen tongues and zonal currents in the equatorial Atlantic, *J. Geophys. Res.*, 113,  
654 C04012, doi:10.1029/2007JC004435.

655 Brandt, P., V. Hormann, A. Körtzinger, M. Visbeck, G. Krahnemann, L. Stramma, R. Lumpkin,  
656 and C. Schmid(2010), Changes in the ventilation of the oxygen minimum zone of the tropical  
657 North Atlantic, *J. Phys. Oceanogr.*, 40, 1784-1801, doi:10.1175/2010JPO4301.1.

658 Brandt, P., A. Funk, V. Hormann, M. Dengler, R. J. Greatbatch, and J. M. Toole (2011),  
659 Interannual atmospheric variability forced by the deep equatorial Atlantic Ocean, *Nature*,  
660 473, 497–500, doi:10.1038/nature10013.

661 Carton, J. A., and B. H. Huang (1994), Warm events in the tropical Atlantic. *J. Phys. Oceanogr.*,  
662 24, 888-903.

663 Carton, J. A., X. H. Cao, B. S. Giese, and A. M. daSilva (1996), Decadal and interannual SST  
664 variability in the tropical Atlantic Ocean, *J. Phys. Oceanogr.*, 26, 1165-1175.

665 Chang, P., L. Ji, and H. Li (1997), A decadal climate variation in the tropical Atlantic ocean from  
666 thermodynamic air-sea interactions, *Nature*, 385, 516–518.

667 Chang, P., T. Yamagata, P. Schopf, S. K. Behera, J. Carton, W. S. Kessler, G. Meyers, T. Qu, F.  
668 Schott, S. Shetye, S. P. Xie (2006), Climate Fluctuations of Tropical Coupled Systems—The  
669 Role of Ocean Dynamics, *J. Clim.*, 19, 20, 2006, 5122-5174, doi:10.1175/JCLI3903.1

670 Chang, P., Y. Fang, R. Saravanan, L. Ji, and H. Seidel (2006), The cause of the fragile  
671 relationship between the Pacific Nino and the Atlantic Nino, *Nature*, 443, 324-328, doi:  
672 10.1038/nature05053.

673 Chang, C-Y, J. C. H. Chang, M. F. Wehner, A. R. Friedman and R. Ruedy (2011), Sulfate  
674 Aerosol Control of Tropical Atlantic Climate over the Twentieth Century, *J. Clim.*, 24, 2540-  
675 2555, Doi: 10.1175/2010JCLI4065.1.

676 Cheney, R., L. Miller, R. Agreen, N. Doyle, and J. Lillibridge (1994), TOPEX/Poseidon: The 2-  
677 cm solution, *J. Geophys. Res.*, 99, 24,555-24,563.

678 Cochrane, J. D., F. J. Kelly Jr., and C. R. Olling (1979), Subthermocline counter-currents in the  
679 western equatorial Atlantic Ocean, *J. Phys. Oceanogr.*, 9, 724–738.

680 Conkright, M. E., et al. (2002), World Ocean Database 2001, vol. 1, Introduction, NOAA Atlas  
681 NESDIS 42, 159 pp., U.S. Gov. Print. Off., Washington D. C.

682 Czaja, A., P. van der Vaart and J. Marshall (2002), A diagnostic study of the role of remote  
683 forcing in Tropical Atlantic variability, *J. of Climate*, vol 15, 3280-3290, 2002.

684 Didden, N., and F. Schott (1992), Seasonal variations in the Western Tropical Atlantic: surface  
685 circulation from geosat altimetry and WOCE model results, *J. Geophys. Res. C* 97 (3), 3529–  
686 3541.

687 Doi, T., T. Tozuka, H. Sasaki, Y. Masumoto, and T. Yamagata (2007), Seasonal and interannual  
688 variations of oceanic conditions in the Angola Dome, *J. Phys. Oceanogr.*, 37(11), 2698–  
689 2713.

690 Ducet, N., P.-Y. Le Traon, and G. Reverdin (2000), Global high resolution mapping of ocean  
691 circulation from Topex/Poseidon and ERS-1 and -2, *J. Geophys. Res.*, 105 (C8), 19477-  
692 19498.

693 Düing, W., P. Hisard, E. Katz, J. Meincke, L. Miller, K. V. Moroshkin, G. Philander, A. A.  
694 Ribnikov, K. Voigt, and R. Weisberg (1975), Meanders and long waves in the equatorial  
695 Atlantic, *Nature*, 257, 280-284, doi:10.1038/257280a0.

696 Enfield, D. B., and D. A. Mayer (1997), Tropical Atlantic sea surface temperature variability and  
697 its relation to El Niño-Southern Oscillation, *J. Geophys. Res.*, 102(C1), 929–945,  
698 doi:10.1029/96JC03296.



699 Fischer, J., V. Hormann, P. Brandt, F. A. Schott, B. Rabe, and A. Funk (2008), South Equatorial  
700 Undercurrent in the western to central tropical Atlantic, *Geophys. Res. Lett.*, 35, L21601,  
701 doi: 10.1029/2008GL035753.

702 Fofonoff, N. P., and C. Froese Fisher, (1958) Program for oceanographic computations and data  
703 processing on the electronic digital computer ALWAC II-E, PSW-1 Programs for properties  
704 of sea water. Manuscript Rep. Series, *Limnol. Oceanogr.*, J. Fish. Res. Bd. Canada, 27, 39  
705 pp.

706 Foltz, G. R., S. A. Grodsky, and J. A. Carton (2003), Seasonal mixed layer heat budget of the  
707 tropical Atlantic Ocean. *J. Geophys. Res.*, 108, 3146, doi:10.1029/2002JC001584.

708 Foltz, G., S. Grodsky, and J. Carton (2004a), Seasonal salt budget of the northern tropical  
709 Atlantic Ocean along 38°W, *J. Geophys. Res.*, 109, C03052, doi: 10.1029/2003JC002111.

710 Foltz, G., J. Carton, and E. P. Chassignet (2004b), Tropical instability vortices in the Atlantic  
711 Ocean, *J. Geophys. Res.*, 109, C03029, doi: 10.1029/2003JC001942.

712 Foltz, G. R., and M. J. McPhaden (2010), Interaction between the Atlantic meridional and Niño  
713 modes, *Geophys. Res. Lett.*, 37, L18604, doi:10.1029/2010GL044001.

714 Foltz, G.R., M.J. McPhaden, and R. Lumpkin (2012), A strong Atlantic Meridional Mode event  
715 in 2009: The role of mixed layer dynamics, *J. Clim.*, 25, 363-380, doi:10.1175/JCLI-D-11-  
716 00150.1.

717 Fonseca, C., G. Goni, W. E. Johns, and E. J. D. Campos (2004), Investigations of the North  
718 Brazil Current retroflexion and North Equatorial Countercurrent variability, *Geophys. Res.*  
719 *Lett.*, 31, L21304, doi: 10.1029/2004GL020054.

720 Frantantoni, D., W. E. Johns, T. L. Townsend, and H. E. Hurlburt (2000), Low latitude  
721 circulation and mass transport pathways in a model of the tropical Atlantic Ocean. *J. Phys.*  
722 *Oceanogr.*, 30, 1944–1966.

723 Garzoli, S. L., and P. L. Richardson (1989), Low-frequency meandering of the Atlantic North  
724 Equatorial Countercurrent, *J. Geophys. Res.-Oceans*, 94(C2), 2079-2090.

725 Garzoli, S. L. (1992), The Atlantic Equatorial Countercurrent: Models and observations. *J.*  
726 *Geophys. Res.*, 97(C11),17,931-17,946.

727 Garzoli, S. L., and E. J. Katz (1983), The forced annual reversal of the Atlantic North Equatorial  
728 Countercurrent. *J. Phys. Oceanogr.*, 13, 2082-2090.

729 Giarolla, E., P. Nobre, M. Malagutti, and L. P. Pezzi (2005), The Atlantic Equatorial  
730 Undercurrent: PIRATA observations and simulations with GFDL Modular Ocean Model at  
731 CPTEC, *Geophys. Res. Lett.*, 32, L10617, doi:10.1029/2004GL022206.

732 Gilson, J., D. Roemmich, B. Cornuelle, and L.-L. Fu (1998), Relationship of TOPEX/Poseidon  
733 altimetric height to steric height and circulation in the North Pacific, *J. Geophys. Res.*,  
734 103(C12), 27,947–27,965, doi:10.1029/98JC01680.

735 Goes, M., and I. Wainer (2003), Equatorial Currents transport changes for extreme warm and  
736 cold events in the Atlantic Ocean, *Geophys. Res. Lett.*, 30 (5), doi: 10.1029/2002GL015707.

737 Goes, M., R. Molinari, I. C. A. Silveira, and I. Wainer (2005), Retroreflections of the North Brazil  
738 Current during February 2002, *Deep Sea Res. Part 1*, 52, 3, 647-667, doi:  
739 10.1016/j.dsr.2004.10.010.

740 Goldenberg, S. B., C. Landsea, A. M. Mestas-Nunez, and W. M. Gray (2001), The recent  
741 increase in Atlantic hurricane activity, *Science*, 293, 474– 479.

742 Goni, G., and Baringer, M. (2002), Surface currents in the tropical Atlantic across high density  
743 XBT line AX08, *Geophys. Res. Lett.* 29(24), doi: 10.1029/2002GL015873.

