

1 Climate response to tropical cyclone-induced ocean mixing in an
2 Earth system model of intermediate complexity

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19 **1. Abstract**

20 We introduce a parameterization of ocean mixing by tropical cyclones (TCs) into
21 an Earth system model of intermediate complexity. The parameterization is based on
22 previously published global budgets of TC-induced mixing derived from high-resolution
23 satellite measurements of anomalous sea surface temperatures along storm tracks.
24 Recognizing the caveats introduced, for example, by the simplified model structure, we
25 find that the representation of realistic TC-induced mixing substantially alters the
26 equilibrium conditions of (i) the thermal structure of the upper ocean, (ii) the surface
27 energy budget, and (iii) the circulation in the equatorial to subtropical Pacific ocean.
28 These changes result in warmer upwelling regions in the eastern equatorial Pacific and an
29 overall increase in ocean heat content consistent with the recent TC heat pump
30 hypothesis. Spatial variability in the mixing appears to be a key factor in the modeled
31 response. We find no substantial influence of the considered TC-induced mixing on
32 poleward ocean heat transport in the analyzed model. Our results suggest that climate-
33 sensitive feedbacks are plausible, however, the large-scale effect is mainly confined to
34 the subtropical Indo-Pacific region for present-day TC climatology.

35

36 **2. Introduction**

37 Understanding the role of tropical cyclones (TCs) within the climate system is an
38 active area of research. There is strong evidence linking low frequency variability in TC
39 activity to tropical sea surface temperatures during the past 60 years [*Emanuel, 2005*].
40 However, the projected response in activity to future warming is less clear [*Emanuel et*
41 *al., 2008; Knutson et al., 2008; Bender et al., 2010; Sabbatelli et al., submitted*].
42 Furthermore, feedbacks may exist which enable TCs to actively contribute to the
43 dynamics of the climate system, rather than passively responding to changes in the large-
44 scale mean state.

45 *Emanuel* [2001] proposed that vertical ocean mixing induced by TC winds may
46 be responsible for the majority of the present-day poleward ocean heat transport. In
47 general, TC winds generate near-inertial internal waves which eventually break [*Black*
48 *and Dickey, 2008*], mixing warm surface water down into the thermocline where it is
49 available to be advected away from the storm regions by the larger scale ocean
50 circulation. *Emanuel* [2001] hypothesized that this oceanic heat convergence is
51 eventually carried poleward by the meridional overturning circulation, estimating the
52 majority of present-day poleward ocean heat transport can be attributed to TC-induced
53 ocean mixing. Several observation-based studies support this hypothesis [*Sriver and*
54 *Huber, 2007; Sriver et al., 2008*]. However, these studies find estimates of TC-induced
55 oceanic heat convergence, and the heat available to be transported poleward, is more
56 conservative than the original *Emanuel* [2001] estimate (~30% of peak heat transport
57 values at storm latitudes). Nonetheless, if TCs are capable of influencing tropical
58 temperature patterns through feedbacks associated with ocean mixing and transport, then

59 these events may be an important factor for understanding the nature of climate
60 variability.

61 The amount of wind energy available to mix the ocean depends on the power
62 dissipated at the surface by friction [*Emanuel*, 2005]. The power dissipation is an
63 integrated measure of TC intensity. It represents the convolution of several cyclone
64 characteristics including wind speed, size, duration, and frequency. While much debate
65 currently focuses on understanding and predicting changes in single metrics such as
66 intensity and frequency, integrated quantities such as power dissipation appear to be more
67 important for describing potential impacts of changes in TC activity on climate. Thus,
68 inferring TC-induced impacts on ocean mixing based on any single metric is incomplete,
69 though there is evidence that certain TC characteristics co-vary (e.g. intensity and
70 duration/frequency) [*Sriver and Huber*, 2007b].

71 Several recent modeling studies have sought to determine the importance of TC-
72 induced ocean mixing on upper ocean properties and transport. Notably, *Korty et al.*
73 [2008] show that including an interactive mixing parameterization, based on TC
74 maximum potential intensity, into an intermediate complexity climate model positively
75 influences poleward heat transport in climate scenarios with increased atmospheric
76 carbon dioxide, suggesting TCs could be an important factor for sustaining warm
77 climates with a small equator to pole temperature gradient. *Jansen and Ferrari* [2009]
78 show that meridional variability in prescribed TC-induced ocean mixing can inhibit
79 poleward ocean heat transport by influencing the subtropical overturning circulation.
80 They find increased heat convergence at the edge of the subtropics associated with
81 enhanced vertical mixing is eventually transported equatorward in the return branch of

82 the subtropical cells, thus limiting the influence of TC-induced mixing on extratropical
83 transport. *Sriver and Huber* [in press] test the sensitivity of an ocean general circulation
84 model to satellite-based global TC surface winds. They find transient, extreme surface
85 wind forcing alters the subtropical overturning, consistent with *Jansen and Ferrari*
86 [2009]. Furthermore, under scenarios with enhanced TC wind forcing, *Sriver and Huber*
87 [in press] find increased oceanic heat convergence in the tropics and warmer
88 temperatures in upwelling regions, resulting in a permanent El Niño-like climate state.
89 These findings support the idea that increased tropical ocean mixing may have
90 contributed to sustaining a permanent El Niño during the Pliocene (3-5 million years ago)
91 [*Brierley et al.*, 2009; *Fedorov et al.*, 2010].

92 Here we diagnose the impact of TCs on the large-scale state by incorporating a
93 global parameterization of ocean mixing by these events into an Earth System Model of
94 Intermediate Complexity (EMIC). This parameterization is based on TC-induced mixing
95 budgets developed previously from satellite measurements of anomalous sea surface
96 temperature along storm tracks [*Sriver and Huber*, 2007; *Sriver et al.*, 2008]. The mixing
97 parameterization varies horizontally and vertically. Our aim is simply to test the first-
98 order equilibrium response of large-scale model properties to the inclusion of a simplified
99 (yet arguably realistic) representation of the present-day, observation-based climatology
100 of TC-induced mixing. Furthermore, using an intermediate complexity model enables us
101 to perform a suite of long-term simulations to full equilibrium (including the deep ocean),
102 with relatively low computational burden, and to analyze the effects of varying levels of
103 prescribed background mixing in order to ascertain the relative contribution of TCs to
104 other mixing processes not resolved by the model.

