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NATIONAL HURRICANE RESEARCH PROJECT

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REPORT NO. 71

Energy Generation and Flux Processes Associated With a Weakening Depression Over the Gulf of Mexico

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U. S. DEPARTMENT OF COMMERCE Luther H. Hodges, Secretary WEATHER BUREAU Robert M. White, Chief

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Energy Generation and Flux Processes Associated With a Weakening Depression Over the Gulf of Mexico

by

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ENERGY GENERATION AND FIUX PROCESSES ASSOCIATED WITH A WEAKENING DEPRESSION OVER THE GULF OF MEXICO

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ABSTRACT

The dissipation on September 21, 1963, of a depression over the Gulf of Mexico, has been investigated in terms of kinetic energy transport and generation processes. To evaluate these processes, data provided by research flights at the 800- and 600-mb. levels were used. Results show that the limited atmospheric region under study not only was subject to internal conversion of kinetic to potential and internal energy but also transported kinetic energy to the surrounding atmosphere.

1. INTRODUCTION

In recent years there have been several investigations of the energy and momentum budgets of tropical cyclones. In particular, decay of tropical storm circulation over land has been ascribed to surface friction (Gangopadhyaya and Riehl [1]), or removal of oceanic heat source (Miller [3]), or changes in kinetic energy exchange processes (Yanai [4]). The purpose of this study is to compute some of the processes involved in the loss of kinetic energy and weakening of the circulation of a depression which on September 21, 1963 was over the Gulf of Mexico. This depression was located off Key West on the morning of September 20, 1963. During the next 48 hours it moved northward over the Gulf of Mexico. At 1800 GMT on September 21, the surface center of the depression lay near latitude $28.5^{\circ}N.$, longitude $86.5^{\circ}W.$ Within the next few hours, the intensity of circulation in the lower troposphere diminished rapidly and by 0600 GMT, September 22, only a trough remained.

Cyclonic circulation was clearly evident from sea level to the 500-mb. surface throughout the life of the depression; a shear line could be identified at the 700 and 500-mb. levels during the weakening stage. As far as could be ascertained from available upper air data, the kinetic energy of the system remained nearly constant during its movement prior to 1800 GMT of September 21, from Key West to the northeastern Gulf of Mexico.

Between 1500 and 2100 GMT, September 21, two research aircraft of the National Hurricane Research Project were flown at 800 mb. and 600 mb. over the eastern Gulf of Mexico and the flight paths covered a region which included a major portion of the depression. Although no further flights were undertaken over the region and the data collected did not cover the entire system, it was felt that the availability of data at these two pressure levels justified

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an investigation of the magnitudes of the terms that enter into the kinetic energy budget. The computed values of the various terms should indicate the relative importance of the physical processes contributing to the rapid loss of kinetic energy of the depression and the adjacent atmosphere.

The analysis area is shown in figure 1.

2. THE KINETIC ENERGY BALANCE EQUATION

The scalar product of the horizontal wind ${\bf W}$ and the equation of motion leads to

$$\partial \mathbf{K}/\partial \mathbf{t} + \mathbf{V} \cdot \nabla \mathbf{K} + \omega \partial \mathbf{K}/\partial \mathbf{p} = - \mathbf{V} \cdot \nabla \mathbf{p} + \mathbf{V} \cdot \mathbf{F}$$
 (1)

where ω (= dp/dt) represents vertical motion, K (= $\mathbb{V} \cdot \mathbb{V}/2$) is the kinetic energy per unit mass, ϕ is geopotential, and [F is the frictional force per unit mass.

If we consider the mass, M, bounded by the perimeter of the grid and extending from 1000 mb. to 600 mb. in the vertical, integration of (1) over this domain yields

$$\frac{\partial}{\partial t} \int_{M} KdM = -\frac{1}{g} \iint \nabla \cdot K \nabla dAdp - \frac{1}{g} \int \frac{\partial}{\partial p} \left(\int_{A} K\omega dA \right) dp$$
$$-\frac{1}{g} \iint \nabla \cdot \phi \nabla dAdp - \frac{1}{g} \int \frac{\partial}{\partial p} \left(\int_{A} \phi \omega dA \right) dp - \frac{1}{g} \iint \omega \alpha dAdp + \int_{M} \nabla \cdot FdM$$
(2)

where A is area of the grid and **d** is specific volume of air.