744 Grodsky, S. A., J. A. Carton, C. Provost, J. Servain, J. A. Lorenzetti, and M. J. McPhaden  
745 (2005), Tropical instability waves at 0°N, 23°W in the Atlantic: A case study using Pilot  
746 Research Moored Array in the Tropical Atlantic (PIRATA) mooring data, *J. Geophys. Res.*,  
747 110, C08010, doi:10.1029/2005JC002941. Guinehut, S., P.-Y. Le Traon, and G. Larnicol  
748 (2006), What can we learn from Global Altimetry/Hydrography comparisons?, *Geophys.*  
749 *Res. Lett.*, 33, L10604.

750 Hanawa, K., P. Raul, R. Bailey, A. Sy, and M. Szabados (1995), A new depth-time equation for  
751 Sippican or TSK T-7, T-6 and T-4 expendable bathythermographs (XBT), *Deep-Sea Res. I*,  
752 42, 1423-1451.

753 Hazeleger, W., P. de Vries, and Y. Friocourt (2003), Sources of the Equatorial Undercurrent in  
754 the Atlantic in a high-resolution ocean model, *J. Phys. Oceanogr.*, 33, 677-693. Hisard, P.,  
755 and C. Henin, (1987), Response of the Equatorial Atlantic Ocean to the 1983-1984 wind  
756 from the Programme Francais Ocean et Climat Dans l'Atlantique Equatorial cruise data set, *J.*  
757 *Geophys. Res.*, 92, C4, 3759-3768.

758 Hormann, V., and P. Brandt (2007), Atlantic Equatorial Undercurrent and associated cold tongue  
759 variability, *J. Geophys. Res.*, 112, C06017, doi:10.1029/2006JC003931.

760 Hormann, V., R. Lumpkin and G. Foltz (2012), Interannual North Equatorial Countercurrent  
761 variability and its relation to Tropical Atlantic climate modes, *J. Geophys. Res.*, 117,  
762 C04035, doi:10.1029/2011JC007697.

763 Hua, B. L., F. Marin, and R. Schopp (2003), Three-dimensional dynamics of the subsurface  
764 countercurrents and equatorial thermostat. Part I: Formulation of the problem and generic  
765 properties, *J. Phys. Oceanogr.*, 33, 2588–2609.

766 Hüttl-Kabus, S., and C. W. Böning (2008), Pathways and variability of the off-equatorial  
767 undercurrents in the Atlantic Ocean, *J. Geophys. Res. - Oceans*, 113, C10018, Doi:  
768 10.1029/2007JC004700.

769 Jochum, M., P. Malanotte-Rizzoli, and A. Busalacchi (2003), Tropical instability waves in the  
770 Atlantic Ocean, *Ocean Modell.*, 7, 146–163, doi:10.1016/S1463-5003(03)00042-8.

771 Jochum, M., and P. Malanotte-Rizzoli (2004), A new theory for the generation of the equatorial  
772 subsurface undercurrents, *J. Phys. Oceanogr.*, 34, 755–771.

773 Johnson, G. C., and D. W. Moore (1997), The Pacific subsurface countercurrents and an inertial  
774 model, *J. Phys. Oceanogr.*, 27(11), 2448–2459.

775 Katz, E. J. (1987), Seasonal response of the sea surface to wind in the equatorial Atlantic, *J.*  
776 *Geophys. Res.*, 92, 1885–1893.

777 Katz, E. J., and S. Garzoli(1982) Response of the western equatorial Atlantic Ocean to an annual  
778 wind cycle, *J. Mar. Res.*, 40, 307–327.

779 Katz, E. J. (1993), An interannual study of the Atlantic North Equatorial Countercurrent, *J. Phys.*  
780 *Oceanogr.*, 23, 1, 116-123.

781 Lagerloef, G., G. T. Mitchum, R. B. Lukas, and P. Niiler (1999), Tropical Pacific near surface  
782 currents estimated from altimeter, wind and drifter data, *J. Geophys. Res.*, 104, 23,313-  
783 23,326.

784 Le Traon P.-Y., F. Nadal, N. Ducet (1998), An improved mapping method of multisatellite  
785 altimeter data, *J. Atmos. Oceanic Technol.*, 15, 522-533, doi:10.1029/2005GL025551.

786 Lumpkin, R. and S. L. Garzoli (2005), Near-surface circulation in the tropical Atlantic Ocean,  
787 Deep-Sea Res., Part I, 52(3), 495-518.

788 Ma, H. (1996), The dynamics of North Brazil Current retroflexion eddies, J. Mar. Res., 54, 35–  
789 53.

790 Marin, F., F. Schopp, and B. L. Hua (2003), Three-dimensional dynamics of the subsurface  
791 countercurrents and equatorial thermostat. Part II: Influence of the large-scale ventilation  
792 and equatorial winds, J. Phys. Oceanogr., 33, 2610– 2626.

793 Metcalf, W., and M. C. Stalcup (1967), Origin of the Atlantic Equatorial Undercurrent, J.  
794 Geophys. Res., 72(20), 4959– 4975.

795 McCarthy, M. C., L. D. Talley, and D. Roemmich (2000), Seasonal to interannual variability  
796 from expendable bathythermograph and TOPEX/Poseidon altimeter data in the South Pacific  
797 subtropical gyre, J. Geophys. Res., 105(C8), 19,535–19,550, doi:10.1029/2000JC900056.

798 Moore, D. W., and S. G. H. Philander (1977), Modelling of the tropical oceanic circulation. The  
799 Sea, E. D. Goldberg, Ed., Marine Modelling, Vol. 6, Wiley and Sons, 319–361.

800 Mayer, D., R. L. Molinari, and M. O. Baringer (2001), Transition regions and their role in the  
801 relationship between sea surface height and subsurface temperature structure in the Atlantic  
802 Ocean, Geophys. Res. Lett., 28, 20, 3943-3946.

803 McCreary, J. P., P. Lu, and Z. Yu (2002), Dynamics of the Pacific subsurface countercurrents, J.  
804 Phys. Oceanogr., 32, 2379–2404.

805 Menkes, C. E., S. C. Kennan, P. Flament, Y. Dandonneau, S. Masson, B. Biessy, E. Marchal, G.  
806 Eldin, J. Grelet, Y. Montel, A. Morlière, A. Lebourges-Dhaussy, C. Moulin, G. Champalbert,  
807 and A. Herbland (2002), A whirling ecosystem in the equatorial Atlantic, Geophys. Res.  
808 Lett., 29, 1553, doi:10.1029/2001GL014576.

809 Moore, D. W., and S. G. Philander (1986), Modelling of the tropical Atlantic circulation. *The*  
810 *Sea*, E. D. Goldberg et al., Eds., Marine Modeling, Vol. 6, John Willey and Sons, 319-361.

811 Moura, A.D., and J. Shukla (1981), On the dynamics of Droughts in Northeast Brazil—  
812 observations, theory and numerical experiments with a general-circulation model. *J. Atmos.*  
813 *Sci.* 38, 2653–2675.

814 Nobre, P., and J. Shukla (1996), Variations of sea surface temperature, wind stress and rainfall  
815 over the tropical Atlantic and South America, *J. Clim.*, 9, 2464– 2479.

816 Pedlosky, J. (1988), Entrainment and the Termination of the Equatorial Undercurrent, *J. Phys.*  
817 *Oceanogr.*, 18, 880–886.

818 Peterson, R. G., and L. Stramma (1991), Upper-level circulation in the South Atlantic Ocean,  
819 *Prog. Oceanogr.*, 26, 1-73.

820 Philander, S. G. H., and R. C. Pacanowski (1986), A model of the seasonal cycle in the tropical  
821 Atlantic Ocean, *J. Geophys. Res.*, 91(C12), 14,192–14,206.

822 Phillips, H. E., and S. R. Rintoul (2002), A mean synoptic view of the subantarctic front south of  
823 Australia, *J. Phys. Oceanogr.*, 32(5), 1536-1553.

824 Pond, S. and G. Pickard (1983), *Introductory Dynamical Oceanography*, 2<sup>nd</sup> Ed., Pergamon Press  
825 Inc., NY 10523, USA.

826 Reverdin, G., P. Rual, Y. du Penhoat, and Y. Gouriou (1991), Vertical structure of the seasonal  
827 cycle in the central equatorial Atlantic Ocean: XBT sections from 1980 to 1988, *J. Phys.*  
828 *Oceanogr.*, 21, 277-291.

829 Reynolds, R. W., N. A. Rayner, T. M. Smith, D. C. Stokes, and W. Wang (2002), An improved  
830 in situ and satellite SST analysis for climate, *J. Clim.*, 15, 1609-1625.

831 Richardson, P.L., and D. Walsh (1986), Mapping climatological seasonal variations of surface  
832 currents in the tropical Atlantic using ship drifts, *J. Geophys. Res.*, 91, 10537-10550.

833 Richardson, P., and Reverdin, G., (1987), Seasonal cycle of velocity in the Atlantic NECC as  
834 measured by surface drifters, current meters and ship drifts. *J. Geophys. Res.*, 92, 3691–  
835 3708.

836 Richardson, S. Arnault, S. Garzoli, and J. G. Bruce, (1992), Annual cycle of the Atlantic North  
837 Equatorial Countercurrent. *Deep-Sea Res.*, 39, 997–1014.

838 Ridgway, K. R., R. C. Coleman, R. J. Bailey, and P. Sutton (2008), Decadal variability of East  
839 Australian Current transport inferred from repeated high-density XBT transects, a CTD  
840 survey and satellite altimetry, *J. Geophys. Res.*, 113, C08039, doi:10.1029/2007JC004664.