105 Because the adopted EMIC (similar to many other EMICS) does not yet contain a
106 dynamic atmosphere, we do not capture the full extent of possible atmospheric feedbacks
107 using this modeling approach. For example, TC-induced changes in tropical sea surface
108 temperature patterns may have important implications for the large-scale mean
109 atmospheric circulations such as the Hadley [*Sriver and Huber, in press*] and Walker
110 circulations. Furthermore, changes in regional and tropical surface temperature can
111 influence basin-wide TC activity metrics, such as frequency, intensity, and spatial
112 distribution [*Wang et al., 2008; Zhao et al., 2009*]. Since our TC mixing parameterization
113 is prescribed, we cannot account for feedbacks that can potentially affect overall TC
114 activity, and more importantly the induced ocean mixing. Here we focus primarily on
115 modeling the impacts of realistic TC-induced mixing on the ocean. This approach is
116 useful for testing the first order response of the ocean to spatially-varying vertical mixing
117 from the present-day global TC climatology. We are presently working on incorporating
118 a climate-sensitive component to our mixing parameterization, utilizing a more
119 sophisticated fully-coupled ocean/atmosphere general circulation model, in order to
120 assess impacts on the coupled ocean-atmosphere system.

121 The paper is organized as follows: section 3 describes the climate model and
122 experimental design, section 4 contains the results and discussion of the model
123 experiment (including effects on thermal structure, surface energy budget, ocean heat
124 content and transport, and circulation dynamics), and section 5 provides a short
125 description of our main conclusions and the implications.

126

127 **3. Model**

128 We use the University of Victoria Earth System Model (UVic) [*Weaver et al.*,
129 2001] version 2.8, which features a 3 dimensional ocean general circulation model based
130 on the Modular Ocean Model (MOM) version 2 [*Pacanowski, 1995*]. The UVic model
131 includes a simple energy-moisture balance atmosphere model with prescribed, diagnosed
132 winds, as well as thermodynamic/dynamic sea-ice and thermomechanical land-ice
133 components. This version of the model also includes terrestrial vegetation and carbon
134 cycling [*Meissner et al., 2003*], and ocean biogeochemistry, based on the ecosystem
135 model of *Schmittner et al.* [2005].

136 The ocean model is coarsely resolved (1.8° latitude x 3.6° longitude and 19
137 vertical levels) and features several key mixing parameterizations, including the Gent-
138 McWilliams isopycnal mixing parameterization [*Gent and McWilliams, 1990*], diapycnal
139 mixing over rough topography by tidal forcing [*Simmons et al., 2004*], and increased
140 vertical mixing rates ($1 \text{ cm}^2/\text{s}$) below 500 meters depth in the Southern Ocean (south of
141 40°S) [*Schmittner et al., 2009*]. In addition, the model uses a prescribed background
142 vertical ocean diffusivity (K_v) to simulate the effects of sub-grid scale mixing processes
143 not captured by the parameterizations listed above. Recent studies attempt to constrain
144 the uncertainty of K_v used in the UVic model [e.g., *Schmittner et al., 2009*], suggesting
145 values between 0.2 and $0.3 \text{ cm}^2/\text{s}$ yield best agreement with observed tracer fields. Here
146 we perform an ensemble of simulations that span a range of K_v values from 0.1 to 0.5
147 cm^2/s , in order to examine the relative contribution of K_v in combination with the added
148 prescribed vertical mixing by TCs.

149 The TC ocean mixing parameterization is based on global mixing budgets
150 developed from satellite-based measurements of anomalous surface temperature along

151 storm tracks [*Sriver and Huber, 2007; Sriver et al., 2008*]. While TCs are transient
152 events, we seek to simplify their modeled representation by characterizing the mixing as
153 annualized diffusivities applied in combination with the K_v values. Figure 1A shows a
154 map of the TC mixing rates at the surface. The pattern of mixing is spatially variable
155 with the largest values typically occurring in the regions with the most TC activity. The
156 mixing depths also vary, and we derive these values from the estimated changes in annual
157 mixed layer depth shown previously [*Sriver et al., 2008*]. Because anomalous mixed
158 layer depth represents only a portion of the total depth affected by TC-induced mixing,
159 we apply an idealized correction by multiplying anomalous mixed layer depth values by
160 3x. This simplistic approach yields mixing length scales on the order of ~100 meters for
161 diffusivities equal to $1 \text{ cm}^2/\text{s}$. Maximum mixing depths penetrate to ~250 meters in areas
162 with the largest TC diffusivities. These mixing rates and length scales are similar to
163 recent estimates based on theoretical arguments [*Korty et al., 2008*] and modeling results
164 testing the sensitivity of upper ocean properties to TC wind forcing [*Sriver and Huber, in*
165 *press*]. The model is coarsely resolved in the vertical direction, therefore, we linearly
166 interpolate diffusivity in the deepest grid boxes in order to reflect decreased mixing
167 where mixing depths occur between model levels.

168 We perform an ensemble of 10 model simulations that span a wide range of UVic
169 K_v values (0.1, 0.2, 0.3, 0.4, 0.5 cm^2/s). For each K_v , we perform 2 simulations
170 corresponding to cases with TC-induced mixing and without. The model simulations are
171 initiated from modern-day climatology and run for 3000 years to approximate
172 equilibrium. In the final 1000 years of the simulation, the ensemble members are run in a
173 carbon-coupled mode, which couples the atmospheric carbon to the land and ocean

174 models. Atmospheric carbon dioxide levels are prescribed to pre-industrial levels
175 throughout the simulations. We do not address transient climate change forced by
176 increasing atmospheric greenhouse gases in this study. The only difference between each
177 pair of runs for a given K_V value is the addition of the TC mixing parameterization that is
178 applied to the oceanic vertical mixing budget. This methodology provides a simple, and
179 flexible, global representation of realistic TC mixing rates suitable for diagnosing the first
180 order importance of these events within the current climate. The simulation results and
181 analysis routines are available from the lead author upon request.