The first and second terms on the right side of (2) represent net horizontal and vertical fluxes of existing kinetic energy at the boundaries of the domain, while the third and fourth terms represent boundary generation processes. The fifth, or the - $\omega \mathbf{a}$ term, represents conversion from potential and internal energy to kinetic energy through the ascent of warm air and descent of cold **air**. The last term represents generation and/or dissipation of kinetic energy by motions of smaller scale than are included in the analysis.

Yanai [4] has shown, in the case of a domain with lateral boundaries but extending from the surface to the top of the atmosphere, that the "value of $-\omega \alpha$ term is partly canceled with exchange of the mean potential energy." In the case of a domain with finite areal and vertical extent, it can be shown that

$$\overline{\mathbf{v}} \cdot \mathbf{\phi} \,\overline{\mathbf{w}} + \partial/\partial \mathbf{p} \ (\overline{\mathbf{\phi}} \,\overline{\mathbf{w}}) = \overline{\mathbf{v}} \cdot \mathbf{\phi}' \,\overline{\mathbf{w}} + \partial/\partial \mathbf{p} \ (\overline{\mathbf{\phi}' \,\overline{\mathbf{w}'}}) - \overline{\mathbf{w}} \,\overline{\mathbf{a}}$$
(3)

where $(\overline{}) = \frac{1}{A} \int (\overline{}) dA$ and we let $\phi = \overline{\phi} + \phi'$, $\omega = \overline{\omega} + \omega'$ and $\mathbf{q} = \overline{\mathbf{q}} + \mathbf{q}'$.



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Figure 1. - Base map used for analyses and grid points at which computations have been made. The crosses indicate grid points. The boundaries of blocks are shown in heavy lines.

If we substitute (3) into (2) and apply a two-dimensional form of the divergence theorem, we get

where L = perimeter of the grid,

v = horizontal wind component perpendicular to the vertical sides, positive outward,

 p_6 = upper pressure level (600 mb., in our case), p_0 = lower pressure level (1000 mb., in our case).

3. COMPUTATIONAL PROCEDURES

The basic data consisted of wind direction and speed, temperature, and geopotential heights at each of the grid points shown in figure 1 for the 1000,

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Figure 2. - Streamlines and isotachs (at intervals of 5 kt.) at the 1000-mb. surface, 1800 GMT, September 21, 1963.





800, and 600-mb. levels. Wind data for the 1000-mb. surface were obtained through conventional streamline isotach analysis of 1800 GMT surface observations from ships and coastal stations. Based on the mean height values over the Gulf of Mexico (Hebert and Jordan [2]), the sea level isobars were relabeled as contours for the 1000-mb. surface. Thus, a sea level pressure of 1012 mb. represented a 1000-mb. geopotential height of 108 m. with increments of 1 mb. in pressure amounting to height changes of 9 m. The wind and height patterns for the 1000-mb. surface are shown in figures 2 and 3. The circulation center and pressure minimum at the surface were located on considerations of continuity of position and intensity as indicated by previous synoptic charts of the District Meteorological Office at Miami.

The height, temperature, and wind data at 800 mb. and 600 mb. were obtained from the aircraft observations as follows. The u and v components, temperatures, and height values observed at 1-min. intervals during the flight were averaged over a flight period of 6 min. and were assigned to the aircraft position at the middle of this period. The average u and v components yielded the average wind direction and speed during the 6 min. For the scale of motions under study, it was not considered necessary to use a shorter smoothing time interval. The height values were, however, adjusted for hydrostatic and gradient wind consistency. Figures 4 to 7 show the wind, height, and temperature patterns at the 800-mb. and 600-mb. pressure surfaces.

The base map used for all analyses was a mercator projection with a scale of 1:5 x 10^{-6} at 22.5°N. The values of wind direction and speed, height



Figure 4. - Streamlines and isotachs (at intervals of 5 kt.) at the 800-mb. surface, 1800 GMT, September 21, 1963.



Figure 5. - 800-mb. geopotential and temperature analysis 1800 GMT, September 21, 1963. Contours are labeled in geopotential meters, isotherms are shown by dashed lines labeled in °C.