841 Rintoul, S. R., S. Sokolov, and J. Church (2002), A 6 year record of baroclinic transport  
842 variability of the Antarctic Circumpolar Current at 140°E derived from expendable  
843 bathythermograph and altimeter measurements, *J. Geophys. Res.*, 107(C10), 3155,  
844 doi:10.1029/2001JC000787.

845 Rio, M. H., S. Guinehut, and G. Larnicol (2011), New CNES-CLS09 global mean dynamic  
846 topography computed from the combination of GRACE data, altimetry, and in situ  
847 measurements, *J. Geophys. Res.*, 116, C07018, doi:10.1029/2010JC006505.

848 Roemmich, D., and J. Gilson (2001), Eddy Transport of Heat and Thermocline Waters in the  
849 North Pacific: A Key to Interannual/Decadal Climate Variability?, *J. Phys. Oceanogr.*, 31,  
850 675–687.

851 Rowe, G. D., E. Firing, and G. C. Johnson (2000), Pacific equatorial subsurface countercurrent  
852 velocity, transport, and potential vorticity, *J. Phys. Oceanogr.*, 30, 1172– 1187.

853 Schott, F., L. Stramma, and J. Fischer (1995), The warm water inflow into the western tropical  
854 Atlantic boundary regime, spring 1994, *J. Geophys. Res.*, 100, 24745–24760.

855 Schott, F., J. Fisher, and L. Stramma (1998), Transports and pathways of the upper-layer  
856 circulation in the western tropical Atlantic, *J. Phys. Oceanogr.*, 28, 1904–1928.

857 Schott, F. A., M. Dengler, P. Brandt, K. Affler, J. Fischer, B. Bourlès, Y. Gouriou, R. L.  
858 Molinari, and M. Rhein (2003), The zonal currents and transports at 35°W in the tropical  
859 Atlantic, *Geophys. Res. Lett.*, 30, 7, 1349, doi:10.1029/2002GL016849.

860 Schott, F., J. McCreary, and G. Johnson (2004), Shallow Overturning Circulations of the  
861 Tropical-Subtropical Oceans, In *Earth Climate: The Ocean-Atmosphere Interaction*, C.  
862 Wang, S.-P. Xie and J.A. Carton (eds.), AGU Geophysical Monograph Series, 147, 261–304.

863 Servain, J., I. Wainer, J.P. McCreary, and A. Dessier (1999), A relationship between the  
864 equatorial and meridional modes of climatic variability in the tropical Atlantic, *Geophys.*  
865 *Res. Lett.*, 26(4), 485-488, doi:10.1029/1999GL900014.

866 Stramma, L. S., and M. England (1999), On the water masses and mean circulation of the South  
867 Atlantic. *J. Geophys. Res.*, 104, 20,863-20,833.

868 Stramma, L., S. Hüttl, and J. Schafstall (2005), Water masses and currents in the upper tropical  
869 northeast Atlantic off northwest Africa, *J. Geophys. Res.*, 110, C12006,  
870 doi:10.1029/2005JC002939.

871 Sutton, R. T. and D. L. R. Hodson (2005), Atlantic Ocean forcing of North American and  
872 European summer climate. *Science*, 309 (5731), 115-118, doi: 10.1175/JCLI4038.1.

873 Tabata, S., B. Thomas, and D. Ramsden (1986), Annual and interannual variability of steric sea  
874 level along line P in the northeast Pacific ocean, *J. Phys. Oceanogr.*, 16, 1378-1398.



875 Torrence, C., and G. P. Compo (1998), A Practical Guide to Wavelet Analysis, Bulletin of the  
876 American Meteorological Society, 79, 1, 61-78. Software available at URL:  
877 <http://atoc.colorado.edu/research/wavelets/>.

878 Urbano, D. F., M. Jochum, and I. C. A. Silveira (2006), Rediscovering the second core of the  
879 Atlantic NECC, Ocean Modell.,12, 1–15, doi:10.1016/j.ocemod.2005.04.003.

880 Voituriez, B. (1981), Les sous courants equatoriaux Nord et Sud et la Formation des domes  
881 thermiques tropicaux, Oceanolog. Acta., 5, 497-506.

882 Wacongne, S. (1989). Dynamical regimes of a fully nonlinear stratified model of the Atlantic  
883 equatorial undercurrent. J. Geophys. Res., 94, 4810–15.

884 Weisberg, R. H., J. H. Hickman, T. Y. Tang, and T. J. Weingartner (1987), Velocity and  
885 temperature observations during the Seasonal Response of the Equatorial Atlantic  
886 Experiment at 0°,28°W, J. Geophys. Res., 92, 5061-5075.

887 Weisberg, R., and T. Weingartner (1988), Instability waves in the equatorial Atlantic Ocean, J.  
888 Phys. Oceanogr., 18, 1641– 1657.

889 Xie, S.-P., and S.G.H. Philander (1994), A coupled ocean-atmosphere model of relevance to the  
890 ITCZ in the eastern Pacific, Tellus, 46A, 340-350.

891 Yang, J., and T M. Joyce, (2006), Local and Equatorial Forcing of Seasonal Variations of the  
892 North Equatorial Countercurrent in the Atlantic Ocean, J. Phys. Oceanogr.,26, 238-254.

893 Zhang, D., M. J. McPhaden, and W. E. Johns (2003), Observational Evidence for Flow between  
894 the Subtropical and Tropical Atlantic: The Atlantic Subtropical Cells, J. Phys. Oceanogr., 33,  
895 1783-1797.

896

897

898

899 **Figure Captions:**

900

901 **Figure 1:** a) Root mean square of SLA (cm) in the tropical Atlantic (filled contours), with  
902 superimposed AX08 sections: selected sections (black), sections not included in the analysis  
903 (gray), and mean of the selected sections (red). b) Probability density function of the averaged  
904 longitude between 10°S-10°N for all the AX08 XBT sections. c) Histogram of the monthly  
905 distribution of the number of the sections before (blue) and after (red) the selection of the  
906 sections that fit into the 68 percentile of b).

907 **Figure 2:** Typical temperature (°C) and salinity (psu) distributions for January (a, c) and June (b,  
908 d) 2010. Salinity is derived from WOD01 climatological T-S relationships. The mean locations  
909 of the eastward currents are marked in (a), and the approximate depths of the water masses are  
910 indicated in (c).

911 **Figure 3:** Anomalies of surface dynamic height ( $DH_0'$ , red), with respect to the annual mean of  
912 the XBT data, along with corresponding thermosteric (blue) and halosteric (green) components  
913 at (a) 4.5°N, (b) the equator, and (c) 4.5°S.

914 **Figure 4:** Comparison between SLA and surface dynamic height anomalies ( $DH_0'$ ): Longitude-  
915 time diagrams of (a) SLA, (b)  $DH_0'$ , and (c)  $SLA - DH_0'$ ; dots on the right-hand side of (c)  
916 mark the realizations of the AX08 XBT transect. (d) Linear fit between  $DH_0'$  and SLA.

917 **Figure 5:** Correlation at each depth and latitude between surface dynamic height anomalies  
918 ( $DH_0$ ) and: a) density anomalies ( $\sigma_\theta$ ), and b) dynamic height anomalies ( $DH'$ ). The thick black  
919 lines in b) mark the mean location of the isopycnals  $\sigma_\theta = 24.5 \text{ kg m}^{-3}$  and  $\sigma_\theta = 26.8 \text{ kg m}^{-3}$  that  
920 define the upper and lower dynamic layers, respectively.

921 **Figure 6:** Mean seasonal cross-sectional velocities derived from XBT data (coloring), with  $\sigma_\theta =$   
922  $24.5 \text{ kg m}^{-3}$  and  $\sigma_\theta = 26.8 \text{ kg m}^{-3}$  isopycnals overlaid (black lines): (a) Dec, Jan; (b) Mar, Apr,  
923 May; (c) Jun, Jul, Aug; and (d) Sep, Oct, Nov. The seasonal-mean absolute surface dynamic  
924 height is shown on top of each panel.

925 **Figure 7:** Geostrophic transports estimated from the synthetic method (black line with open  
926 dots) and from the XBT data (blue dots), with corresponding linear fit between the two  
927 transport estimates: (a, b) NECC, (c, d) NEUC, (e, f) SEUC, and (g, h) EUC.

928 **Figure 8:** Seasonal cycle of geostrophic transport ( $Sv$ ), core velocity (cm/s), and latitudinal  
929 position (deg): (a, e, i) NECC, (b, f, j) NEUC, (c, g, k) SEUC, and (d, h, l) EUC. Red dots are  
930 monthly XBT averages and red line represents the corresponding fit of annual and semi-annual  
931 harmonics; green and blue lines mark the average synthetic estimates for the periods 1993-  
932 2000 and 2000-2010, respectively. Shown error bars are the standard errors for the synthetic  
933 estimates, and the standard error plus fitting error for the XBT data.

934 **Figure 9:** Time series and respective wavelet transforms for the monthly (a) NECC transport, (b)  
935 NEUC transport, (c) EUC transport, and (d) zonal pseudo-wind stress averaged over the region  
936  $5^\circ\text{S}$ - $5^\circ\text{N}$ / $30^\circ\text{W}$ - $20^\circ\text{W}$ . Transport time series are generated by the synthetic method. The  
937 wavelet power spectra are based on a Morlet transform and regions above the 95% significance  
938 level are encircled by black contours, with the bowl-shaped black lines indicating the cone of  
939 influence.