182

183 **4. Results and Discussion**

184

185 **4.1 Upper-Ocean Thermal Response**

186 The addition of TC-induced mixing in the model significantly alters the upper
187 thermal structure of the global ocean (Figure 1B). The largest effect is seen in the Pacific
188 basin. The northwestern Pacific region exhibits pronounced cooling in the areas
189 experiencing the largest amount of TC activity, consistent with the cyclone-induced
190 ocean heat pump mechanism [*Sriver and Huber, 2007*]. The model exhibits warmer
191 near-surface temperatures in the eastern equatorial Pacific region which is mostly devoid
192 of TCs, consistent with recent independent modeling studies [*Jansen and Ferrari, 2009*;
193 *Fedorov et al, 2010*; *Sriver and Huber, in press*]. These upwelling regions are a key
194 component of the subtropical overturning circulation. This shallow meridional
195 overturning consists of poleward surface Ekman transport, sinking at the edge of the
196 subtropical gyres, and equatorward flow of cooler water at depths of ~200 meters

197 [McCreary and Lu, 1994; Klinger and Marotzke, 2000]. In the Pacific, the equatorward
198 flow feeds into the Equatorial Undercurrent, where it is eventually upwelled in the
199 eastern equatorial Pacific region. Thus, our results suggest a substantial portion of the
200 heat pumped into the interior ocean in the Pacific basin is carried equatorward by the
201 return branch of the subtropical cells, leading to anomalously warm upwelling regions in
202 the equatorial cold tongue in the eastern Pacific. This feature suggests Pacific TC
203 activity could provide a mechanism for sustaining a permanent El Niño [Brierley *et al.*,
204 2009; Fedorov *et al.*, 2010] and is consistent with independent model results testing the
205 sensitivity of an ocean general circulation model to TC winds [Sriver and Huber, in
206 press].

207 While we find some influence on near-surface temperature in other regions
208 experiencing TCs such as the North Atlantic (Figure 1B), the magnitude of the
209 anomalous temperature is substantially less than in the Pacific basin. This result indicates
210 that, within the current global climatology, the effects of TC-induced mixing are largely
211 confined to the dynamics of the Pacific ocean.

212

213 **4.2 Surface Energy Budget**

214 Although UVic does not contain a dynamic atmosphere component, the simplified
215 energy-moisture balance atmosphere model allows us to diagnose the first-order response
216 of the surface energy budget to changes in upper ocean temperatures forced by TC-
217 induced mixing. The anomalous downward surface heat flux over the ocean between an
218 equilibrated case with TC-induced mixing and the corresponding control is displayed in
219 Figure 2A (for background diffusivity $K_v=0.2 \text{ cm}^2/\text{s}$). The addition of TC-induced

220 mixing has a strong effect on the model's global surface heat budget over the ocean.
221 Generally, TCs cause a net increase in downward surface heat flux in regions where they
222 occur. The most prominent region of increased downward surface heat flux occurs in the
223 northwestern Pacific (Figure 2A), corresponding to the region experiencing the largest
224 amount of TC activity and coolest temperature response (Figure 1). Figure 2A shows
225 decreased downward heat flux in the eastern equatorial Pacific, consistent with the
226 anomalously warm near-surface temperatures in the eastern equatorial Pacific discussed
227 previously (Figure 1B).

228 Figure 2B represents the zonally integrated downward surface fluxes from Figure
229 2A. We find increased oceanic heat convergence at latitudes experiencing TCs, which is
230 again consistent with the cyclone-induced heat pump hypothesis. The low latitudes
231 experience decreased downward surface heat flux associated with warmer eastern
232 equatorial temperatures. The ocean-to-atmosphere heat flux in the mid-latitudes is
233 increased (shown by negative downward flux in Figure 2B), representing warmer
234 temperatures at those latitudes caused by increased poleward heat transport associated
235 with TC-induced ocean mixing.

236

237 **4.3 Ocean Heat Content and Transport**

238 Following *Emanuel's* [2001] hypothesis (discussed in the introduction), ocean
239 heat convergence by TC-induced mixing should be balanced by increased heat transport
240 under equilibrium conditions. Originally, it was presumed this heat transport would be
241 poleward, thus linking TCs to the wind-driven subtropical gyres and the global
242 meridional overturning circulation, which carry heat and mass to high latitudes. While

243 there is evidence of enhanced poleward heat transport, which increases the ocean-to-
244 atmosphere heat flux in the mid-latitudes (Figure 2B), a substantial portion of this heat
245 remains in the tropics, thus lowering ocean heat uptake in the equatorial oceans. The
246 reason for this modeled response is due to the current-day climatology of TC-induced
247 mixing, which occurs primarily in the Pacific basin between 10 and 30 degrees north. As
248 described previously (section 4.1), the subtropical overturning dominates upper ocean
249 transport in these regions. This circulation tends to trap TC-induced ocean heat in the
250 tropics, which warms the eastern boundary upwelling regions. The tropical warming
251 associated with this TC mixing parameterization appears to limit the effectiveness of TCs
252 as a possible tropical thermostat for the present-day climatology (e.g. distribution,
253 strength, and depth of mixing).

254 We diagnose global impacts on heat distribution and transport by examining the
255 influence of cyclone mixing on upper ocean heat content (Figures 2C and 2D) for the
256 case with background mixing $K_v=0.2 \text{ cm}^2/\text{s}$. We find a net increase in heat content
257 between 0 and 500 meters depth for much of the tropics, with the largest warming
258 occurring along the eastern Pacific boundaries. Conversely, the Pacific basin exhibits
259 cooling at intermediate depths between 500 and 1500 meters. Near the surface, the
260 positive anomalous heat content in the Atlantic basin is less than for the other regions.
261 However, the spatial distribution is more uniform in the Atlantic compared to the Pacific
262 and Indian basins. Figure 2C depicts increased heat content along the western boundary
263 in the north Atlantic downstream of the largest amount of TC-induced mixing in that
264 region. This suggests the interaction of the subtropical gyre with the TC-induced mixing
265 is responsible for redistributing some heat poleward in the Atlantic ocean. In addition,

266 we find the mean barotropic streamfunction to be sensitive to TC-induced mixing in both
267 the Atlantic and Pacific basins, though there is only a slight positive influence on the
268 strength of the global meridional overturning circulation.