Figure 6. - Streamlines and isotachs (at intervals of 5 kt.) at the 600mb. surface, 1800 GMT, September 21, 1963.



Figure 7. - 600-mb. geopotential and temperature analysis 1800 GMT, September 21, 1963.

and temperature were read over a square array of 64 points (fig. 1) with a grid interval of 96 km.

From the interpolated values of wind direction and speed, the u and v components at the grid points were readily computed. The procedure adopted for computing vertical motion was as follows.

The values of horizontal divergence at the center of each 96-km. square were first computed kinematically. The scheme of computation for a square defined by the grid points (i, j), (i, j + 1), (i + 1, j + 1) and (i + 1, j), with i increasing northward, and j increasing eastward, was

$$b = \frac{1}{2x96km.} [(u_{i,j+1} + u_{i+1,j+1}) + (v_{i+1,j} + v_{i+1,j+1}) - (u_{i,j} + u_{i+1,j}) - (v_{i,j} + v_{i,j+1})]$$

$$(5)$$

The vertical motion profile at the center of each square was represented

as

$$\omega = A \ln \frac{1000}{p} + B \left(\ln \frac{1000}{p} \right)^2$$
(6)

where A and B are parameters which were determined by the divergence values at 800 mb. and 600 mb. Equation (6) implies $\omega = 0$ at 1000 mb., which is a reasonable lower boundary condition for ω . Differentation of (6) gives

$$D_{\text{SOO mb.}} = (-\partial \omega / \partial p)_{\text{SOO mb.}} = (A/p_{0}) + (2B/p_{0}) \ln 10/8$$

$$D_{\text{GOO mb.}} = (-\partial \omega / \partial p)_{\text{GOO mb.}} = (A/p_{0}) + (2B/p_{0}) \ln 10/6$$
(7)

These two equations when solved simultaneously yielded values for A and B and thus for ω as a function of pressure at each of the 49 central grid points (fig. 1). The divergence and vertical motion fields are shown in figures 8 to 11.

Area and line integrals involved in the computation of terms on the right side of (4) were approximated by summations of values representative of area and perimeter increments. Vertical integration was approximated by the trapezoidal rule.

With the above approximations and with input data consisting of wind direction, wind speed, height, and temperature for each grid point, it was possible to compute the terms on the right side of (4), with the exception of the dissipating term. It may be pointed out that the computed values of kinetic energy and flux and generation processes are only appropriate to motions on a scale comparable to the grid interval used.



Figure 8. - Horizontal divergence at 800 mb., 1800 GMT, September 21, 1963. Isopleths of divergence values are labeled in units of 10⁻⁵ sec.⁻¹ Shaded areas represent convergence



Figure 10. - Vertical motion ω at 800 mb. for 1800 GMT, September 21, 1963. Shaded areas represent upward motion. Isopleths are labeled in units of 10⁻³ mb. sec.⁻¹, or approximately equivalent to 1.4 cm. sec.⁻¹



Figure 9. - Horizontal divergence at 600 mb., 1800 GMT, September 21, 1963. Isopleths of divergence values are labeled in units of 10⁻⁵ sec.⁻¹ Shaded areas represent convergence.



Figure 11. - Vertical motion ω at 600 mb. for 1800 GMT September 21, 1963. Shaded areas represent upward motion. Isopleths are labeled in units of 10-3 mb. sec.⁻¹ or approximately 1.4 cm. sec.⁻¹

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Table 1. - Computed values of terms in the kinetic-energy balance equation, in units of 10^{18} ergs per second. Positive value indicates process acted to increase the kinetic energy of the region.

(K v _n)	(Kw)	(ø'v _n)	(¢'w')	(- ϣ , ≮ ')
-1.52	0.35	-3.13	2.81	-1.72

4. TERMS IN THE KINETIC ENERGY BALANCE EQUATION

Table 1 gives the values of the computed terms in (4) for the atmospheric region under study.