940 **Figure 10:** a) Instantaneous correlation between the standardized NECC transport anomalies and  
941 both SSTA (coloring) and pseudo-wind stress anomalies (vectors). Only the statistically  
942 significant values are shown. b) Time series of the standardized NECC transport anomalies  
943 (red), SSTA index (blue), and meridional wind index (y-wind, black). The SSTA index is

944 defined by subtracting the averages over the northern ( $35^{\circ}\text{W}-17^{\circ}\text{W}/0-7^{\circ}\text{N}$ ) and southern  
945 ( $30^{\circ}\text{W}-15^{\circ}\text{W}/11^{\circ}\text{S}-23^{\circ}\text{S}$ ) boxes as marked in (a), and the wind index is defined as the average  
946 of the meridional pseudo-wind stress anomalies in the northern box. c) Lagged correlations of  
947 the NECC with the SST (blue) and wind indices (black), with significant values marked with  
948 bold lines.

949 **Figure 11:** Same as Figure 10, but for NEUC transport anomalies. The SST index (blue) is  
950 defined here as the difference between the northern box average over the Guinea Dome region  
951 ( $10^{\circ}-25^{\circ}\text{N}/15^{\circ}-35^{\circ}\text{W}$ ) and the southern box average ( $8^{\circ}\text{S}-5^{\circ}\text{N}/30^{\circ}\text{W}-10^{\circ}\text{E}$ ) as marked in (a),  
952 and the y-wind index is defined as the average of meridional pseudo-wind stress anomalies  
953 over the northern box.

954

955

956 **Tables captions:**

957 **Table 1** – Latitudinal and isopycnal ranges used in the volume transport calculations of the  
958 Atlantic equatorial eastward currents.

959 **Table 2** – Percentage of the variance of transport, core velocity, and position explained by the  
960 annual and semi-annual harmonics for each current using XBT estimates and the synthetic  
961 method for the periods 1990-2000 (S/1990-2000) and 2000-2010 (S/2000-2010). The synthetic  
962 method is not analyzed for the SEUC.

963

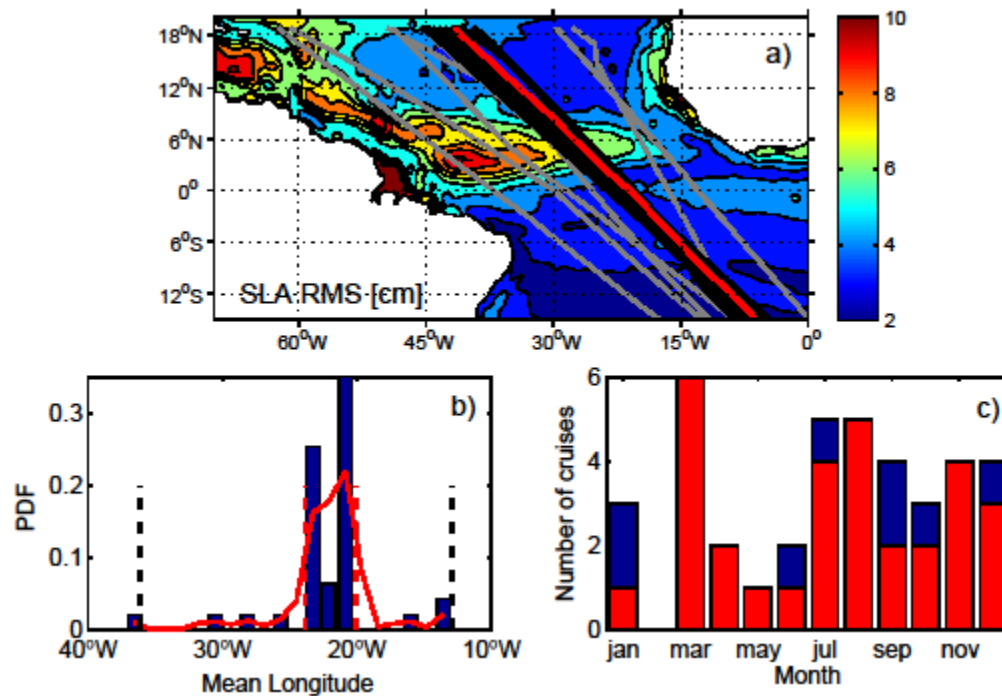
964

965 **Figures:**

966

967

968



969

970

971 Figure 1: a) Root mean square of SLA (cm) in the tropical Atlantic (filled contours), with

972 superimposed AX08 sections: selected sections (black), sections not included in the analysis

973 (gray), and mean of the selected sections (red). b) Probability density function of the averaged

974 longitude between 10°S-10°N for all the AX08 XBT sections. c) Histogram of the monthly

975 distribution of the number of the sections before (blue) and after (red) the selection of the

976 sections that fit into the 68 percentile of b).

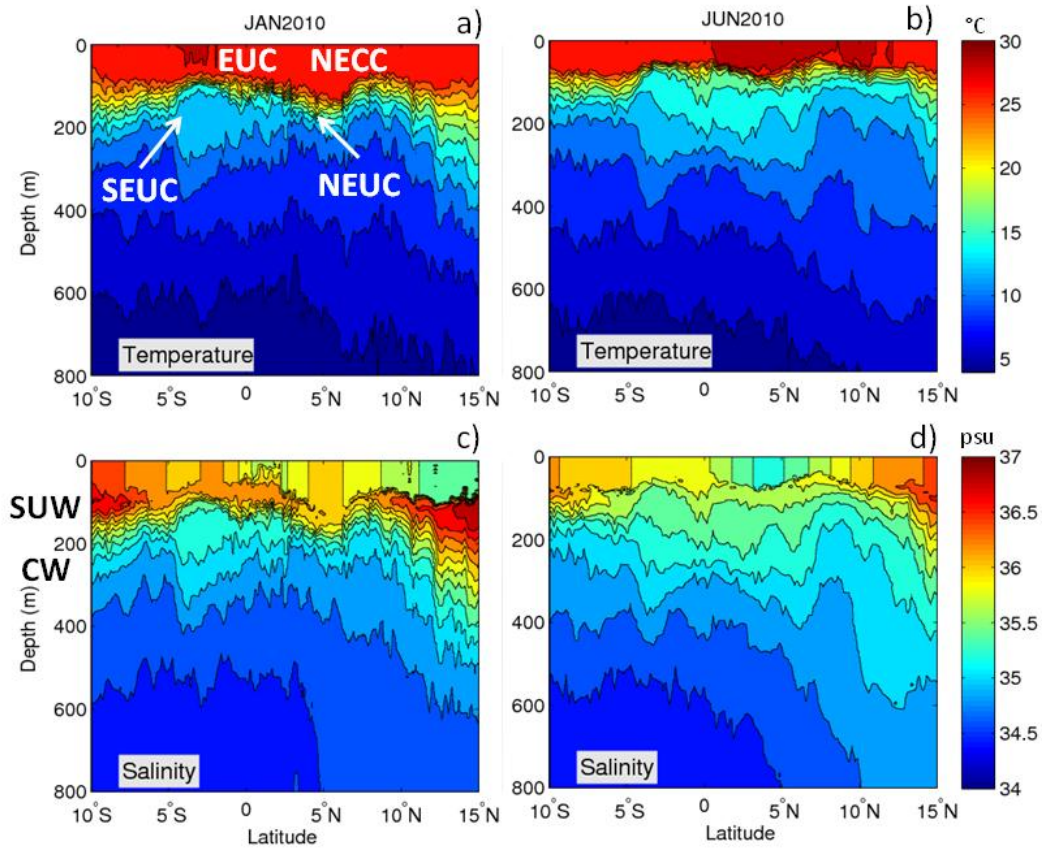
977

978

979

980

981



982

983 Figure 2: Typical temperature (°C) and salinity (psu) distributions for January (a, c) and June (b,

984 d) 2010. Salinity is derived from WOD01 climatological T-S relationships. The mean locations

985 of the eastward currents are marked in (a), and the approximate depths of the water masses are

986 indicated in (c).