269 The modeled total poleward ocean heat transport scales with K_v in the simulations
270 without TC-induced mixing, but it is relatively insensitive to the TC mixing
271 parameterization (Figure 3). TC-induced mixing tends to inhibit poleward heat transport
272 out of the tropics in the northern hemisphere, consistent with previous model results
273 [*Jansen and Ferrari, 2009*]. Subtropical poleward transport is increased in the both
274 hemispheres, which suggests at least some effect of TC-induced mixing on heat transport
275 by the wind-driven subtropical gyres. However, the relative contribution is small
276 compared to the peak poleward fluxes at these latitudes.

277 As discussed previously, the heat content in the uppermost 500 meters is
278 increased in the Pacific basin for the simulations with TC-induced mixing, especially in
279 the tropical and subtropical latitudes (Figure 2C). However, below 500 meters, we find
280 decreased ocean heat content in the northern hemisphere maximizing in the central
281 equatorial Pacific (Figure 2D). This feature can be understood by analyzing two
282 simulations with idealized scenarios of prescribed TC-induced mixing. In the first
283 simulation, we applied a constant mixing of $1 \text{ cm}^2/\text{s}$ across the entire tropics (30S to 30N)
284 to a depth of 200 meters. In the second simulation, we applied the zonally averaged
285 diffusivities (and mixing depths) from our observation-derived estimates (Figure 1A),
286 reflecting a simple zonally-invariant representation of TC-induced mixing that preserves
287 meridional variability. Both idealized simulations were performed for a background
288 diffusivity of $K_v=0.2 \text{ cm}^2/\text{s}$ and were carried out identically to the other ensemble

289 members (see Methods for details). These additional test cases allow us to diagnose the
290 mixing characteristic responsible for the altered thermal structure (e.g. spatial variability
291 in the zonal, meridional, and/or vertical directions).

292 We find that both simulations with idealized mixing result in more uniform
293 warming of the tropics at all levels in the model, as well as increased poleward ocean
294 heat transport and meridional overturning strength. Thus, there is no cooling below 500
295 meters in the equatorial Pacific. These simulations suggest that the decrease in equatorial
296 Pacific temperature below 500 meters (Figure 2D) in the model is caused by zonal
297 variability in TC-induced mixing. In other words, zonal mixing gradients in the upper
298 200 meters can potentially impact the thermal structure of the deep ocean. Changes in
299 these gradients associated with variability in TC-induced mixing may affect heat uptake
300 by the ocean, though timescales at which these effects become important are unclear.
301 Additionally, these results suggest that simplified approaches used to simulate TC-
302 induced mixing, such as broadly increasing vertical diffusivity equally across the entire
303 tropics or certain latitude bands, may over-estimate the influence of these events on the
304 large-scale dynamics. More realistic representations that capture the meridional, zonal,
305 and vertical variability of the mixing may be necessary to accurately simulate the climate
306 response, particularly in scenarios where the characteristics of the mixing may change as
307 a function of the climate state.

308

309 **4.4 Dynamical Impacts**

310 To examine the influence of TC-induced mixing on the dynamics of the equatorial
311 Pacific, we now focus on the upper ocean temperature and velocity for that basin. Figure

312 4A shows the zonally averaged temperature and meridional velocity for the control run
313 with background diffusivity $K_v=0.2 \text{ cm}^2/\text{s}$, and Figure 4B displays the difference between
314 the corresponding run with TC-induced mixing and the control shown in 4A. Figure 4B
315 reflects the substantial increase in temperature within the uppermost 500 meters in the
316 tropics and the decreased temperature in the deeper ocean. The meridional structure of
317 the subtropical overturning circulation is apparent in Figure 4A, with poleward transport
318 in the upper-most 100 meters in both hemispheres and the deeper equatorward flow from
319 ~100 to ~300 meters depth. We find a positive response in the strength of this
320 overturning in the northern hemisphere Pacific, where the majority of TC-induced mixing
321 occurs. There is little influence on the circulation strength in the southern hemisphere
322 cell.

323 Vertical profiles of the zonally averaged anomalous temperature and eastward
324 velocity for all the TC runs are compared to the corresponding controls for the Pacific
325 basin in Figure 5. The temperature response in Figure 5A is consistent with Figure 4B,
326 reflecting the warm anomaly between the surface and ~400 meters depth. We find that
327 the Equatorial Undercurrent is enhanced in the TC runs, along with intensified westward
328 flow beneath 400 meters depth (Figure 5B). It is important to note peak values of the
329 modeled Equatorial Undercurrent are considerably less than observations (15 cm/s
330 compared to ~100 cm/s), which is a typical limitation in coarse-resolution ocean general
331 circulation models.

332 The vertical profiles (Figure 5) show that the general response in modeled
333 temperature and velocity are relatively insensitive to the choice of K_v , though the
334 magnitude of the effect depends on the magnitude of K_v . The magnitudes of the

335 temperature anomalies are inversely proportional to K_v , with an increased upper ocean
336 warm anomaly for lower background mixing values. The response in the strength of the
337 equatorial near-surface currents depends on the magnitude of the K_v , however, the
338 westward velocity anomaly below 200 meters does not. This result suggests that the
339 equatorial westward flow at depth is not a result of the overall mixing strength in the
340 tropics, rather it is due to the spatial variability of the mixing in the zonal directions.