An immediately obvious feature of these results is the tendency for compensation between the horizontal and vertical fluxes of existing kinetic energy and between the boundary generation processes. The loss of kinetic energy due to conversion to potential energy is apparently not compensated for either by inward transport of kinetic energy through the boundaries or by generation due to the $\phi'v_n$ or $\phi'\omega'$ processes. Thus kinetic energy was not only being destroyed within the region but was also being lost to the surrounding atmosphere. The net depletion of kinetic energy due to these flux and conversion processes amounts to 3.21×10^{18} ergs sec. ⁻¹ for the region, or 622.6 ergs cm. ⁻² sec. ⁻¹ for the 1000- to 600-mb. layer. The actual kinetic energy in the region was computed to be 27.71 $\times 10^{22}$ ergs at 1800 GMT. Hence, if only these processes had been operating, kinetic energy of the region would have been reduced to zero in about 24 hours, unless radical changes occurred in some of these processes.

It is more likely that the frictional dissipation action was actually small in the region while certain eddies, not detected on the scale used in the present analysis, produced kinetic energy. The rate of production was not, however, high enough to overcome the loss due to processes on a larger scale. Since the scale of motions included in the present analysis was of the order of 200 km., the eddies which might have produced kinetic energy in the region were perhaps of convective scale.

. 5. PRODUCTION OF KINETIC ENERGY IN THE DEPRESSION AND ITS SURROUNDINGS

It is of interest to investigate the nature of kinetic energy production in the portion of the region covered by the depression as related to production in other parts of the region. For this purpose, the analysis area was divided into four blocks (fig. 1). Within each block the generation term $-\Psi \cdot \nabla \phi$ of (1) was evaluated at each pressure level, using finite differences over the grid intervals. The results given in table 2 show that kinetic energy was generated largely in the surface layer and destroyed in the upper layers. It must be pointed out that since the number of observations at the 1000-mb. level is small, the values of the generation term for this level are likely to contain more uncertainties than those for the higher levels.

Block	1000 mb.	800 mb.	600 mb.
I	1.76	-0.57	-0.91
II	2.12	-7.09	- 3 .67
III	-1.32	2.17	0.82
IV	0.94	-2.38	- 2.65

Table 2. - Work term $(-\mathbf{W} \cdot \mathbf{\nabla} \phi)$, in units of 10^{15} ergs sec.⁻¹

As in the case of Blocks II and IV, which incorporate the main stream of southwesterly flow to the east of the depression, kinetic energy was generated. in the area of the depression (Block 1) in the low levels but was destroyed in the upper levels. In view of the nearly symmetric flow within Block I and the small wind speeds around the depression, it would seem reasonable to neglect energy changes in Block I due to horizontal flux of kinetic energy. Since upward motion was predominant at the 800- and 600-mb. levels over the same area, any contribution due to vertical flux of kinetic energy was certainly negative. Further, the work term $(-W \cdot \nabla \phi)$, was such that kinetic energy was destroyed in the upper levels. The small positive rate of generation allowed the kinetic energy in the lower levels to be dissipated easily by frictional stresses at the surface. The depression was thus characterized by a net loss of kinetic energy which could not, perhaps, be overcome by the possible generation of energy by small-scale eddies.

6. GENERAL CONCLUSIONS

For the portion of the lower troposphere over the Gulf of Mexico bounded by the 1000- and 600-mb. pressure surfaces, flux and generation terms in the kinetic energy balance equation have been evaluated. Since the region under study was open, energy exchanges with the surrounding atmosphere were considered.

The system was not only subject to internal conversion of kinetic to potential and internal energy but was also acting as a source for the surrounding atmosphere. Such a state presumably resulted from the choice of this particular analysis region which covered only a portion of a short wave. Further, estimates of the so-called frictional term were not possible in the absence of data on actual kinetic energy changes in the region. It is likely that motions on a scale not detected in the present analysis were such that they actually generated kinetic energy.

This study has been based on data at three pressure levels in the lower troposphere and over a limited area at one particular time. The analysis region could not be chosen large enough for computations to be made over the entire extent of the cyclonic flow in the lower troposphere. Moreover, a certain amount of uncertainty in the analyses was inevitable in areas not actually covered by the aircraft flights. Nevertheless, it is of interest to note that the kinetic energy production rate (represented by the integral of $-\mathbf{W} \cdot \mathbf{\nabla} \phi$) of -1.863×10^{18} ergs sec.⁻¹ compares favorably with the total rate of -2.045 x 10^{18} ergs sec.⁻¹ represented by the sum of the generation terms in table 1. There is need for further studies similar to the current one in order to gain more insight into the processes associated with changes in intensity of cyclones.

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