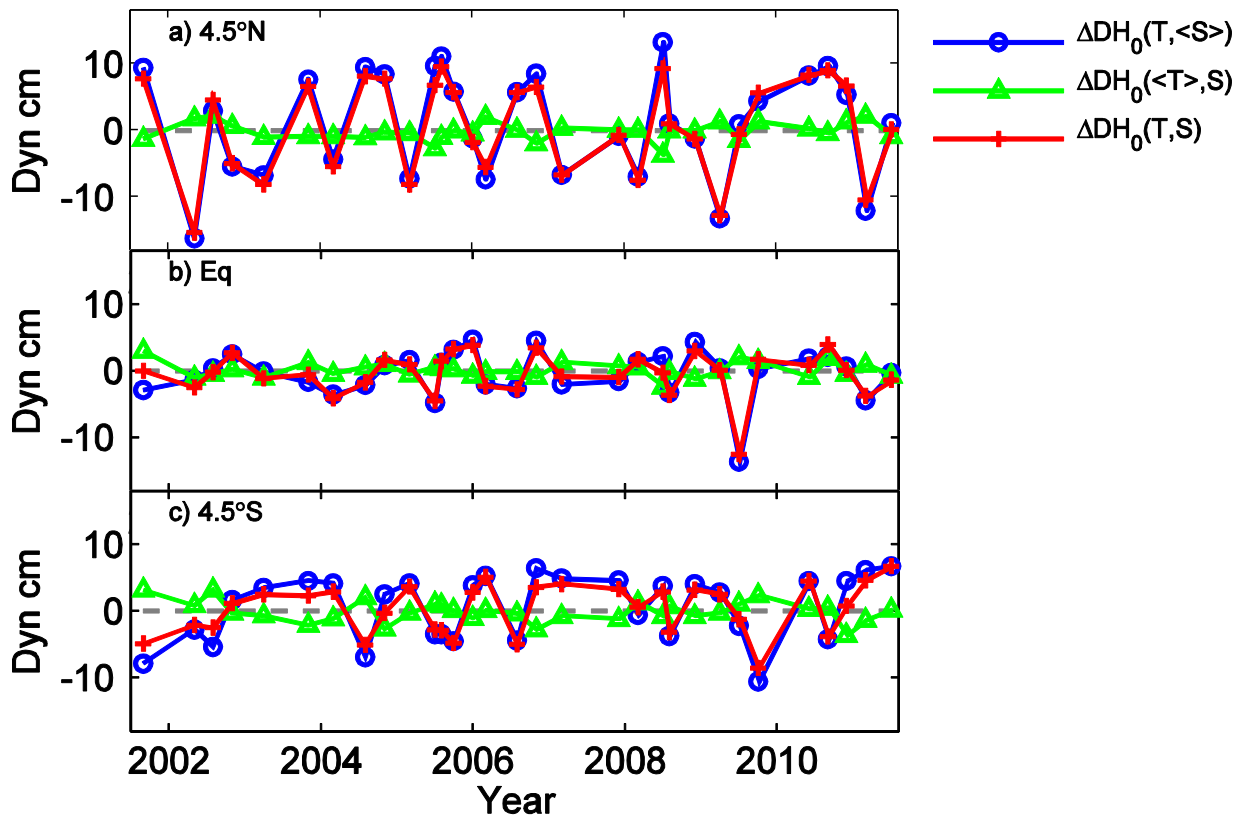
987

988

989

990

991



992

993 Figure 3: Anomalies of surface dynamic height ( $\Delta DH_0'$ , red), with respect to the annual mean of  
994 the XBT data, along with corresponding thermosteric (blue) and halosteric (green) components  
995 at (a)  $4.5^\circ\text{N}$ , (b) the equator, and (c)  $4.5^\circ\text{S}$ .

996

997

998

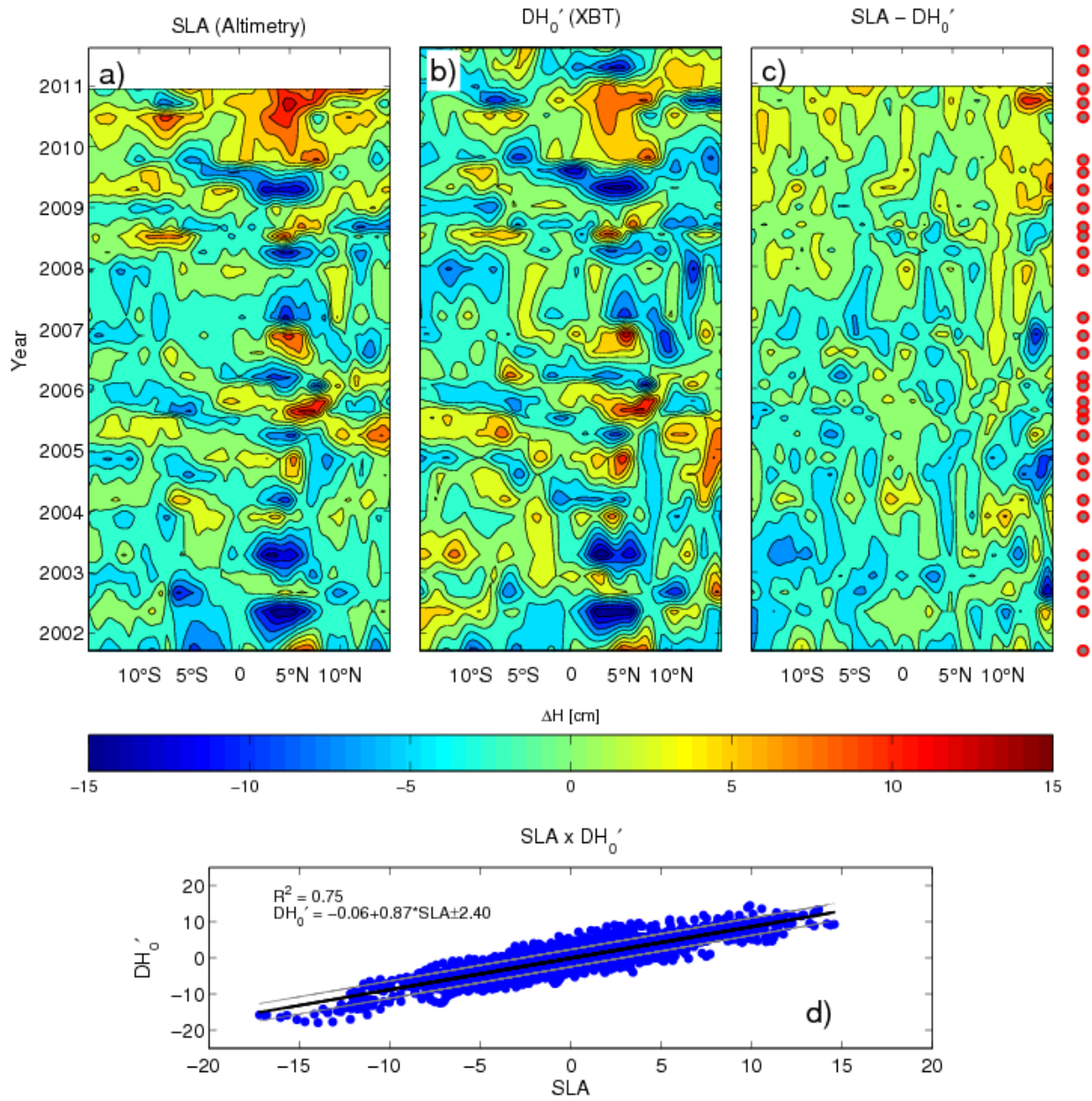
999

1000

1001

1002





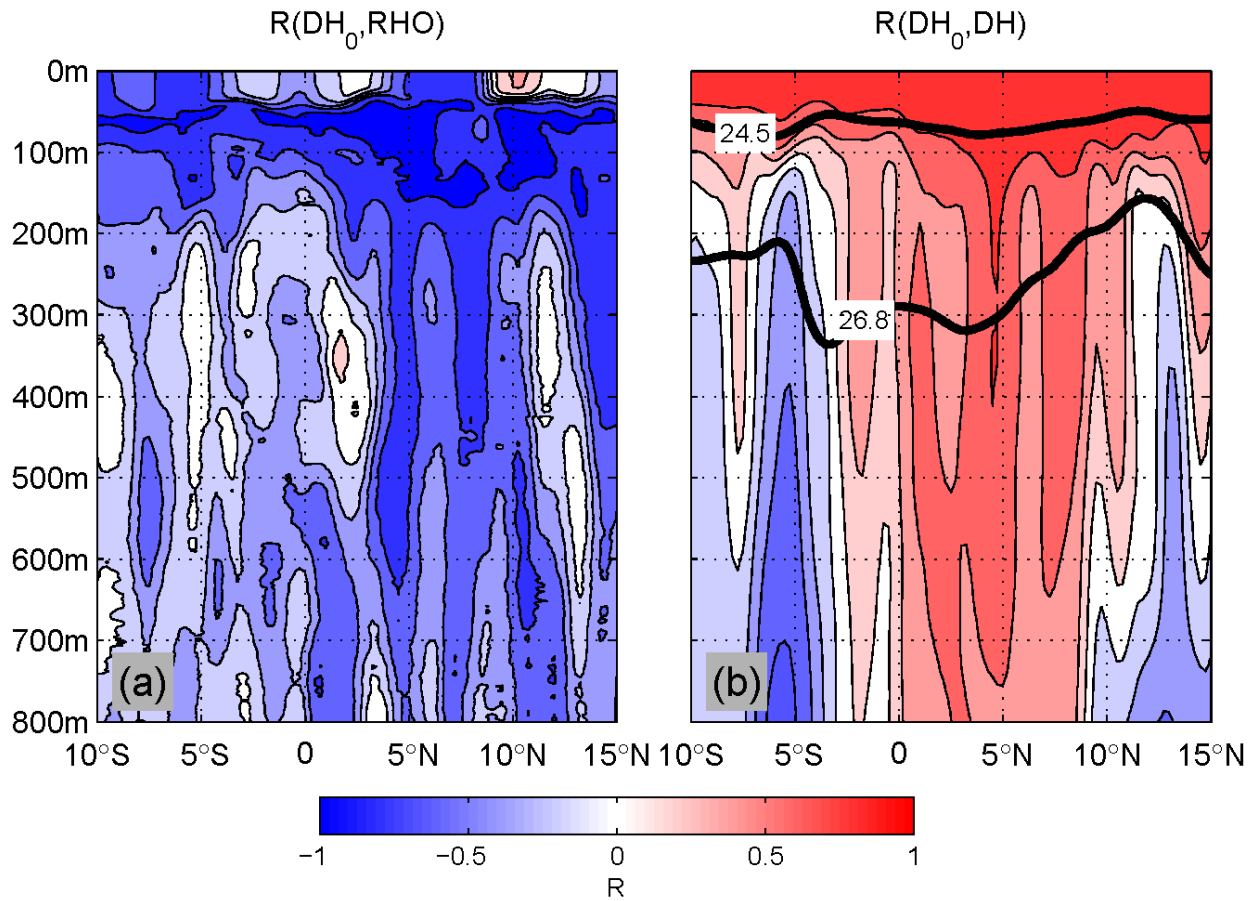
1003

1004

1005 Figure 4: Comparison between SLA and surface dynamic height anomalies ( $DH_0'$ ): Longitude-  
 1006 time diagrams of (a) SLA, (b)  $DH_0'$ , and (c)  $SLA - DH_0'$ ; dots on the right-hand side of (c) mark  
 1007 the realizations of the AX08 XBT transect. (d) Linear fit between  $DH_0'$  and SLA.