341 We isolate the zonal structure of the equatorial velocity profile in Figure 5 for the
342 case with background diffusivity $K_v=0.2$. The control case (Figure 6A) shows a well-
343 defined EUC with peak flow from the central to eastern Pacific regions. In Figure 6B, we
344 see the influence of TC-induced mixing on the velocity structure. The model exhibits
345 pronounced eastward flow out of the western Pacific at ~200 meters depth, which feeds
346 warmer water into the Equatorial Undercurrent. As a result, the 20 °C isotherm is
347 deepened by ~10 meters in the TC cases. The intensified equatorial undercurrent is
348 accompanied by deeper anomalously westward flow throughout much of the central
349 Pacific that reaches a maximum in the western Pacific, which is caused primarily by the
350 enhanced TC-induced mixing along the western boundary. This mixing may have
351 impacts on the dynamics on the Indo-Pacific warm pool that could affect transport into
352 the Indian Ocean via the Indonesian Throughflow [*Wijffels et al., 2008*].

353 In the model, the TC mixing parameterization increases heat transport into the
354 Indian Ocean via the Indonesian Throughflow, while total volume transport through the
355 Indonesian Throughflow is decreased. In the upper 200 meters, volume transport is
356 increased, which carries anomalously warm water from the western Pacific to the Indian
357 Ocean. Below 200 meters, the westward flow is decreased, and temperature gradients

358 between the basins are small. The net effect is a small decrease in the column-integrated
359 volume transport, along with increased heat transport, within the Indonesian
360 Throughflow. Given the resemblance of the modeled thermal structure to El Niño-like
361 conditions, decreased volume transport is consistent with recent observational results that
362 indicate decreased transport during El Niño events [*Tillinger and Gordon, 2009*]. Thus,
363 our findings suggest TC-induced mixing may influence the dynamics of the Indo-Pacific
364 warm pool, contributing to redistributing heat and mass via the Indonesian Throughflow.
365 However, it is not clear yet whether this result is realistic, given the simplified
366 topography of the Indonesian Throughflow contained in the coarse resolution UVic
367 model. Future experiments with higher-resolution ocean models should better elucidate
368 the robustness of this mechanism.

369

370 **5. Conclusions**

371 We implement a realistic representation of present-day TC-induced ocean mixing
372 in the UVic Earth System Model. We find the thermal structure of the modeled ocean to
373 be sensitive to our parameterization of this process for all considered values of oceanic
374 vertical background mixing. The thermal response results in altered near-surface
375 temperatures and surface fluxes, along with a redistribution of ocean heat by the large-
376 scale circulation. The main dynamical impacts are confined to the Pacific ocean, where
377 the mixing modifies the subtropical overturning circulation. This results in a decreased
378 zonal temperature gradient with cooler temperatures in the western region and warmer
379 temperatures in the eastern equatorial Pacific. This El Niño-like near-surface temperature
380 pattern is maintained by intensified eastward flow of warm water out of the western

381 equatorial-to-subtropical Pacific, ultimately feeding into the Equatorial Undercurrent and
382 upwelling regions in the eastern Pacific.

383 Our results suggest that TC-induced mixing based on present-day climatology can
384 have dynamical implications for the global ocean, but the impacts are primarily limited in
385 our model to the tropical (and subtropical) latitudes. We find little evidence that TC-
386 induced mixing influences equator-to-pole poleward ocean heat transport or the
387 meridional overturning circulation within the context of the current climatology.
388 However, our findings do support an active role of TCs within climate that is capable of
389 modifying the thermal and dynamical structure of the ocean, and the possibility of
390 impacts on the meridional overturning in the context of climate change cannot be ruled
391 out. Moreover, the spatial variability of this mixing appears to be a key factor in
392 determining the extent of its impact. TC-induced mixing may have important
393 implications for surface fluxes, tropical temperature distributions and circulation patterns
394 in the tropical and subtropical oceans. A better understanding of how climate-induced
395 changes in TC activity (and the associated changes in ocean mixing) could impact ocean
396 properties and dynamics is needed, though it is unclear at what timescales these processes
397 become important—an issue that will be addressed in future work.

398 Our current results and interpretations are adorned by the caveats introduced by
399 the lack of a dynamic atmosphere component in UVic. We are, for example, unable to
400 address coupling feedbacks such as changes in trade wind forcing caused by large-scale
401 surface temperature anomalies. Because this forcing is important for regulating the
402 strength of the subtropical overturning, potential feedbacks are likely missing. However,
403 the use of flexible intermediate complexity models such as UVic allows us to examine

404 the first order climate response to TC-induced mixing using an ensemble approach,
405 where we can rigorously test various TC mixing scenarios spanning the full parameter
406 space of applicable background ocean mixing values. Future work is needed to address
407 the possibility of climate feedbacks using a fully-coupled ocean-atmosphere general
408 circulation model.

409

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497

498 **8. Figure Captions**

499 **Figure 1. A.** Surface map of tropical cyclone diffusivity used as input for UVic climate
500 model simulations. Diffusivities are derived from satellite-based climatology developed
501 in *Sriver et al.* [2008]. **B.** Upper ocean temperature difference between the equilibrium
502 tropical cyclone simulation and the corresponding control for a background diffusivity
503 $K_v=0.2 \text{ cm}^2/\text{s}$. Difference is averaged over the uppermost 80 meters.

504

505 **Figure 2. A.** Difference in the total downward surface heat flux between the equilibrium
506 tropical cyclone simulation and the corresponding control for background diffusivity
507 $K_v=0.2 \text{ cm}^2/\text{s}$. Units are Watts/meter². **B.** Zonal integration of the heat flux difference
508 shown in Figure 2A. Units are Watts/meter divided by 10^8 . **C.** Difference in ocean heat
509 content, integrated from the surface to 500 meters depth, between the equilibrium tropical
510 cyclone simulation and the corresponding control for background diffusivity $K_v=0.2$
511 cm^2/s . Units are Joules/meter² divided by 10^9 . **D.** As in Figure 2C but integrated from
512 500 meters to 1500 meters depth.

513

514 **Figure 3. A.** Total northward ocean heat transport for control cases with varying
515 background mixing (K_v). Units are in Petawatts (1 Petawatt = 10^{15} Watts). **B.**
516 Difference between cases with TC mixing parameterization and the corresponding
517 control for each K_v .