1008

1009



1010

1011 Figure 5: Correlation at each depth and latitude between surface dynamic height anomalies  
 1012 ( $DH_0$ ) and: a) density anomalies ( $\sigma_\theta$ ), and b) dynamic height anomalies ( $DH'$ ). The thick black  
 1013 lines in b) mark the mean location of the isopycnals  $\sigma_\theta = 24.5 \text{ kg m}^{-3}$  and  $\sigma_\theta = 26.8 \text{ kg m}^{-3}$  that  
 1014 define the upper and lower dynamic layers, respectively.

1015

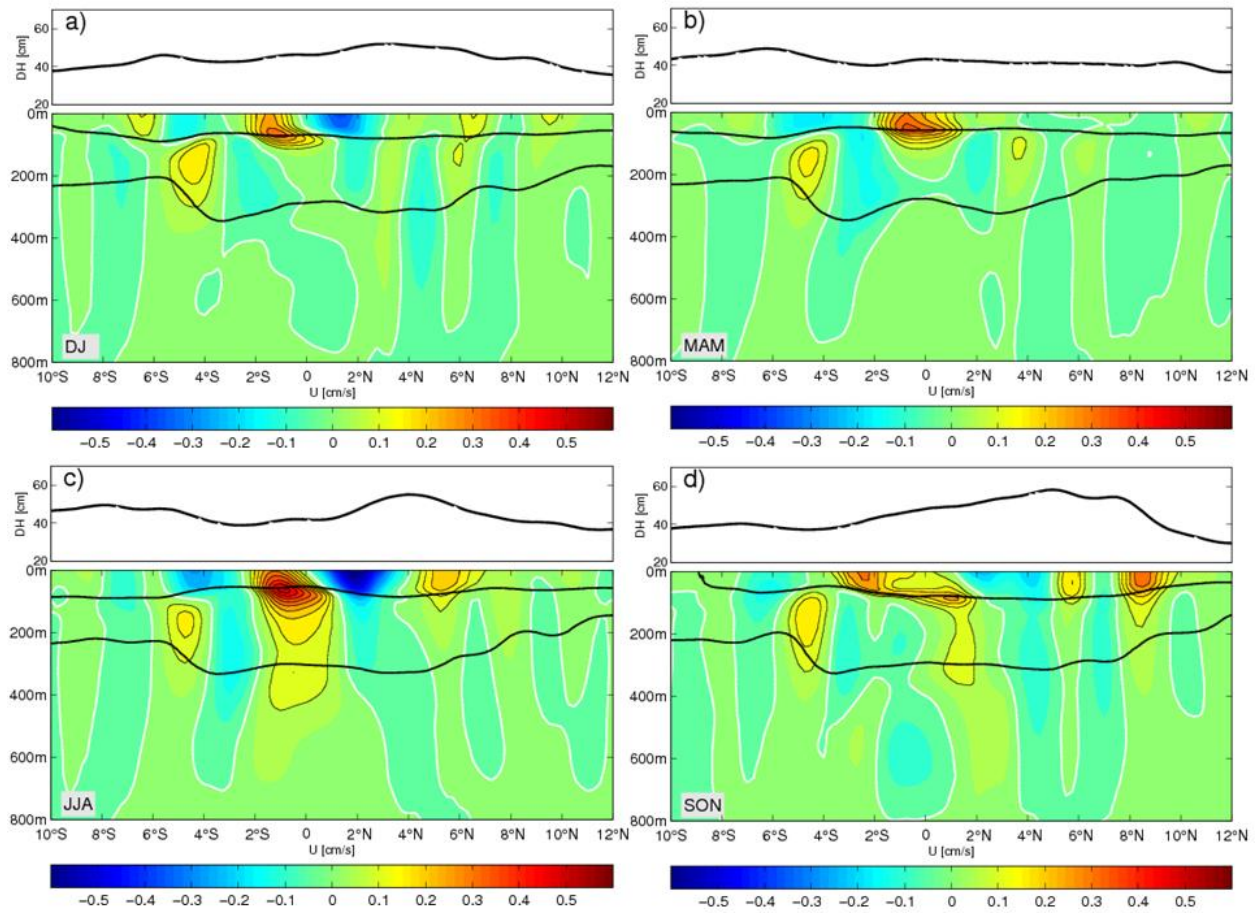
1016

1017

1018

1019

1020



1021

1022

1023 Figure 6: Mean seasonal cross-sectional velocities derived from XBT data (coloring), with  $\sigma_\theta =$

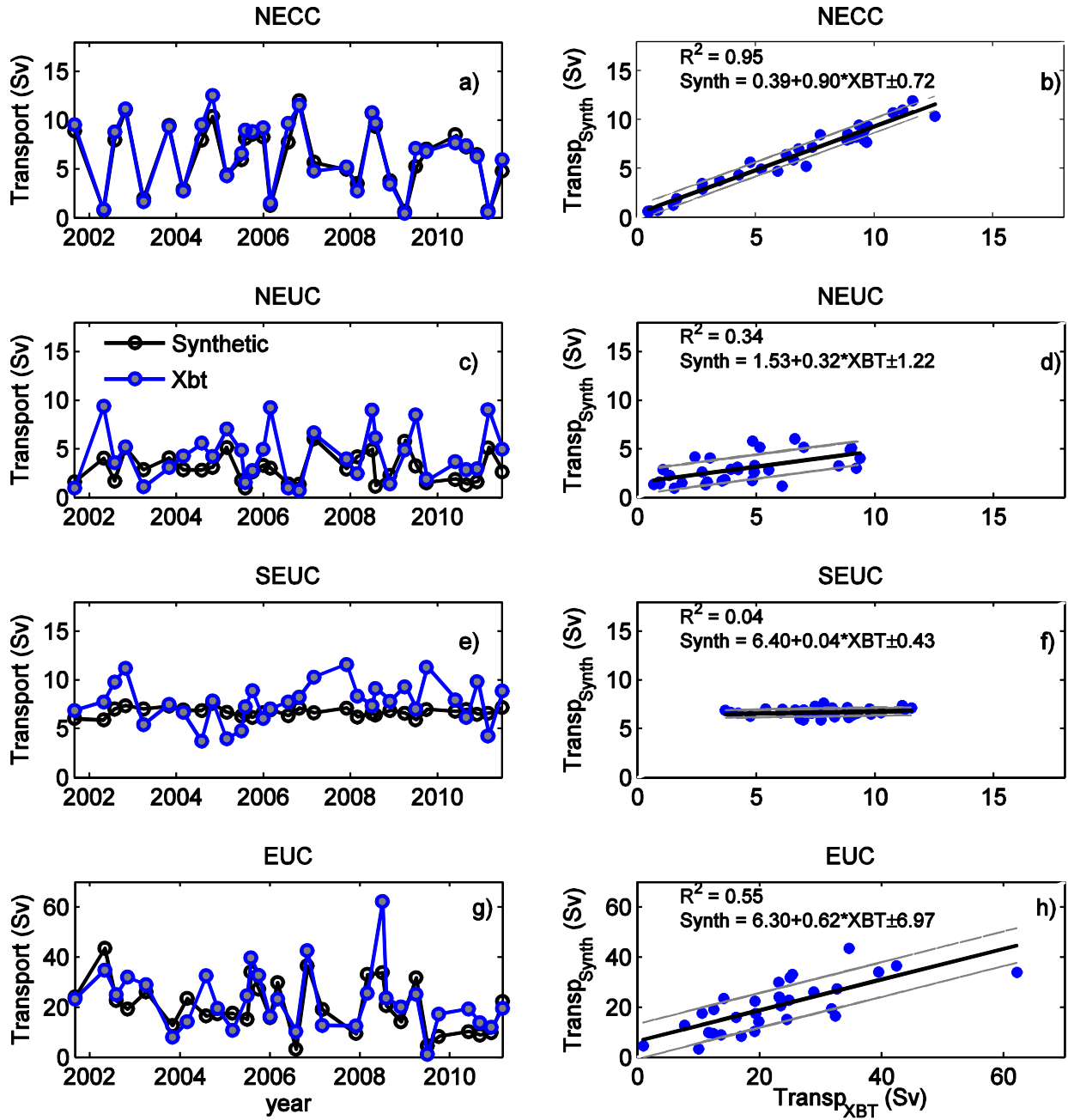
1024  $24.5 \text{ kg m}^{-3}$  and  $\sigma_\theta = 26.8 \text{ kg m}^{-3}$  isopycnals overlaid (black lines): (a) Dec, Jan; (b) Mar, Apr,

1025 May; (c) Jun, Jul, Aug; and (d) Sep, Oct, Nov. The seasonal-mean absolute surface dynamic

1026 height is shown on top of each panel.

1027

1028



1029

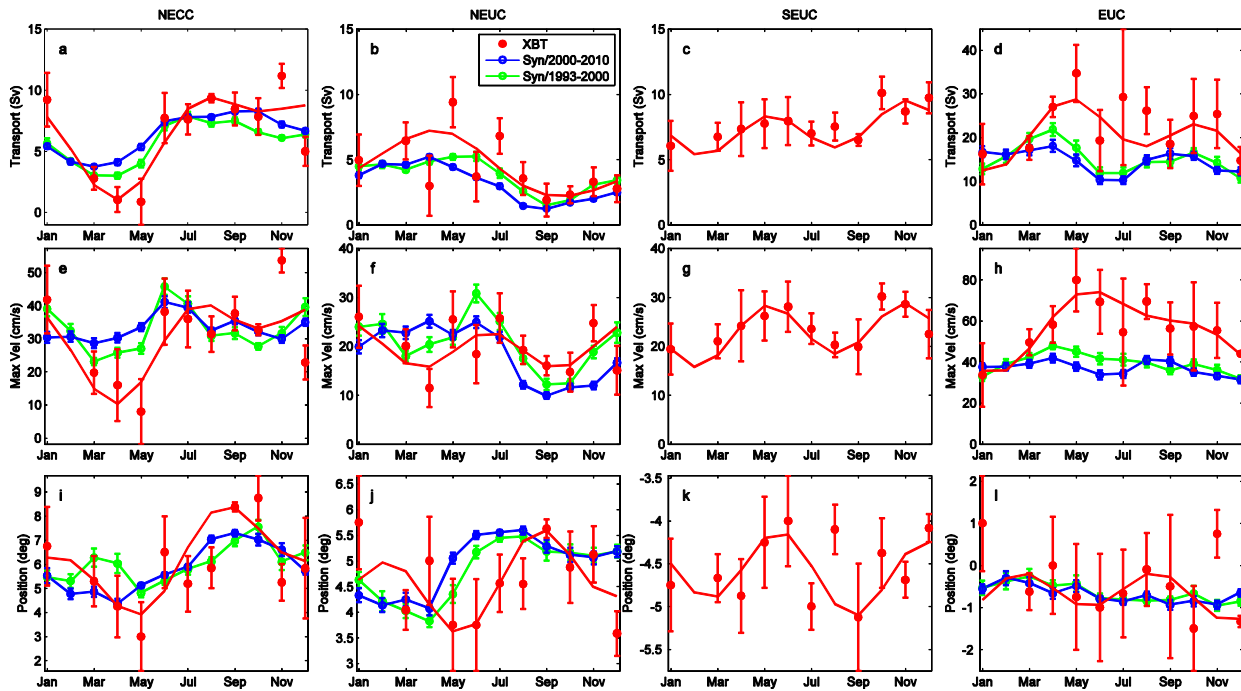
1030 Figure 7: Geostrophic transports estimated from the synthetic method (black line with open dots)  
 1031 and from the XBT data (blue dots), with corresponding linear fit between the two transport  
 1032 estimates: (a, b) NECC, (c, d) NEUC, (e, f) SEUC, and (g, h) EUC.

1033

1034

1035

1036



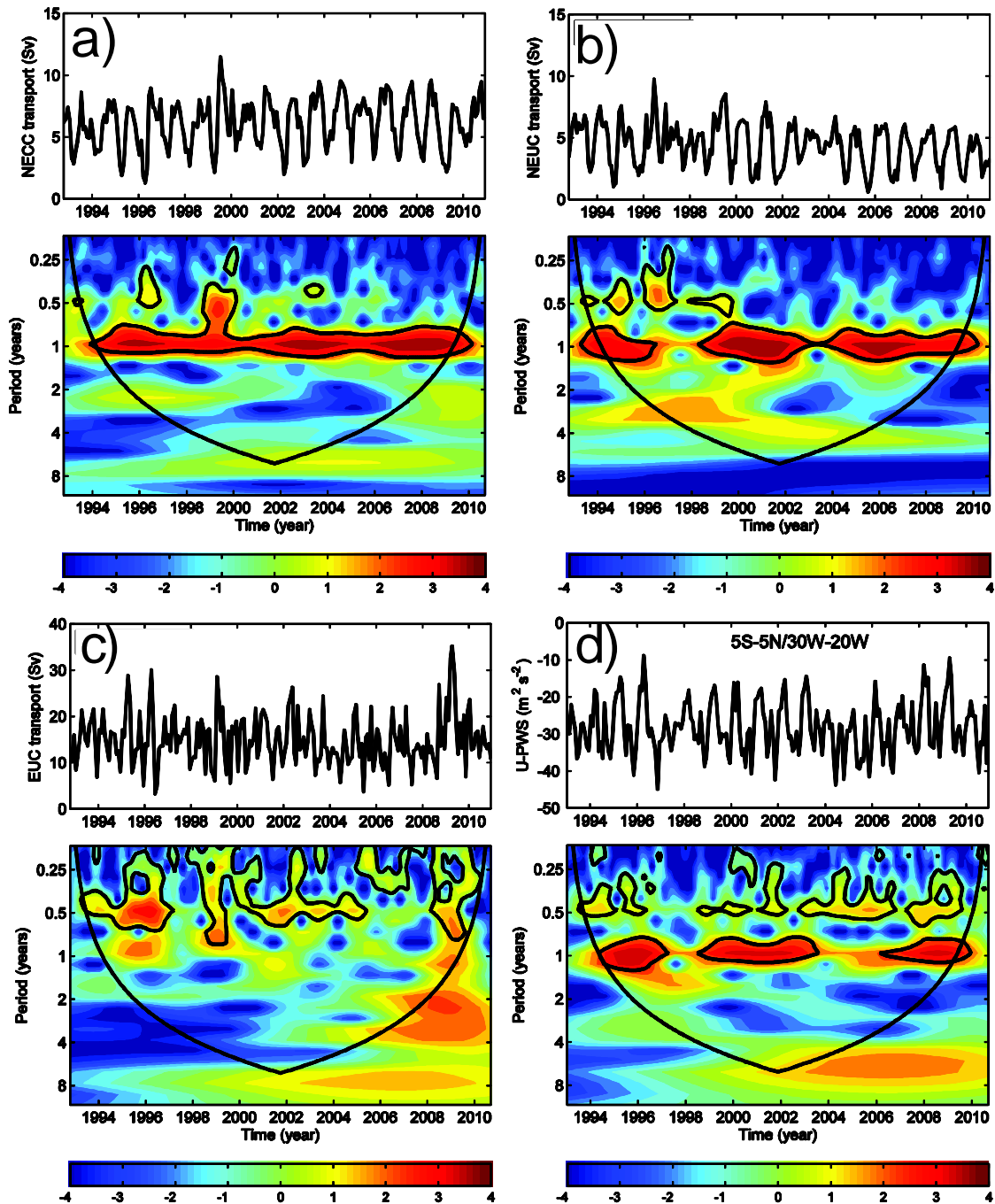
1037

1038

1039 Figure 8: Seasonal cycle of geostrophic transport (Sv), core velocity (cm/s), and latitudinal  
1040 position (deg): (a, e, i) NECC, (b, f, j) NEUC, (c, g, k) SEUC, and (d, h, l) EUC. Red dots are  
1041 monthly XBT averages and red line represents the corresponding fit of annual and semi-annual  
1042 harmonics; green and blue lines mark the average synthetic estimates for the periods 1993-2000  
1043 and 2000-2010, respectively. Shown error bars are the standard errors for the synthetic estimates,  
1044 and the standard error plus fitting error for the XBT data.