518

519 **Figure 4. A.** Zonally-averaged potential temperature (color contours) in the Pacific basin
520 for the control case with background diffusivity $K_v=0.2 \text{ cm}^2/\text{s}$. Black contours represent

521 the zonally-averaged meridional velocity for the same region (solid – northward, dashed
522 – southward). Meridional velocity contour spacing is 0.5 centimeters/second. **B.** As in
523 Figure 4A but for the difference between the tropical cyclone case and the control.
524 Meridional velocity contour spacing is 0.05 centimeters/second.

525

526 **Figure 5 A.** Vertical profile of the zonally-averaged equatorial potential temperature
527 difference in the Pacific basin. Each curve represents the difference between the tropical
528 cyclone case and the corresponding control for the full range of background diffusivities.

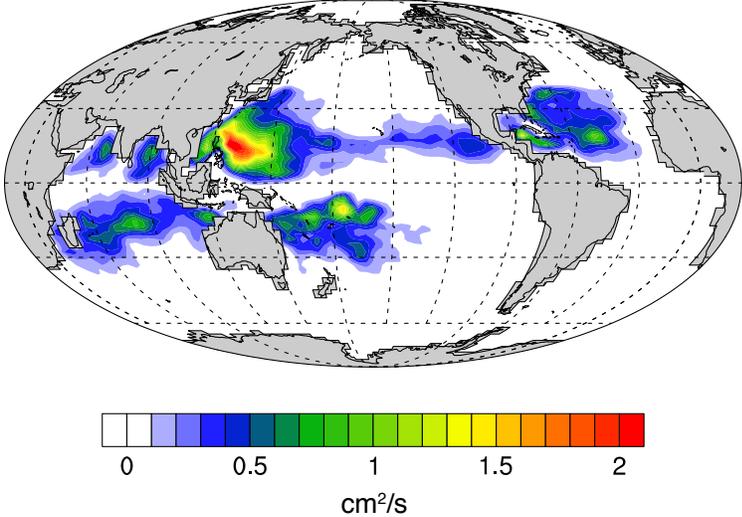
529 **B.** As in Figure 5A but for eastward velocity.

530

531 **Figure 6. A.** Equatorial transect of eastward velocity (color contours) in the Pacific basin
532 for the control case with background diffusivity $K_v=0.2 \text{ cm}^2/\text{s}$. Black contours represent
533 surfaces of constant temperature for the same region. Temperature contour spacing is 2
534 °C. **B.** As in Figure 6A but for the difference between the corresponding tropical
535 cyclone case and the control (color contours). The black contours denote the 20 °C
536 isotherm for the tropical cyclone case (dashed contour) and the control (solid contour).

Figure 1

A. Tropical Cyclone Diffusivity



B. Anomalous Ocean Temperature (upper 80 meters)

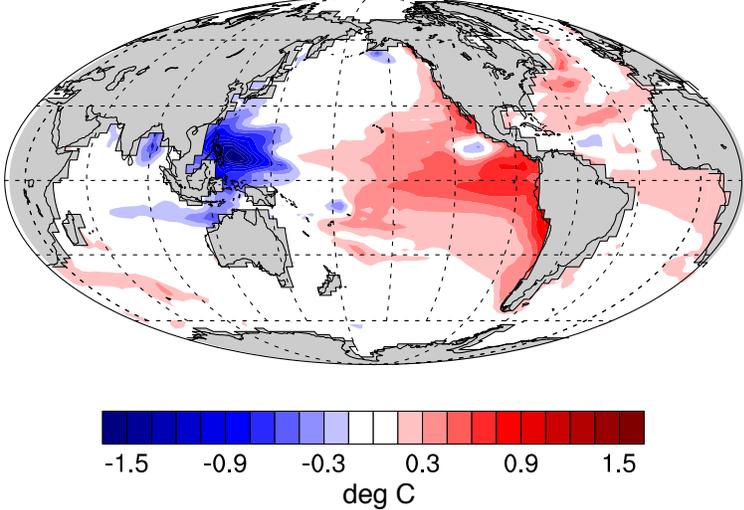


Figure 2

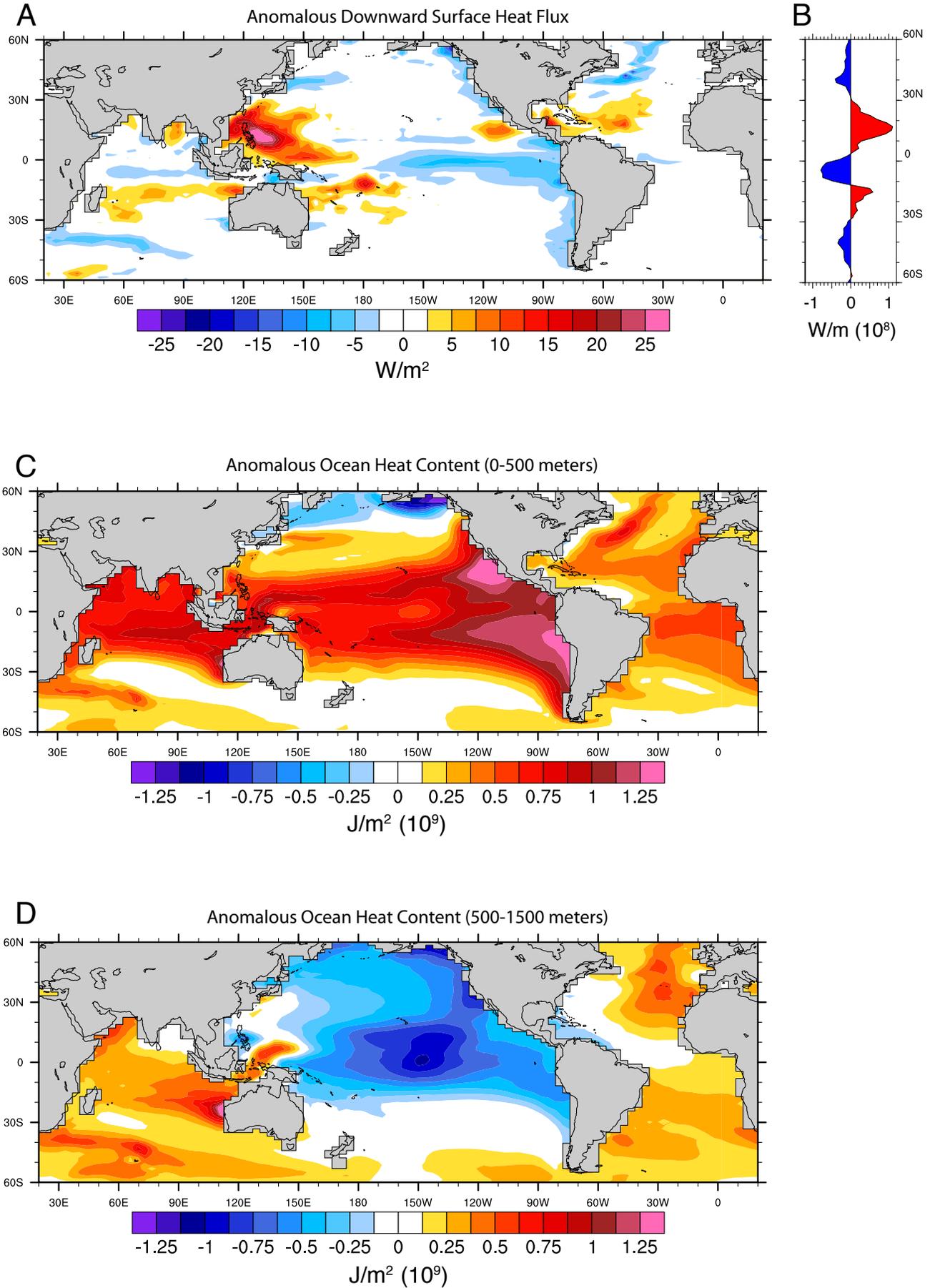


Figure 3