1045

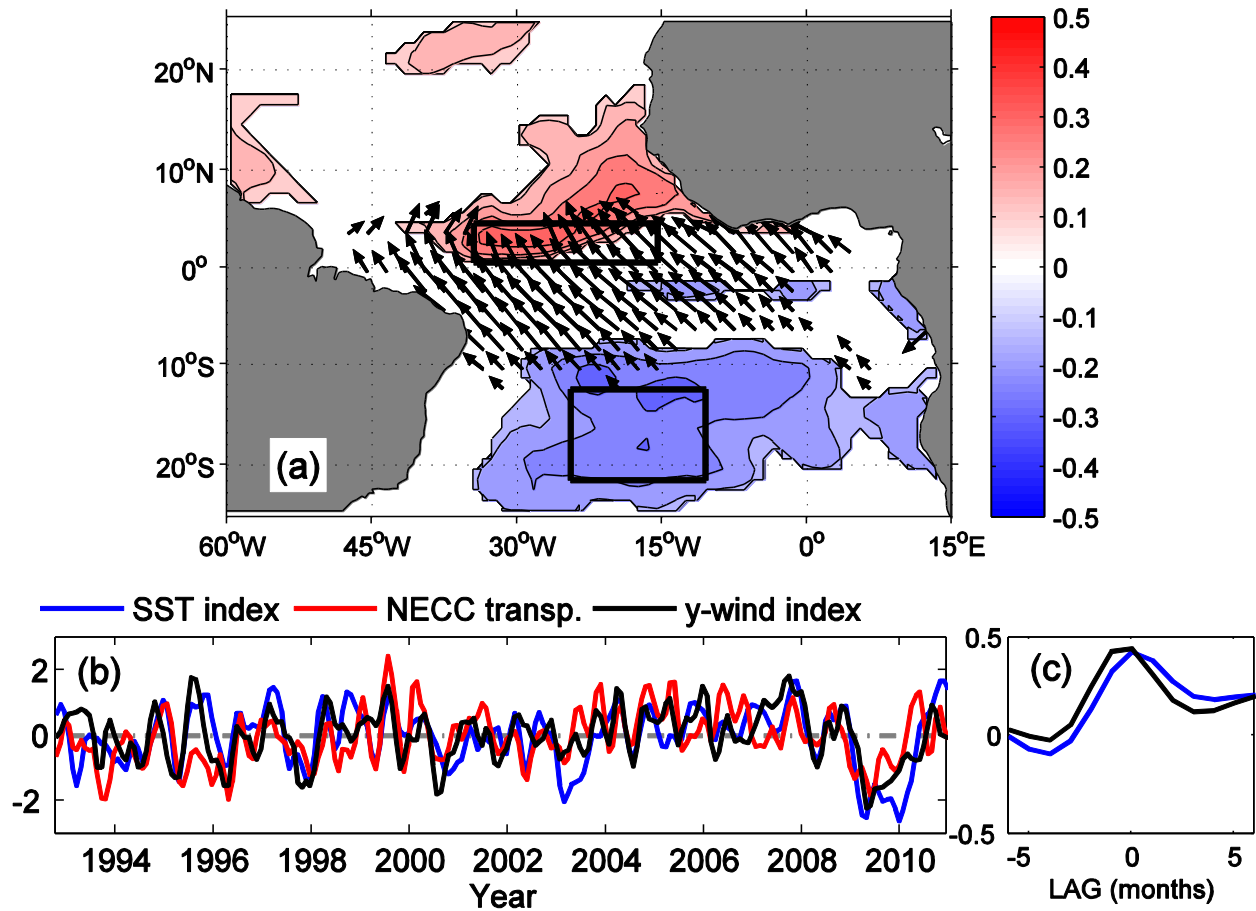
1046



1047

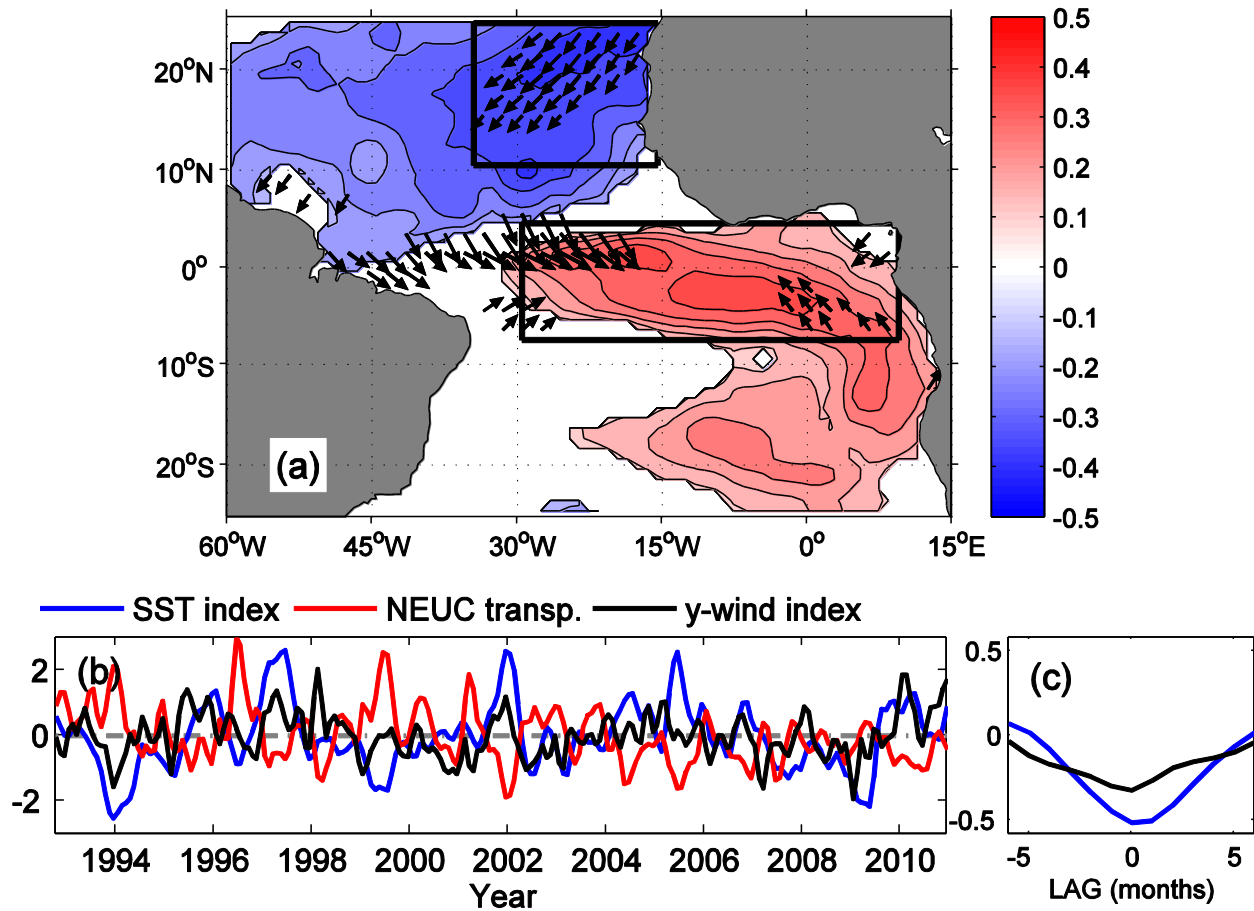
1048 Figure 9: Time series and respective wavelet transforms for the monthly (a) NECC transport, (b)  
 1049 NEUC transport, (c) EUC transport, and (d) zonal pseudo-wind stress averaged over the region  
 1050  $5^{\circ}\text{S}-5^{\circ}\text{N}/30^{\circ}\text{W}-20^{\circ}\text{W}$ . Transport time series are generated by the synthetic method. The wavelet

1051 power spectra are based on a Morlet transform and regions above the 95% significance level are  
 1052 encircled by black contours, with the bowl-shaped black lines indicating the cone of influence.  
 1053



1054  
 1055 Figure 10: a) Instantaneous correlation between the standardized NECC transport anomalies and  
 1056 both SSTA (coloring) and pseudo-wind stress anomalies (vectors). Only the statistically  
 1057 significant values are shown. b) Time series of the standardized NECC transport anomalies (red),  
 1058 SSTA index (blue), and meridional wind index (y-wind, black). The SSTA index is defined by  
 1059 subtracting the averages over the northern (35°W-17°W/0-7°N) and southern (30°W-  
 1060 15°W/11°S-23°S) boxes as marked in (a), and the wind index is defined as the average of the  
 1061 meridional pseudo-wind stress anomalies in the northern box. c) Lagged correlations of the

1062 NECC with the SST (blue) and wind indices (black), with significant values marked with bold  
 1063 lines.  
 1064



1065  
 1066 Figure 11: Same as Figure 10, but for NEUC transport anomalies. The SST index (blue) is  
 1067 defined here as the difference between the northern box average over the Guinea Dome region  
 1068 ( $10^{\circ}$ - $25^{\circ}$ N/ $15^{\circ}$ - $35^{\circ}$ W) and the southern box average ( $8^{\circ}$ S- $5^{\circ}$ N/ $30^{\circ}$ W- $10^{\circ}$ E) as marked in (a), and  
 1069 the y-wind index is defined as the average of meridional pseudo-wind stress anomalies over the  
 1070 northern box.

1071  
 1072



1073

1074

1075

1076 Table 1: Latitudinal and isopycnal ranges used in the volume transport calculations of the

1077 Atlantic equatorial eastward currents.

1078

Current	Latitude	$\sigma_\theta$ (kg m <sup>-3</sup> )
NECC	3°N - 10°N	0 - 24.5
NEUC	3°N - 6°N	24.5 - 26.8
EUC	2.5°S - 2.5°N	0 - 26.8
SEUC	6°S - 3°S	24.5 - 26.8

1079

1080

1081

1082

1083

1084

1085

1086

1087

1088

1089

1090

1091  
 1092  
 1093  
 1094  
 1095  
 1096  
 1097  
 1098  
 1099

Table 2: Percentage of the variance of transport, core velocity, and position explained by the annual and semi-annual harmonics for each current using XBT estimates and the synthetic method for the periods 1990-2000 (S/1990-2000) and 2000-2010 (S/2000-2010). The synthetic method is not analyzed for the SEUC.

Current	Data	Transport (%)		Core Velocity (%)		Position (%)	
		Annual	Semi-annual	Annual	Semi-annual	Annual	Semi-annual
NECC	XBT	58	13	41	20	52	19
	S/1990-2000	78	17	5	66	64	14
	S/2000-2010	93	3	31	23	95	1
NEUC	XBT	63	1	1	35	22	41
	S/1990-2000	75	21	23	64	80	17
	S/2000-2010	99	1	83	5	76	13
SEUC	XBT	14	90	8	92	5	60
EUC	XBT	24	40	72	16	2	21
	S/1990-2000	28	64	79	22	70	15
	S/2000-2010	19	68	13	54	69	10

1100