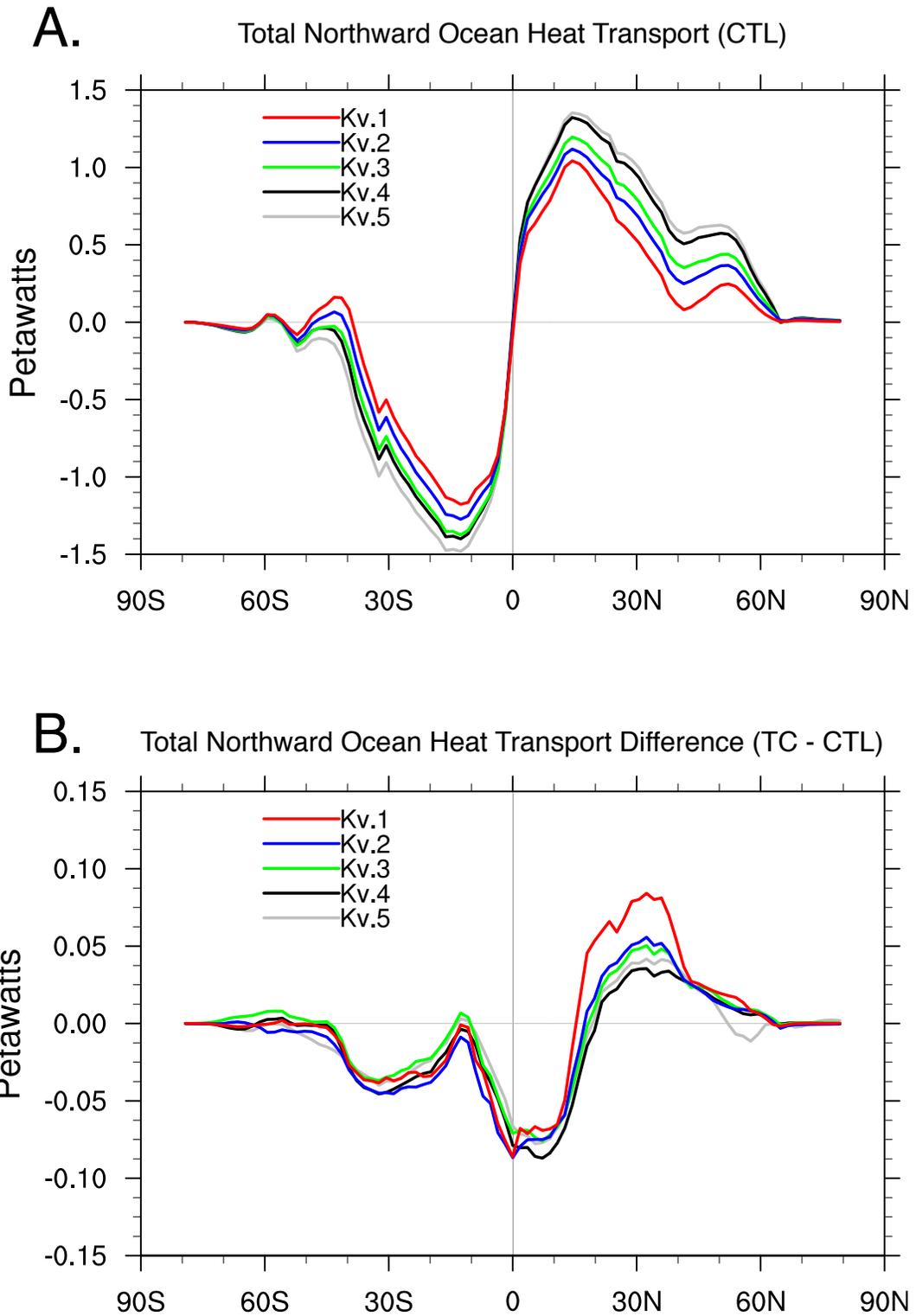


Figure 4

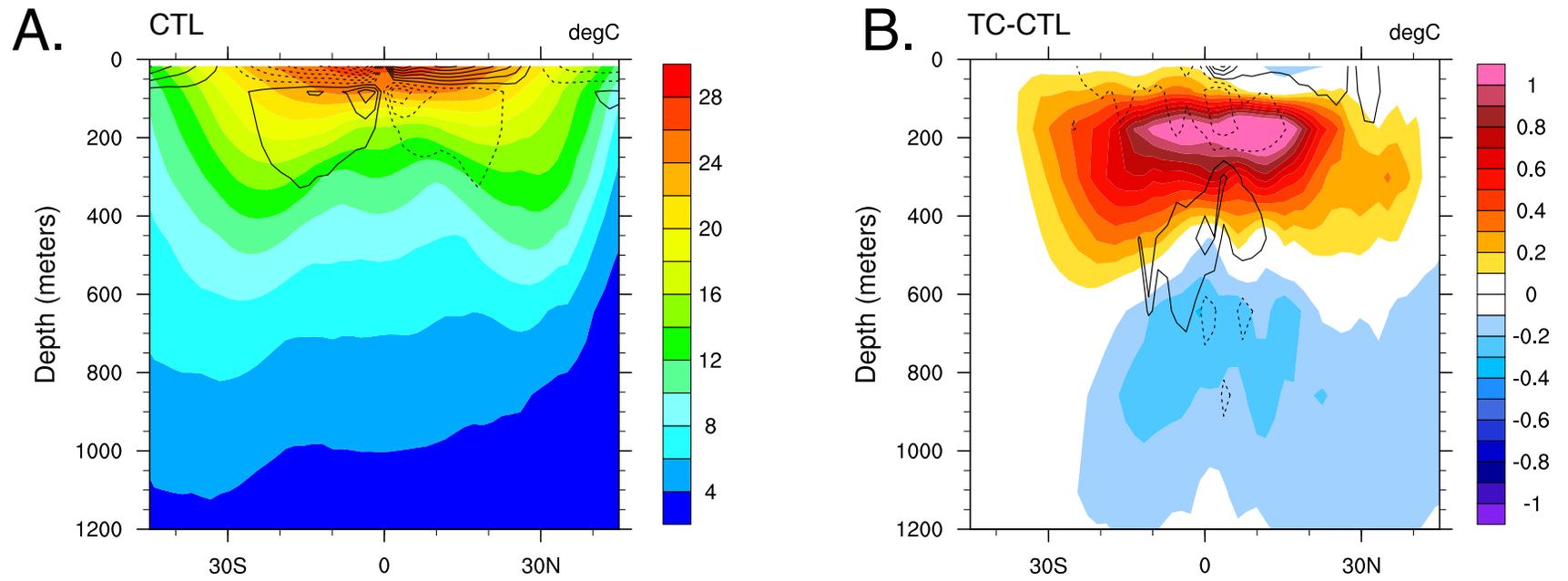


Figure 5

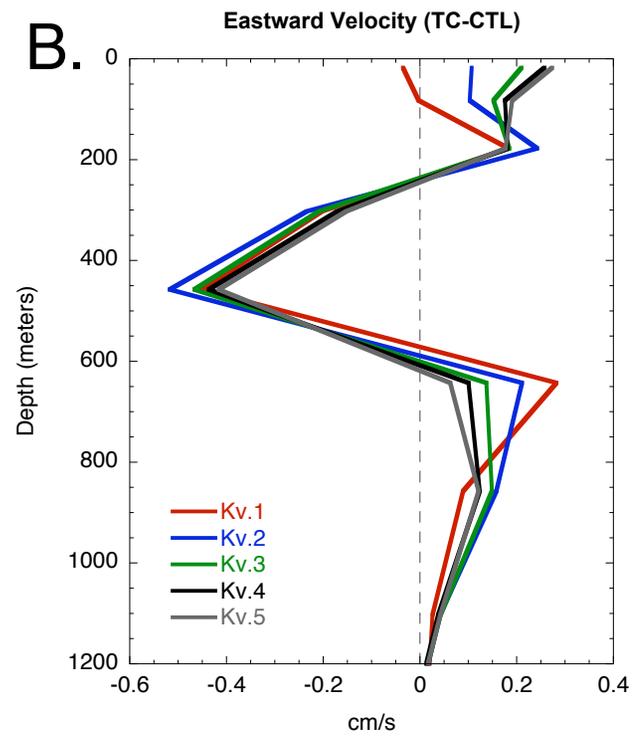
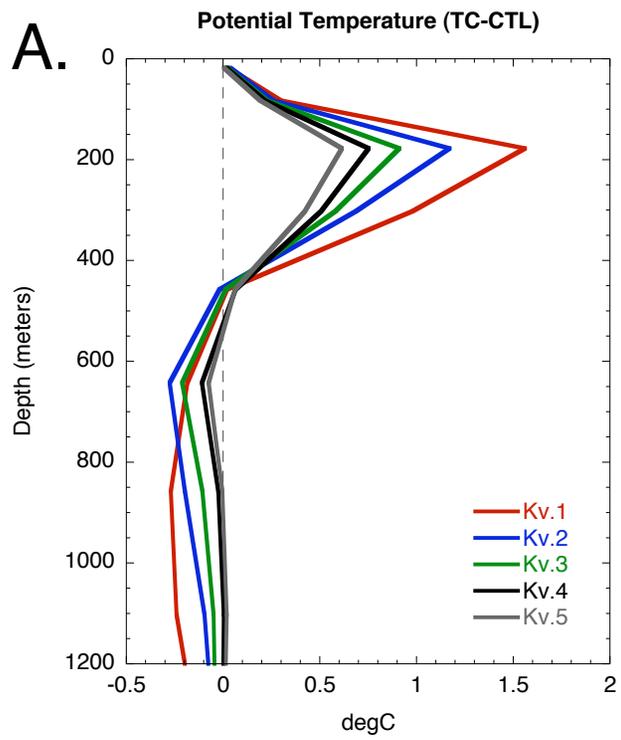


Figure 6

